THE DESERVATION AND ANALYSIS OF LUNAR OCCULTATIONS OF STARS WITH AN EMPHASIS ON IMPROVEMENTS TO DATA ACQUISITION INSTRUMENTATION AND REDUCTION TECHNIQUES

Bу

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A DISSERTATION PRESENTED TO THE GRADUATE SCHOOL OF THE UNIVERSITY OF FLORIDA IN PARTIAL FULFILLMENT OF THE REQUIREMENTS FOR THE DEGREE OF DOCTOR OF PHILOSOPHY

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This work is dedicated to

Rose Wenig

and

Alice Schneider,

my late, beloved grandmothers.

If the stars should appear one night in a thousand years, how would men believe and adore; and preserve for many generations the remembrance of the city of God which had been shown!

"Emerson", Nature; Addresses, and Lectures

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THE OBSERVATION AND ANALYSIS OF LUNAR OCCULTATIONS OF STARS WITH AN EMPHASIS ON IMPROVEMENTS TO DATA ACQUISITION INSTRUMENTATION AND REDUCTION TECHNIQUES

Bу

Glenn H Schneider

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Chairman: John P. Oliver Major Department: Astronomy

A program of observation and analysis of lunar occultations was conceived, developed, and carried out using the facilities of the University of Florida's Rosemary Hill Observatory (RHO). The successful implementation of the program required investigation into several related areas. First, after an upgrade to the RHO 76-cm. reflecting telescope, a microprocessor controlled fast photoelectric data acquisition system was designed and built for the occultation data acquisition task. Second, the currently available model-fitting techniques used in the analysis of occultation observations were evaluated. A number of numerical experiments on synthesized and observational data were carried out to improve the performance of the numerical techniques. Among the numerical methods investigated were

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solution schemes employing partial parametric adjustment, parametric grouping into computational subsets (randomly and on the basis the correlation coefficients), and preprocessing of the observational data by a number of smoothing techniques for a variety of noise conditions. Third, a turn-Key computational software system, incorporating data transfer. reduction, graphics and dislplay, was developed to carry out all the necessary and related computational tasks in an interactive environment.

Twenty-four occultation observations were obtained during the period March 1983 to March 1984. The observational data and the solutions resulting from the subsequent reductions are presented graphically and tabularly for each of the occultation events. Several angular diameter determinations were made. Among those of particular interest were 32 Librae (12.1, +/- 1.9 milliseconds of arc), 1 Geminorum-B1 (5.9. +/- 0.8 milliseconds of arc), and X07598 (5.5. +/- 2.0 milliseconds of arc). The visual/spectroscopic binary 1 Geminorum was discovered to have a fourth. previously undetected, component. Two other stars, X13534 and X13607, were found to be binary with compainions closer than 15 milliseconds of arc. Previously unknown faint companions were discovered for ZC1221 and ZC0126. Times of geometrical occultation for all the events (including the secondary components of the binary systems) were reported as part of a cooperative astrometric project to the International Lunar Occultation Center.

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CHAPTER I INTRODUCTION

Lunar Occultations: A Historical Synopsis

When the moon, as a result of its orbital motion, moves in front of a star as viewed by an Earth-based observer, the light from the star is obscured. Such an event is known as a lunar occultation. At the time of the star's disappearance (or reappearance) the moon's limb is seen to move across the stellar disc. MacMahon (1909) proposed that the angular diameter of the occulted star could be determined by measuring the time interval of lunar limb passage across the star. Immediately thereafter, Eddington (1909) correctly pointed out that this was not possible, as the star's light would be diffracted by the moon's limb, hence the problem could not be approached from the standpoint of simple geometrical optics. He noted that the time-scale of the disappearance phenomenon is essentially unchanged as a function of stellar diameter.

Williams (1939) showed that while the time-scale of the event is not affected by the angular diameter of the star the diffraction intensity as a function of time most certainly is. The time-scale of variation of the diffraction fringes projected on the Earth's surface resulting from a lunar occultation is on the order of tens of milliseconds, thus

requiring fast photometric observations. Whitford (1939), reported on observations of occultations of Nu Aquarii and Beta Capricorni using a photocell with a Cesium-Oxygen-Silver cathode on the 100-inch telescope. Neither of these events showed any deviation in the diffraction pattern from a point source, as his instrumental detection limit was approximately 5 milliseconds-of-arc. Yet the foundation for a powerful new technique for the acquisition of fundamental astronomical information, i.e. stellar diameters, was laid.

Over the ensuing three decades additional photoelectric occultation observations were carried out. The first angular diameter measurement was reported by Evans (1951) for the star Antares. This was followed, also by Evans (1959), with the determination of the angular diameter of Mu Geminorum. Over the next two decades other observations had been made. and additional theoretical work on the extraction of information from lunar occultation observations progressed. It was not until the advent of electronic computers and reliable fast photometric equipment that occultation observations could truly begin to be exploited. In a now classic series of papers by Nather and Evans (1970), Nather (1970), Evans (1970 and 1971), and Nather and McCants (1970) the theoretical and observational aspects of lunar occultations are discussed in detail. The methods reported in these papers serve as the foundation for modern investigation of lunar occultations.

Only in the last few years, with the revolution in both microcomputer and opto-electronic technology, have the tools essential to bringing the observation and analysis of lunar occultations come to fruition. The problems are still many, but the instrumental hurdle, at least, may now be cleared.

Information Which May Be Learned From the Analysis of Lunar Occultations

The analysis of the lunar occultation intensity curve, obtained from a fast photoelectric record of an occultation event can yield, in principle, a wealth of information. The degree to which any observation can be exploited depends upon a large number of variables. The geometry of the relative position of the moon and the star, the quality of the sky during the observation (seeing and scintillation), the physical nature of the source, and the response of the instrumental system, to mention only a few, can help or hinder the discovery of information hidden in the intensity curve.

In the case of an occultation of a single star the angular diameter of the star can be determined. Coupled with either a knowledge of the stellar parallax, or the V and R stellar flux (Barnes and Evans, 1976) this angular measurement can be transformed into an actual linear diameter. The observational techniques for direct measurement of stellar diameters are severely limited. Speckle interferometry (Lohman and Weigelt, 1980), Phase Correlation Interferometry (Brown, 1968), and Michelson

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Interferometry (Brown, 1980), are the only other currently available techniques. All of these are restricted, instrumentally, to the measurement of diameters of very bright stars.

The size of extended non-stellar sources such as the emission shells of Be stars can be determined (White and Slettebak, 1980) as was investigated in the case of Zeta Tauri (Schmidtke and Africano, 1984). In ideal cases, with multiple observations the brightness distributions of such sources could be found.

The angular optical resolution achievable through occultation observations, on the order of milliseconds-of-arc, often leads to the discovery of stellar duplicity of stars previously thought to be single. An accurate determination of the separation of the components in a binary (Evans, 1971) or multiple system (Evans et al., 1977) can be found from simultaneous observations from more than one site (or a projected separation from a single observation). The individual brightnesses of otherwise "unresolved" binaries, or multiple systems, can be determined.

The field of chronometry is dependent upon, and enhanced enormously by, the precise measurement of the times of occultation events. These event timings lead to an accurate determination of the moon's longitude from which Ephemeris Time is derived. Time intervals are easily obtained with a precision of one part in 10 trillion through the use of

atomic clocks. The observation of dynamical phenomena must have a zero point reference to couple dynamical events to atomic time (Van Flandern, 1974). Hence the accumulation of a database of accurate timing measurements of occultation events (dynamical phenomena) is essential in helping provide corrections to Co-ordinated Universal Time (C. U. T.).

The determination of stellar positions, the fundamental concern of the field of astrometry can benefit from occultation observations. The predicted time of an occultation event for a given topocentric location depends upon a Knowledge of the moon's actual longitude, the geometry of the point of contact of the lunar limb, and the position in a given co-ordinate system of the star under study. Prediction errors in the times of occultation can lead to an improved Knowledge of the positions of stars (Van Flandern, 1975). In some cases, gross errors in the predicted times of occultation can be attributed to positional errors in the stellar catalogs.

Goals of the Program of Occultation Observation

The primary intent of this program was to concentrate the investigation in areas which could have lead to astrophysically significant, or interesting results. While it could not have been determined, a priori, which stars or stellar systems would yield the most useful information, a judicious choice of program objects was called for. Many observational selection effects and instrumental limitations

dictated the types of stars for which the occultation method would be fruitful. The general nature of these constraints has been addressed by Taylor (1966) and Ridgway (1977). Hence, in order to be considered for the observing program, candidate stars had to meet a variety of selection criteria. Unfortunately this left only a handful of stars each year to be considered for this investigation.

As a result of the stringent limitations a great deal of attention was paid to improving observational techniques and the development of new microelectronic data acquisition instrumentation. Once implemented, improved instrumentation allowed relaxation, to some degree, of the observational constraints placed on candidate star selection, thereby increasing the number of stars which were available for study. In addition, several non-traditional methods of data reduction were tried, which in some cases proved to be advantageous over the more conventional procedures.

The primary emphasis of the program was to extract as much information as possible from lunar occultation observations. The determination of stellar angular diameters, the discovery of unsuspected stellar duplicity or multiplicity, and the elucidation of the parameters of double or multiple systems were paramount. Of secondary importance, but an item to which no concessions were made during evolution of the program, was the accurate measurement of the times of occultation events. These measurements were reported, as part of an international co-operative project.

to the United States Naval Observatory, and the International Lunar Occultation Center.

An outgrowth of the lunar occultation program has been the observation and analysis of occultations of stars by asteroids. Such observations typically yield information about the size and shape of the occulting body, as well as better astrometric positions of the occulted stars. Over two thousand asteroids are known and have been catalogued (Bender, 1979), and several hundred have orbits determined with a sufficient degree of precision to allow topocentric predictions of asteroidal occultations to be carried out well in advance of the anticipated event. Hence, the third aspect of this investigation of high speed occultation photometry extended the observational domain to asteroidal, as well as lunar occultations.

CHAPTER II INSTRUMENTATION

Optical Equipment

The Seventy-Six Centimeter Telescope

Location and description. All observations, unless otherwise noted, were carried out at the University of Florida's Rosemary Hill Observatory (RHO) in Bronson, Florida, utilizing the seventy-six centimeter Tinsley reflecting telecope. This telescope is located at Latitude of 29 degrees 23 minutes 59.4 seconds North, Longitude 82 degrees 35 minutes 11.0 seconds West, at an altitude of 44 meters above mean sea level, based on the 1927 North American Datum. The primary mirror has a focal ratio nominally of f/4. In its Cassegrain configuration the telescope provides a nominal focal ratio of f/16. Figure 2-1 shows the seventy-six centimeter telescope and the associated optical and electronic equipment employed for the observations of lunar occultations.

<u>The telescope light baffle</u>. Observing immediately adjacent to the bright lunar disk increases the background sky illumination considerably. For a given spectral region there is nothing which can be done to enhance the signal to noise ratio considering noise due to atmospherically scattered light. It was found, however, that the Cassegrain



Figure 2-1. The seventy-six centimeter Tinsley relecting telescope at the University of Florida's Rosemary Hill Observatory.


optical baffling for the seventy-six centimeter reflector could be improved upon greatly. A redesign of the optical baffle permitted a reduction in light scattered by the telescope optics and supporting structure. In addition, a tighter baffle system enabled the rejection of off-axis rays, preventing them from reaching the focal plane.

The primary design considerations for a new baffle system were two-fold. First, any baffles and/or field stops would have to be easily removable, so as not to impact on the telescope configuration required by concurrently running observing programs. Second, the unvignetted field of view at the focal plane had to marginally exceed the field obtained with the available diaphragm while minimizing extraneous light. An Astromechanics dual channel photometer was the primary instrument to be used for observing lunar occultations. This instrument has a maximum diaphragm opening corresponding to a field of view with a diameter of thirty two seconds of arc.

Many designs were considered by constructing models of the optical path, and ray tracing paraxial and marginal rays. It became quite obvious early in this investigation that no single tube/field stop design would be satisfactory for the task. The new baffle system would have to consist of two major elements: a tube with concentric annular field stops (rings), and an annular ring placed around the Cassegrain secondary mirror. A schematic representation of the optical path is presented in Figure 2-2.





To satisfy the first design objective a system of field stops was built into a baffle tube identical at the mounting base to the existing tube. This would allow easy change-over to either the new or old baffle tube. The Cassegrain secondary ring would be be installed on the already existing locking pins on the secondary mirror baffle, which serve to hold the mirror cover in place. The new baffle tube, with the inner concentric rings removed to show its construction, is shown in Figure 2-3 along with the secondary ring.

Since the position of the focal plane varies with respect to the fixed position of the telescope superstructure and baffle assembly as a function of temperature, some degree of tolerance had to be allowed in the placement and size of field stops in the baffle tube. The system as designed would allow a one minute of arc unvignetted field of view at the focal plane at 15 degrees Celsius. This rather liberal tolerance was felt prudent considering the wide variation of climatic conditions experienced in northern Florida. Table 2-1 gives the specifications for the spacing and sizes of the inner ring stack. The secondary baffle ring has an outside diameter of 11.375 inches.

To actually perform the ray tracing, the optical path of the telescope had to be determined. Although the blueprints from Tinsley Laboratories claim the primary is an f/4 paraboloid this had to be tested as the engineering specifications did not necessarily represent the state of the completed telescope.

Figure 2-3. The occultation light-baffle tube and ring.



TABLE 2-1 SPECIFICATIONS FOR THE OPTICAL BAFFLE TUBE

Ring	Hole Diameter	Ring	Hole Diameter
1	3.56	5	2.24
2	3.23	6	1.91
3	2,90	7	1.67
4	2.57	8	1.24
Outsi Outsi Thick Spaci Bevel	de diameter or r de diameter of r ness of rings 1- ng between rings angle for all 8	ing 1: ings 2- 8: : holes:	5.000 7: 4.812 0.0625 5.938 45 degrees
Note:	All linear measu	rements	are in inches.

Using photographic plates of regions routinely monitored by the quasar research program, the plate scale of the primary mirror was found to be 67.065 (S.E. 0.06) seconds of arc per millimeter, corresponding to a focal length of 117.505 (S.E. 0.11) inches. The scale was determined by Pollock (1980) by measuring the positions of the quasar 0922+14 (denoted "0") and the star SAO 098559 (denoted "S") for six plates taken over a period of seven years. Table 2-2 shows the plate measures used in determining the true scale of the primary mirror. The distance to the focal plane from the front surface of the Cassegrain secondary was measured at 15 degrees Celsius. From this the focal length of the Cassegrain system was found to be 1205 centimeters.

The assembled baffle tube is inserted through the central hole in the primary mirror and is screwed into place. The old, wide-field baffle tube is removed by grabbing the outside end and twisting sharply counter-clockwise. This will loosen its seating. When unscrewing the old baffle tube one hand is kept below the lower end of the tube to prevent it from hitting the primary mirror when it is completely released. When screwing the occultation baffle tube into

TABLE 2-2

DETERMINATION OF THE PRIMARY MIRROR FOCAL LENGTH

PLATE DATE	StoQ (mm)	Q to S	MEAN DIFF (mm)	SCALE
4/2-3/72	9.212	9.215		
	12.988	12.987		
	9.208	9.217		
	12.982	12.991	3.774	67.01
1/30-31/73	10.024	10.032		
	13.792	13.793		
	10.026	10.028		
	13.796	13.796	3.767	67.14
2/27-28/76	7.940	7.988		
	11.740	11.758		
	7.979	7.985		
	11.750	11.753	3.769	67.10
2/27-28/76	2.782	2.791		
	6.555	6.567		
	2.788	2.790		
	6.560	6.567	3.774	67.01
2/13-14/77	3.995	4.000		
	7.775	7.776		
	4.000	4.005		
	7.769	7.778	3.774	67.01
1/28-29/79	19.972	18.972		
	22.749	22.736		
	18.971	18.975		
	22.735	22.742	3.768	67.12
			MEAN	67.065
				(0.06)

position, care must also be taken to assure the tube is not being cross-threaded. To switch back to the old tube (which must be done if using the infra-red photometer available at RHO) the process is reversed. The Cassegrain secondary baffle ring slips over the end of the secondary containment cylinder after it has been aligned with the cylinder's three positioning pins. A small rotation will secure the ring position insets against the pins.

Photoelectric Photometers

The Astromechanics photometer. Unless otherwise stated, the photometer used throughout this investigation was a dual channel instrument manufactured by Astromechanics. This instrument splits the light path into two beams by means of dichroic filters so two wavelengths can be monitored simultaneously. Though the instrument can be used in dual channel mode, observations of lunar occultations obtained thus far at RHO have been observed only in one color. The photometer employs two dry ice cooled photomultiplier tubes (PMT's). The PMT in channel-1 is a 6256S, while channel-2 uses a red sensitive 9684. In most cases 1200 Volts DC was applied to the PMT used. This photometer has a number of different available diaphragms. The linear diameters and angular field sizes of these diaphragms, along with their respective letter designations are listed in Table 2-3.

TABLE 2-3 DIAPHRAGM DESIGNATIONS, LINEAR AND ANGULAR FIELD SIZES

그 문 눈 날 것 같 같 좀 잘 있는 것 두 안 주 것 같 것 같 것 같 것 같 것 같 것 같 것 같 것 같 것 같 것							
Letter Designation	Diameter (mm)	Field (arc-secs.)					
G	1.98	32.5					
н	1.52	25.0					
I	0.93	15.2					
J	0.51	8.3					
		AND					

Optical filters. Occultation observations made with the Astromechanics photometer employed Johnson V and B filters. as well as intermediate bandwidth yellow and blue interference filters. One inch diameter interference filters were obtained from Pomfret Research Optics, and are designated "y" and "b" respectively. The "y" filter, Pomfret part number 20-5400-1, has a peak spectral transmission at 5400 Anostroms and a Full Width at Half Maximum (FWHM) of 100 Angstroms. The "b" filter, Pomfret part number 20-4700-1 has a peak spectral transmission at 4700 Angstroms and a FWHM also of 100 Angstroms. These filters were selected in spectral regions for which M and K stars are relatively free of major absorption features. Of course, late type stars are riddled with a myriad of spectral lines. Hence the choice of these particular filters was somewhat of a compromise. Spectra typical G. K. and M. stars presented by Keenan and McNeil (1976) were examined, and on average were found least plaqued with absorption lines at wavelengths of 4716 and 5408 Anostroms. These would have been the ideal central wavelengths for selected filters, but the cost of custom made filters was prohibitive. The filters procured were selected

to be as close to these wavelengths as possible from a stock list. The H-Beta line at 4861 Angstroms is outside of the "b" filter passband. While other lines are found at 4716 and 4670 Angstroms, the integrated passband is less subject to absorption losses than adjacent regions. The TiO band at 5448 Angstroms enters into the "y" filter passband, but it is centered close to the lower half-power point.

While an actual set of narrow band filters was not available, a digital spectrum scanner employing Ebert-Fastie optics (Parise, 1978) is part of the standard equipment at RHO. The scanner can be used in a non-scanning mode as a variable-passband tunable filter. This in fact was done with great success in observing the occultation of Zeta Tauri in the passband of its H-Alpha emission.

<u>The portable photometer</u>. A small, lightweight photometer employing a 1P21 photomultiplier tube and built-in Johnson U, B, and V filters was used exclusively for events observed from remote sites. This instrument is discussed in detail by Chen and Rekenthaler (1966).

Data Acquisition Electronics

The SPICA-IV/LODAS System

Design criteria for a new SPICA. The concept of a Small Portable Interactive Computer for Astronomy (SPICA) was first conceived by Dr. John P. Oliver. The first SPICA system was implemented on a KIM-1 computer, and is the precursor to the three generations of SPICAs which have followed. The common

thread linking the first SPICA to the latest version, the SPICA-IV, is the use of a 6502 microprocessor. Though each major upgrade to the SPICA systems has involved more hardware an effort has been made in SPICA-IV to retain portability, or at the very least transportability.

The Lunar Occultation Data Acquisition System (LODAS) is the software control program designed to carry out the task of fast photometric data acquisition on SPICA-IV. It is, in actuality, inaccurate to say that LODAS was designed for SPICA-IV, or SPICA-IV for LODAS. The system requirements were such that the hardware and software grew together in a complementary fashion. The major design criteria for SPICA-IV/LODAS are listed in Table 2-4.

> TABLE 2-4 DESIGN CRITERIA FOR SPICA-IV/LODAS

- Data acquisition rates up to and including 1 KiloHertz must be supported.
- The system must support a minimum of two simultaneous data acquisition channels.
- Memory space must be provided to hold a minimum of two seconds of data in each channel, at a rate of 1kHz.
- 12-bit sample resolution should be used, to give a large dynamic range and eliminate last minute gain switching.
- 5. The system should retain easy transportability.
- The system must function in the abscence of a disk drive, or a disk operating system.
- A user friendly command structure and display must be implemented.

8. The control program must reside in Read Only Memory.

In all cases these criteria were met, and in the first three cases they were exceeded.

The SPICA-IV digital electronics. The heart of the SPICA-IV/LODAS system is an Advanced Interactive Microcomputer, model AIM-65, manufactured by Rockwell International. The AIM-65 has proven to be an invaluable design and development tool for the LODAS system as well as for several other astronomical data acquisition systems implemented at the Rosemary Hill Observatory. The AIM-65 is an 8-bit microcomputer with sixteen address lines incorporating a 6502 microprocessor chip. Up to 4-kilobytes of Random Access Memory (RAM), in the form of paired 1K-by-4 bit chips (i.e. 2114's) can be accommodated on the AIM board. A 6522 Versatile Interface Adaptor (VIA), which is a programmable chip holding 16 bidirectional I/O lines, four control lines, and two timers, serves as an interface to the "outside world" through an expansion connector on back of the AIM board. A standard ASCII keyboard, a 20 character alphanumeric LED display, and a thermal printer are provided for user I/O. The AIM-65 accomodates 24-kilobytes of ROM space in five 4-kilobyte Read Only Memory (ROM) sockets. memory mapped in the areas of \$B000 to \$FFFF. The AIM-65 operating system is normally resident in two ROM's occupying the uppermost 8-kilobytes of address space.

The SPICA-IV/LODAS system uses three boards manufactured by Micro-Technology Unlimited (MTU) to expand its RAM memory and to support peripheral devices. The first of these is a

16-Kilobyte dynamic RAM board, model number K-1016 (MTU, 1979). This board is address assignable only on 8-Kilobyte boundaries. Since the AIM-65 board is designed to hold 4-Kilobytes of RAM, addressing the RAM board on 8-Kilobyte boundaries would leave a 4-Kilobyte hole in the system address space. While this in itself is not a problem, it will be seen that all available 64-Kilobytes of AIM-65 address space must be used in configuring the system to meet the design criteria. This 4-Kilobyte hole would then require a 68-Kilobyte address space which the AIM-65 does not support. With this in mind the 2114 RAM chips were removed from the AIM board and the lower 16 Kilobytes (address range \$0000 to \$3FFF) of system RAM reside contiguously on the dynamic RAM board.

A major design consideration was for LODAS to be able to function even if there were a hardware failure of either the disk drive or its controller board. A Shugart model 801, 8-inch floppy disk drive is used for primary data storage. The second MTU board is a Double Density Disk Controller (DDDC), model number K-1013 (MTU, 1980a), used in conjunction with the disk drive. The Channel Oriented Disk Operating System (CODOS) is distributed by MTU along with the DDDC board. The CODOS software provides many utility functions as well as a Service Call Processor and a Visual Memory Terminal driver program (MTU, 1981) which enhance the overall utility of the SPICA system.

The CODOS and its associated programs occupy address space in the range of \$5000 to \$5FFF, and \$8000 to \$9FFF. The DDDC board also provides an additional 4-kilobytes of RAM which is mapped in the address range \$4000 to \$4FFF. Although the disk/CODOS system is an integral part of SPICA-IV/LODAS, it is modular. Both the disk drive and the DDDC board may be removed from the system without impairing the data acquisition capability of the LODAS.

The final MTU board is a bit-mapped dynamic 8-Kilobyte "visual" RAM high resolution display, model number K-1008 (MTU, 1980b). The display has a horizontal resolution of 320 dots (40 bytes) and is 192 scan lines in length. Thus 61440 bits may be individually controlled on the display within a 192-by-320 dot matrix. The hi-bit of the lowest byte on the board's address space is displayed on the upper-left of a Video Display Unit. The lo-bit of the highest byte occupies the extreme lower-right corner of the display. This board is memory mapped for the address space \$6000 to \$7FFF.

All three MTU boards were designed with on-board voltage regulators to derive +5 Volts from an unregulated +8 Volt supply. Since the SPICA-IV power supply provides regulated +5 Volts, the +8 Volt regulator on each MTU board was bypassed, and regulated +5 Volts distributed to each board. Other than this and the addition of a BNC connector on the visual memory board, the MTU boards are unmodified.

An additional board, referred to as the Analog-to-Digital Converter/Clock (ADCC) board used in the SPICA-IV/LODAS system, was built by Oliver specifically for high speed occultation photometry. However, the devices provided on this board have been conveniently memory mapped and are available for other application programs running in any SPICA system at RHO. An examination of the SPICA-IV/LODAS system memory map, as shown on the LODAS assembly listing in Appendix A , reveals that the system memory space would normally be fully occupied with no addressable locations available for the placement of this board. The AIM-65 reserves the address space \$A000 to \$AFFF for its own on-board devices, but only a small portion is actually address decoded (Rockwell International, 1978). A minor modification on the AIM-65 board was made to free up normally undecoded address space within this range, and has subsequently been made to all SPICA-III and SPICA-IV computer systems in use at the RHO.

The ADCC was laid out on an MTU prototyping board, K-1020 (MTU, 1980c), and holds a real time clock and three analog-to-digital converters (A-to-D's). All components on the board were wire-wrapped. The clock circuit is based on a National Semiconductor MM-58167 chip which Keeps time in month, day-of-month, day-of-week, hour, minute, second, and hundredths of a second. The clock is software readable and writeable, and can generate interrupts at either predetermined intervals or at a specific time/date



Figure 2-4. Schematic diagram of the SPICA-IV/LODAS clock circuit.



Figure 2-5. Schematic diagram of one of the three SPICA-IV/LODAS analog-to-digital converter circuits.

combination. For ease of operation the clock chip was interfaced to the SPICA-IV system through a 6522 on the protoboard. Thus the LODAS control software commands the clock through the protoboard VIA rather than controlling it directly. A schematic diagram showing the implementation of the clock circuit is presented in Figure 2-4.

The ADCC board also holds three 12-bit Analog Devices AD-574 analog-to-digital converters, each of these memory mapped into two contiguous bytes. A voltage conversion is initiated by writing to the A-to-D's. A digitized representation of the presented input voltage is obtained by reading the two memory mapped data bytes. The settling time for these A-to-D's is 35 microseconds. All three A-to-D's may be used in either a bipolar or unipolar mode, as selected by switches placed on the front-left edge of the ADCC board. In the unipolar mode the dynamic input range of the A-to-D's is zero to 10 Volts. In the bipolar mode the range is -5 to +5 Volts. Since the DC output of the photometer amplifiers used at RHO produce a zero-to-1 Volt negative going signal, buffer amplifiers (AD-741L's) are used between the A-to-D inputs and the actual signal input to the protoboard. Figure 2-5 shows one of the three A-to-D converter circuits. The circuits for all three channels are identical.

<u>Ancillary equipment</u>. Special I/O signals, such as the analog inputs to the A-to-D's, are connected through the system to a 44-pin connector on the back-left edge of the

protocard. In addition inputs to each of the A-to-D's may be provided through miniature phone plugs mounted on the front-left side of the ADCC board.

The AIM-65, the three MTU boards and the A-to-D/clock board are mounted on an MTU K-1005 Card File and 5-Slot Motherboard (MTU, 1980d). A Zenith Data Systems 12-inch model ZVM-121-Z monitor is used to display video output. A Radio Shack model CTS-41 cassette recorder is used for secondary data storage.

For operational convenience a remote-control paddle was built. The need for remote control can be critical when observing alone, and obviates the need for continually running up and down telescope access ladders to operate the SPICA-IV Keyboard. The paddle holds twelve pushbutton switches, which can select up to twenty-two Key closures on the AIM Keyboard. It is connected to the SPICA-IV/LODAS through its own cable, and specific functions can be activated by the observer from the telescope.

<u>Power supplies</u>. The DC voltages for the system, except those required by the thermal printer and disk drive, are provided by a Power One, model HBB-512 power supply. This supply can source 5 Volts DC at 3 Amperes, and 12 Volts DC at 1.2 Amperes. The thermal printer is powered by a Power One, model HB-24-1.2 (+24 Volt, 1.2 Ampere) supply. The disk drive employs a Power One, model CP-205 power supply providing 24 Volts DC at 1.5 Amperes, +5 Volts DC at 1 Ampere, and -5 Volts DC at 0.5 Ampere. The CP-205 supply

along with the disk drive, is packaged separately from SPICA-IV/LODAS. For use at RHO the disk drive/power supply unit is mounted on the bottom shelf of a rolling equipment cart (see Figure 2-6).

<u>SPICA-IV system configuration</u>. Figure 2-7 indicates the overall system configuration. All major components including peripheral I/O devices are shown. Figure 2-8 is a photograph of the ADCC board. The polarity-mode switches for the A-to-D's are mounted on the top of the board, as is a trim capacitor to adjust the MM-58167 clock rate.

The MTU K-1005 card file holding the digital electronics boards and the HBB-512 power supply, which comprise the major components of the SPICA-IV/LODAS system, are packaged in a small aluminum chassis. All signal and power cables are brought into the system through connectors on the back plate so as not to be obtrusive during operation. A cooling fan, which can be disabled on cold nights, is also mounted on the back plate of the chassis. Figure 2-9 shows the placement and fucntion of each signal and power connector found on the back plate of the chassis. Figure 2-10 shows the assembled SPICA-IV/LODAS system on its rolling cart in operation at RHO.

<u>Portable use</u>. A key element in the system design was the need for relatively low power utilization. The portable aspect of the SPICA system had to be retained in order to use SPICA-IV/LODAS in the field. Observations of lunar grazing or asteroidal occultations often require setting up a

Figure 2-6. The SPICA-IV/LODAS 8-inch floppy disk drive and drive power supply.





Figure 2-7. The SPICA-IV/LODAS system configuration.

Figure 2-8. The Analog-to-Digital Converter/Clock (ADCC) board.



Figure 2-9. SPICA-IV/LODAS backplate and connector layout.



Figure 2-10. The SPICA-IV/LODAS rolling rack and Cassegrain photometric equipment in use at the Rosemary Hill Observatory.



portable photoelectric station in a dark, secluded site where the availability of AC electric power is often non-existent. Thus one reason for using dynamic (as opposed to static) RAM is its lower overall power utilization, drawing higher current only during periods of active write cycles.

For observing at a remote site AC power is required to operate not only the SPICA-IV/LODAS system, but a Kepco model ABC-2500M high voltage power supply, an Astronomical Time Mechanisms model 240V DC electrometer amplifier, and a True Time Instruments WWVB receiver as well. To provide AC power a Nova model 1260-24 DC-to-AC inverter, running on two 12 Volt DC automobile batteries, has sufficient capacity to operate the entire photoelectric station for 35 hours. The AC inverter can supply approximately one Ampere at 120 Volts AC. Thus, to conserve power, the DDDC board and the Shugart 801 disk drive are not used. Data are saved to cassette tape after an observed event. Although the inverter can also provide power for the ZDS 12-inch monitor, this additional load reduces the working life of the portable power supply system considerably. Hence, for field use a Gold Star 12-inch black and white television with an RF modulator that had been built-in is used in its place. Though the power required for the television is no less than that of the ZDS monitor, it can be run directly on 12 Volts DC. The source of 12 Volts can be derived from one of the two batteries supplying the DC-to-AC inverter. In practice, however, the transporting vehicle's 12 Volt car battery is

used to power the television as well as the telescope drive corrector, slewing motors, electric dew cap, and ancillary equipment.

Figure 2-11 shows the SPICA-IV/LODAS system in field use while observing the asteroidal occultation of 14 Piscium by Nemausa on September 11, 1983 (Dunham et al., 1984). In this case AC power was available at the observing site. Figure 2-12, taken on November 13, 1984, shows the SPICA-IV/LODAS system when it was powered by the portable supply system while observing the asteroidal occultation of BD +08 471 by Ceres from the middle of the Florida Everglades.

Analog Electronics

The UWUB receiver. Nather and Evans (1970) have pointed out that the reduction of photoelectric observations, in principle, can yield the times of geometrical occultation with an uncertainty of only one or two milliseconds. As Table 2-4 has shown, a primary design criterion of SPICA-IV/LODAS was to have a data acquisition rate of least one point per millisecond. The inherent degree of accuracy in overall system timing depends upon both the AIM-65 phase-2 clock and the clock on the ADCC board. Thus, in order to realize absolute timing accuracy of one millisecond a standard time calibration source must be employed.

This requirement precludes the idea of using High Freguency (HF) transmissions from the National Bureau of Standards' radio station WWV (located in Fort Collins,

Figure 2-11. The SPICA-IV/LDPAS, and portable photometer on an 8-inch Schmidt-Cassegrain telescope used to observe the occultation of 14 Piscium by 51 Nemausa. The observation was made on the grounds of the Macon (Georgia) County Museum of Aris and Sciences.



Figure 2-12. The SPICA-IV/LODAS, and portable photometric system used to observe the occultation of BD +08 P471 by Ceres from a remote site in the middle of the Florida Everglades. The 14-inch Schmidt-Cassegrain telescope seen in these photographs uses on loan thom the Tampa Amateur Astronomical Society.





Colorado) as a suitable time base reference. The uncertainty in the HF propagation path arising from variability in ionospheric conditions between Fort Collins and Bronson can lead to timing uncertainties as large as a several milliseconds.

Fortunately, NBS provides a Very Low Frequency Coordinated Universal Time broadcast, via radio station WWVB, which transmits at a standard carrier frequency of sixty kiloHertz (Kamas, 1977). At sixty kiloHertz the mode of propagation is strictly by ground waves; hence, the variability in propagation time is removed (Department of the Army, 1953). The propagation path will simply follow a great circle from transmitter to receiver, amounting to a fixed light-time delay of 7.4 milliseconds at the RHO.

WWVB transmits timing information in a tristated time code. The strength of the carrier wave is modulated by reducing output power for 0.2, 0.5, or 0.8 seconds each second. Encoded in this modulation envelope are the time, date, and current UT1 correction. Each ten-second period and the start of each new minute are identified by encoded framing references.

Detection and interpretation of this signal are precisely what is needed to provide the timing accuracy desired. Several avenues of approach were debated. Rather than having a receiver built, a commercially available unit well suited to the task was procured. The unit, a True Time Instruments model 60-T receiver provides not only a detected carrier output, but a TTL compatible code output as well.

A small modification made to the TTL output, dividing it down to 0.8 Volts, permits feeding the code signal directly into one of the three A-to-D converters available in the SPICA-IV/LODAS system. The digitized time code is sampled simultaneously along with the photometric channels. The receiver is mounted on a 19-inch equipment rack, shown in Figure 2-13, along with a WWV receiver and the high voltage power supply used by the PMT's.

In actual use the LODAS system clock is set manually by the observer to an audio WWV signal. This procedure results in a clock setting accuracy of a few tenths of a second. It is then noted if the clock was set fast or slow. Digitized WWVB second transitions then provide a correction to the nearest millisecond.

The signal strength at RHO rarely exceeds 125 microvolts per meter (True Time Instruments, 1974). An active antenna, model A-60FS, also manufactured by True Time instruments is currently used at the observatory. This is marginal under circumstances of unfavorable reception, and an alternate antenna design is being considered for future use at the observatory. However, it has been found that eighty percent of the time a decodable signal is available while observing. During nights of signal fading, time code is digitized before and after the event as conditions permit and time corrections to the computer's internal clock are interpolated in post-observational reduction.

The receiver introduces a measured electronic delay time of 19 milliseconds from the time of reception of the WWVB


Figure 2-13. The WWV and WWVB receivers, and PMT high voltage power supply.



carrier to code output. This, however, varies slightly as a function of signal strength. By attenuating the input signal from the antenna it was found the delay is lengthened to 21 milliseconds at a level where the time code cannot be reliably detected. This then sets the limit of the absolute timing determination to +/-1 millisecond, meeting the occultation program's allowable tolerance. All reductions then have a final correction of 27.4 milliseconds subtracted from the determined time of geometrical occultation, with an additional error of +/-1 millisecond added to the formal error of the solution.

<u>Radio antennas</u>. For use at the observatory the receiver is mounted in the telescope main-power distribution rack, immediately above the WWV receiver. These two receivers share an antenna cable, so only one receiver can be used at a time. The antenna connector must be switched from the WWV receiver to the WWVB receiver before use. Approximately one minute is required by the WWVB receiver after being powered up to lock onto the time code and produce a readable decoded output. A twenty-five foot cable terminated at one end with a BNC connector and a three conductor phono plug on the other is kept on the rolling cart with the SPICA-IV/LODAS system. The BNC end is connected to the CODE output of the WWVB receiver, and the phono plug end connected to one of the SPICA-IV/LODAS signal inputs. The input channel selected to receive the time code should be switched to unipolar mode.

Because of the frequent lightning strikes, unavoidable on one of the highest hills in Florida, the WWV and WWVB antennas are disconnected at the base of the antenna tower at the end of a night's observing. A PL-259 connector can be found at the tower base to which the WWVB active antenna connects via a plug, and the WWV long wire antenna connects by means of an alligator clip. The WWVB antenna and its 15-foot antenna cable are removed from its 6-foot high mounting stand and stored above the desk on the first floor of the observatory building.

The photometer amplifier. Since millisecond time resolution is desired, the photometer amplifier used must have a response at least as fast at a range of gains useful to the occultation observing program. An Astronomical Time Mechanisms model 240 fast photometric DC electrometer amplifier was choosen. This amplifier described by Astronomical Time Mechanisms (1980) is based on a circuit by Uliver (1976) designed specifically with lunar occultation observations in mind. Amplification is achieved in two stages: first in a current-to-voltage conversion stage; and second in a buffer amplifier. A similar amplifier had been in use at the Rosemary Hill Observatory fourty-six centimeter telescope for several years. Caton (1981) carried out a program of UBV photometry on RS CVn stars using this amplifier. He found eighth magnitude stars could be observed giving a full scale reading with the amplifier switched to the so-called "C" gain. This gain setting employs 10 megOhm

feedback resistance in the first amplification stage. Using the seventy-six centimeter telescope would yield a gain of approximately one magnitude over the forty-six centimeter telescope.

The effective time constant of the amplifier is limited by the high precision megohm feedback resistors and the capacitance of the input signal cable (added to a capacitance of 5 pf., the value of the feedback capacitor used on the signal input). In order to observe ninth magnitude stars with a 2 KHz half power response, twice that of the target time resolution, the capacitance of the input signal cable cannot exceed 45 pf. RG-58 A/AU co-axial cable has a capacitance of 28.5 pf./foot. Thus, to achieve this time resolution for stars of ninth magnitude a cable of this type no longer than 18 inches must be used. Rather than RG-58, which is commonly used as a signal cable, the occultation program uses RG-71/U, with a measured capacitance of 13.1 pf./foot.

For practicality, the amplifier is mounted on the side of the photometer cold box as may be seen in Figure 2-9. This allows a short cable run (only eight inches is needed), and is in an extremely convienient place for an observer operating the photometer. Having the amplifier fixed to the photometer also permits the signal cable from the PMT to be securely fastened down thus preventing any possible cable flexure. Such flexure would result in charge redistribution along the signal cable leading to erroneous fluctuations in the observed signal level.

The amplifier coarse gain steps are in increments of 2.5 magnitudes, and fine gain steps in increments of 0.5 magnitudes. After initial use at the telescope the amplifier was modified to provide a 0.25 magnitude gain switch to boost the effective gain at any coarse and fine combination. This was done to provide the observer with a bit more flexibility in chosing the amplification factor used for the purpose of real-time photometric data display.

The output of the amplifier is connected to the SPICA-IV/LODAS system by means of a fifteen foot signal cable terminated on both ends with phono plugs. This cable is kept on the rolling cart along with the WWVB signal cable. One end is connected to the amplifier output marked RECORDER, and the other end is connected to a SPICA-IV/LODAS input switched to unipolar mode.

Limitations of the Occultation Photometric System

Once obtained, the new amplifier was used to assess the limitations of the seventy-six centimeter telescope photometric system and to confirm that stars of reasonable faintness could be observed while preserving a system time constant on the order of a millisecond. Stars over a range of six magnitudes were observed on the moonless night of May 23, 1981 U.T. with the Astromechanics photometer, cooled with dry ice, and a Johnson V filter. Measurements were taken at both the nominal operating voltage of the PMT of 1200 Volts DC, and at the maximum operating voltage of

1600 Volts DC. Five minutes of settling time was allowed after switching voltages before readings were taken. The observations, listed in Table 2-5, give the star name, U.T. of observation, the star's V magnitude, the photometer diaphragm used. For both voltages the amplifier coarse and fine gains, and the signal level due to the star (normalized to a full scale value of 255) are listed. In all cases the dark current and sky background have been subtracted from the star-plus-sky readings. Gain settings were adjusted to give readings as close to 65 percent of full scale as possible.

TABLE 2-5 STELLAR INTENSITY READINGS WITH RHO SPICA-IV/LODAS

Star Name	U.T.	mŲ	Dia.	1200 Gain	Volts Star	1600 Gain	Volts Star
Gamma Leo-a	0315	2.6	J	87	193	82	173
57 UMa	0338	5.2	J	C7	173	87	153
88 Leo-a	0401	6.1	J	C9	165	89	143
88 Leo-b	0413	8.6	J	D10	158	C10	136

The fine gain steps of 0.5 magnitude run from "1" to "11". Hence stars with a V-magnitude as faint as approximately 7 can be observed at a coarse gain setting of "C", at a PMT voltage of 1200 Volts DC. In order to gain one magnitude observations can be made at 25 percent of full scale. With twelve bit digitization this is still roughly one part in one-thousand, or a photometric precision of 0.001 magnitude. Alternatively, a gain can be achieved by increasing the PMT voltage. As can be seen from Table 2-5 increasing the PMT voltage to 1600 Volts provides a gain of approximately 2.4 magnitudes. However, the penalty of

increased thermal noise (dark current) must be paid if this option is taken. Fortunately, the RMS amplitude of the dark current for the 1600 Volt observation of 88 Leo-b was only 2 percent of the star signal level.

These observations tend to lead to over-optimistic results, as occultation observations will not be made in dark skies; indeed the telescope will be pointed in the direction of the moon. To assess a "worse case" condition. the star SAO 098723 (mV=8.7) was observed on May 11. 1981 when it was only 0.1 degrees away from the dark limb of a 55 percent illuminated moon. Using the same diaphragm and filter as the observations in Table 2-5, a gain setting of D7 was required to bring the star-plus-sky level up to 67 percent (170 counts out of 255). The sky contribution was 130 counts, so the star signal was only 16 percent of full scale. In order to bring the signal into the range of "C" gains, an additional gain of 0.5 magnitude is needed. This would place the gain of the amplifier at C11. At 1600 Volts the star-plus-sky signal was just above half scale on a gain setting of C7. Hence under bright sky conditions the photometric system would allow detection of stars down to ninth magnitude with timing accuracy of one millisecond. The limit is more stringent for observations with either of the intermediate bandwidth filters, whose integrated bandpass is roughly one-tenth that of the Johnson V filter. In this case the limiting magnitude is reduced to roughly 6.5.

The Lunar Occultation Data Acquisition System Software Software Design Considerations

The design criteria specified in Table 2-4 were binding not only for the choice of hardware to be used the SPICA-IV system, but applied equally, if not even to a greater extent, to the design of the data acquisition and process control software. The execution of a repetitive task at a precisely defined rate, which must interact in real-time with the "outside-world" is best accomplished by the technique of interrupt processing. Thus, the major process control task of the LODAS software, real-time data acquisition at a rate of at least one-thousand 12-bit samples per second, in at least two channels, was relegated to an interrupt service routine. Yet, some functions of LODAS do not have the need for either repetitive or regularly scheduled execution. For example, scanning the keyboard for user requests, or updating the 20-character alphanumeric display with system status information can be done at the microprocessor's leisure. Hence LODAS actually operates two concurrent programs. A foreground program handling all input to, and output from the observer runs continuously in a relatively quiescent mode, calling upon system services only when required by the observer or program logic. This program is repeatedly interrupted by the aforementioned interrupt service routine. referred to as the background task, on a regularly scheduled basis.

System timing. The limiting factor which played a major role in the development of the LODAS operating concepts was the execution speed of the 6502 microprocessor. The system (i.e. microprocessor) clock rate for the AIM-65, as for all 6502 systems, is 1-megaHertz. A sampling rate of 1-kiloHertz would require the evocation of the interrupt service routine every millisecond. This constrains the process control tasks to a maximum of one-thousand system clock cycles. However, the system cannot spend all of its available clock cycles executing the interrupt service requests. Some percentage of the total system throughput must be allocated to the foreground task. Fortunately, the reaction time of any person is much longer than the cycle time of a 1-MegaHertz computer. This undeniable physiological fact allows a very low bias in time-slicing for the foreground routine. A comfortable allowance of a minimum of 10 percent was deemed more than adequate.

On average, the execution cycle time for a typical machine instruction on a 6502 microprocessor is 3.5 microseconds. This means that approximately three hundred instructions, at most, could be issued in the interrupt service routine before interrupt pile-up would occur. The need for rapid execution speed is clear. For this reason the LODAS software was implemented in 6502 machine language, rather than a more user-oriented, but slower, high level language.

It was found that while the task of data acquisition and storage requires less than a half-millisecond, other interrupt service requests (such as servicing a video "strip chart" display) would demand a total number of machine instructions well in excess of the three hundred maximum. Because of this the interrupt service routine was multi-phased, handling data acquisition and storage in each phase, and pieces of other service requests in successive phases. In breaking the background task into four phases the execution time for any one phase required less than 450-microseconds. Hence, the basic system interrupt rate was defined as 500-microseconds. This allows data to be sampled in successive pairs and averaged together in real-time before being stored. Thus a 1-kiloHertz data sample is actually comprised of two 500-microsecond pair-averaged samples. This not only effectively increases the signal-to-noise ratio of the acquired data by 41 percent, but the stored 1-kiloHertz samples have a Nyquist cut-off frequency of 1-kiloHertz (half the actual data sampling rate) as well.

<u>Memory usage</u>. The area available for storing data in a circulating event buffer is 18-kilobytes in length. In order to optimize the use of this limited (but sufficient) resource, two 12-bit acquired and averaged data samples are bit-packed into three 8-bit bytes. This packing, accomplished by the interrupt service routine, saves 25 percent over storing the 12-bit data, unpacked, into two 8-bit bytes. The 18-kilobytes of available RAM are

partitioned into three 6-kilobyte regions, each to hold one channel's data. Four seconds of pair-averaged data, acquired at an effective rate of 1-kilohertz, can be stored in each 6-kilobyte region. Hence, the system as built and programmed is capable of holding twice the amount of data originally envisioned, and in an additional data channel as well.

A modified version of the LODAS program, called FASTDAS (Fast Asteroidal Data Acquisition System) partitions RAM into only two data storage areas, thereby gaining 50 percent in the data buffer circulation length. This program has been used in observing asteroidal occultations of stars from remote sites (Dunham et. al, 1984).

The LODAS/E07 program has been assembled to reside in ROM at an address space of \$D000. This allows co-residency with AIM-65/FORTH, which is often used on RHO SPICA systems.

<u>Supplementary program documentation</u>. The overall logic flow for the LODAS foreground and background (interrupt service) programs, as well as the program initialization procedure are shown on the operational flow charts presented as Figures 2-13 and 2-14. A fully annotated assembly listing of the LODAS program is contained in Appendix A. This listing reflects LODAS program revision number E07, the seventeenth incarnation of LODAS since its inception. The assembly listing is preceded by a detailed accounting of the LODAS memory space in terms of 1/0 addressing, AIM-65 monitor utilization, program variable space, and an overall system memory map. Following the assembly listing is a symbol-table

Label	Description of Program Step
	(INITIAL PROGRAM ENTRY)
COLDST	[Establish Address of Interrupt Service Routine]
LODAS0	[Print/Display Program Name and Version Number]
SETVIA	[Initialize T1 for 500 Microsecond Interrupts]
CLINIT	[Setup Access Control for Internal Clock]
SETGAP	[Initialize Cold Start Variables]
	[Interactive Setup of All Control Parameters]
	[Initialize Video Display for "Glass Recorder"]
WARMST	[<u>Initialize Action Code to "Quit"</u>]
SETCTR	↓ [<u>Initialize Flags, Counters and Pointers</u>]
SETSTL	↓ [<u>Initialize Data Buffer Addresses</u>]
	↓ [<u>Clear Possible Interrupt Sources: Clock, T1</u>]
SETBT	↓ [<u>Setup Bit Pattern Table for Video Display</u>] ↓
	[<u>Clear and Draw Background Pattern on Video</u>]
MAIN	(Foreground Program: Command Service, Display)
SCAN	[Scan Keyboard for Command Entry]
	<pre></pre> < <u>"Exit to Monitor" Command Received?</u> >NO
	[<u>Clear both VIA Interrupt Flag Registers</u>] ↓
	(STOP! EXIT TO AIM-65 MONITOR)
NEWCDE	[<u>Execute User Request: See Table 2-8</u>]+
DSPCHK	<pre></pre>
DISTME	<main<[<u>Update LED Display]</main<[<u>

Figure 2-14. LODAS foreground program logic flow chart.



Figure 2-15. LODAS background program logic flow chart.

STEP1	[Read A-to-D's, Average with Previous Value]
	[Store Averages in Circulating Data Buffer]
	↓ <stpend<[<u>Update_Data_Buffer_Pointer]</stpend<[<u>
STEP2	<pre>STEP2B<yes<<u>Is Freeze Flag Clear?></yes<<u></pre>
	(Has Delay Expired?)YES
	<pre>STEP2B<[Update Delay Counters]</pre>
SETSNP	<pre><stpnd1<[set clock]+-<="" flag,="" pre="" read="" snap=""></stpnd1<[set></pre>
STEP2B	[Put A-to-D Readings in Average Registers]
	<pre>STPND1 <[Read_Clock]</pre>
STEP3	[<u>Read A-to-D's, Average with Previous Value</u>] ↓
	[<u>Store Averages in Circulating Data Buffer</u>]
	[<u>Update Data Buffer Pointer</u>]
STPEND	[Increment Current Storage-Byte Pointer]
STPND1	[Increment the Interrupt Service Step Count]
STPND2	[<u>Increment the Sample Counter</u>]
RGET	[Restore Y, X and A-Registers from Stack]
	(EXIT INTERRUPT SERVICE ROUTINE)
Statement	s in the flow chart are interpreted as follows:

[<u>Unconditionally Execute Statement In Brackets</u>] (LABEL<--[<u>Execute Statement in Brackets and Go To LABEL</u>] (LABEL<--NO-(<u>If Proposition Matches YES/NO Go To LABEL</u>) (<u>PRIMARY ENTRY OR EXIT POINT</u>)

Follow flow lines either unconditionally, or propositionally as is applicable:

[<u>Statement</u>]------ ⟨<u>Proposition</u>⟩--YES/NO--↓ or ↓

Figure 2-15--Continued.

and reference list. The LODAS makes extensive use of internal subroutines and data tables. Tables 2-6 and 2-7 provide a synopsis of these routines and data tables respectively.

Peripheral Input/Output

User (Observer) I/O. The LODAS program was written with ease of operation in mind. An observer at the telescope often has enough problems confronting him or her, and an unfriendly computer need not be among them. After powering up the SPICA-IV, and evoking LODAS, the observer is prompted on the 20-character alphanumeric display to enter the parameters salient to the observing session. The date, data acquisition rate and channel assignments are among these input requests. These parameters may be changed at any time by issuance of a LODAS command. The LODAS commands are actuated by a single keystroke (or for safety, by depressing the CTRL Key simultaneously with the command key). The command-Key assignments are, in most cases, mnemonic to the command request. A list of LODAS commands is given in Table 2-8. Some of the LODAS commands are acted upon immediately (as in the case of EXIT) and some require additional information from the observer. Examination of Table 2-8 will reveal the proper responses to any LODAS command request prompt. These responses are identical to those required on system initialization when the parameters are first established.

Once the LODAS system initialization is completed the observer is kept abreast of system status on the 20-character

TABLE 2-6 LODAS/E07 SUBROUTINES

Name	Description of Subroutine
CLENUP	Post tape writing clean-up. Reset interrupt count, clear VIA interrupt flags, re-enable 0.1 second interrupts, restore "Action Code".
CLEAR	Clear the 20-character alphanumeric display.
COMENT	Input comment from user, up to 40-characters in length. Display, print, and store in the tape header buffer, (TBUFF).
DECHEX	Convert packed BCD in A-register to hexidecimal.
DISBUF	Output contents of display buffer (DISBUF) to 20-character LED display.
DRSET	Interactive data acquisition rate setting. Get data acquisition rate from keyboard.
FRZSET	Interactive delay time (freeze) setting. Get delay from Keyboard. Validity of entry is checked.
GETKEY	Input a character from the AIM-65 Keyboard, if no key pressed then wait.
GOCLCK	Reset clock with values stored in MILSEC.
HIOUT	Convert left half-word of A-register to ASCII and store in DISBUF at offset indicated by X-register.
HX1 BCD	Convert hexidecimal byte in A-register to 2-byte BCD. Puts MSD into Y-register, and LSD into A-register.
HXASC2	Convert hexidecimal value in A-register to 2-byte ASCII and store in DISBUF at offset indicated in X-register.
HXASC3	Convert two byte hexidecimal value in A, and Y-registers to three byte ASCII and store in DISBUFF at offset indicated by X-register.
KEYCK	Modified AIM-45 Keyboard scan routine. Does not exit to monitor on ESC.

Table 2-6. Continued.

Name	Description of Subroutine
LHWOFA	Halfword shift to A-register to the right. Zero out the left halfword of A-register.
OBHXAS	Convert a one-digit hexidecimal number to ASCII.
PACK2	Get two digit BDC number from keyboard, display number, store as packed BCD in A-register.
PNDM	Print and display an in-line message.
RDCLIN	Set up clock to accept a read request.
RDCLK	Read current time from clock and store in MILSEC. A, X, and Y registers unaffected.
ROLACT	Clear the 20-character alphanumeric display and restore old "Action Code".
TICSET	Set up tics, data channel markers and screen lines on video display.
TIMEGO	Start clock running from keyboard command.
TIMSET	Interactive clock setting routine. Get time and date from keyboard. Start clock if commanded. Validity of entry is checked.
TOGPRT	Toggle AIM-65 printer on/off.
TAPINT	Write occultation observation to cassette tape. interrupt request servicing disabled.
TVCLEA	Clear the video display.
TVDISP	Output next hi and lo resolution data points on video display.
TVSET	Interactive set-up of video display parameters. Get channel assignments and display rate from keyboard.
TVSETX	Clear video display and redraw background.
WRCLIN	Set up clock to accept a write request.
WRTAPE	Transfer a contiguous block of data to tape.

TABLE 2-7 LODAS/E07 DATA TABLES

Name	Description of Data Table
CLCTBL	A table of packed BCD numbers from decimal 00 to 99, inclusive. Used by clock and intensity display routines to 20-character LED display.
TVTICS	A bit mapped data table containing the video display background pattern.
TVTBLH	The hi byte of the address of each video display line as mapped on the visual memory.
TVTBLL	The lo byte of the address of each video display line as mapped on the visual memory.

TABLE 2-8 LODAS/E07 KEYBOARD COMMANDS

Key=5 LODAS Cold Start:

This command, if executed from the AIM-65 monitor, will transfer control to the Lunar Occultation Data Acquisition System and begin execution. The system will respond by flashing LODAS R65/E07 on the alphanumeric display and logging this on the printer. Program default parameters will be established, and the observer asked for variable set up parameters. This command is valid only from the AIM-65 monitor and will be ignored once LODAS is in control and running.

Key=6 LODAS Warm Start:

To re-enter LODAS from the AIM-65 monitor while preserving previously set parameters use this command. It is assumed that LODAS was previously cold started, and exited (i.e. to CODOS or the AIM-65 monitor). Warm start will not require resetting the internal clock/calander, selecting data acquisition rates, display channels or the "video strip chart" display rate.

Key=Cc Enter a COMMENT:

This command will allow the observer to enter a comment of up to 40 characters in length (two lines) on the printed observing log. The most recent comment is also retained in the data buffer header to be saved on disk or tape on command. Entering a RETURN in the comment field will terminate comment entry.

Key=Dc Exit To The DISK Operating System (CODOS):

This command will terminate LODAS and boot the Channel Oriented Disk Operating System. If a CODOS system diskette is not in the disk drive, and the drive door closed the system will freeze up. CODOS must be entered in order to save observing data to a data diskette, or load another observing program from the system diskette. Note: To go from CODOS to the AIM-65 monitor strike the ESC key. To go from the AIM-65 monitor to CODOS, after CODOS had been previously booted strike the F3 key.

Table 2-8. Continued.

Key=Fc Save FILE On Cassette tape

Data may be saved on cassette tape instead of a CODDS disk by issuing this command. Be sure the cassette tape recorder is set to RECORO, and that a non-write protected cassette is in the tape recorder. LODAS will ask for a data FILE name. Any name up to five characters in length may be entered. Data transfer to tape will begin immediately after the entry of the file name. As each block is writen the current block count will be presented on the alphanumeric display. When all data has been transferred the message TAPE WRITE COMPLETED will be displayed.

Key=Gc Restart System U.T. Clock To Preset Time: (GO):

If the system clock had previously been set using the U.T. Clock Set command, but not started, this command will start the clock running at the preset time.

Key=Ic INITIATE Data Taking (Integrate/Sample):

Upon receipt of this command LODAS will immediately begin sampling the three Analog to Digital Converters at the preset sample rate. Acquired data will be stored in the all three channels of the 40% sample circulating data buffers. The Initiate Data Taking command will reset the data buffer pointers to the beginning of the each of the three data storage buffers will overwrite any previously buffered data.

Key=Pc Toggle PRINTER (On/Off):

Each time this command is entered the 20 column printer will toggle from on to off, or off to on. If toggled off command interaction logging will still be displayed on the 20 column alphanumeric display, but will not be printed. The words ON or OFF will be flashed on the 20 character alphaneumeric display to indicate the new status of the printer.

Key=Q QUIT Data Taking:

Upon receipt of this command LODAS will cease data taking after the previously specified delay time. When data taking stops all system status information related to data taking will be saved in the data header buffer for possible subsequent storage to disk or tape.

Key=Rc Select Data Taking RATE:

This command is used to select the data acquisition rate. LODAS will ask for the desired rate, which is to be entered in milliseconds per point. Three digits must be entered. Thus if data is to be taken at five points per millisecond the entry should be 005. Data may be taken at any rate from one to 256 milliseconds per point. An entry of 000 will result in the 256 millisecond per point rate. Data will actually be taken at twice the specified rate and pair averaged before being stored.

Key=Tc Set Delay TIME:

The Delay TIME parameter set by this command affects the system response time to a GUIT command. LODAS will prompt for the desired time delay to be waited before the system will QUIT data taking when commanded to do so. Delay times are entered in units of 100 times the data acquisition rate in milliseconds. Two digits (00-97) must be entered. Thus an entry of 12, with a data acquisition rate of one point per millisecond, will cause LODAS to wait 1.2 seconds before halting data acquisition on a QUIT command.

Table 2-8. Continued.

Key=Uc Set The UNIVERSAL Time Clock/Calendar:

This command will allow the observe to manually reset the internal U.T. Clock. LODAS will first prompt for the the year, month, and day to which the clock should be seeded. Entry must be in the form of a six digit number. For example 850118 will seed the calendar to January 18, 1985. LODAS will then ask for the day of the week, as a single digit number. Day 1 is Sunday and day 7 is Saturday. LODAS will then request the hour and minute to be entered as a four digit number. Thus 1820 will seed the clock to 18 hours and 20 minutes. After these entries are made LODAS will prompt by displaying START=ANY EXIT=CR. If any Key other than RETURN is struck, the clock/calendar will immediately start running with the seeded values. If RETURN is hit, the seeded values will be stored and the clock/calendar can be commanded to start running at a later time by issuing a GO command.

Key=Vc Set Up VIDEO Display Parameters:

This command allows the observer to select which of the three input data channels is be displayed on the Hi Resolution Graphics Display, and which is to be displayed on the Lo Resolution Graphics Display. LODAS first asks for the input channel number to be assigned to the A (Hi Resolution) display, and then for the channel number to be assigned to the B (Lo Resolution) display. Valid input channel numbers are 1, 2, and 3. LODAS then requests the video display rate. The display rate is coupled to the the data acquisition rate. Entry is in units of the data acquisition speed divided by two per displayed point. Thus if the data acquisition rate is 5 milliseconds per point (1/200 second), an entry of 002 will result in a point being displayed every 1/50 second. Entries must be three digit numbers in the range 001 to 255. An entry of 000 is interpreted as 256.

Table 2-8. Continued.

Key=Xc EXIT To The AIM-65 Monitor:

This command will terminate LODAS and return control to the A1M-65 monitor. LODAS may then be restarted with either a COLD or WARM start. This command should be used if CODOS is to be re-entered rather than booted. After exiting to the A1M-65 monitor use the F3 key to re-enter a previously booted CODOS system.

NOTE: A letter designation of "c" postfixing the command key (i.e. $KE^{+}XC$) indicates that the CTRL Key must be depressed simultaneously along with the specified key.

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alphanumeric display. Figure 2-16 shows the format of this display.

HHMMSS M 000 000 128 . 1 1 . Intensity on A-to-D channel #3 . : : : Intensity on A-to-D channel #2 2 : : : Intensity on A-to-D channel #1 : Current (or Last) Command Code 2 Universal Time

Figure 2-16. Format of the LODAS 20-character alphanumeric system status display.

The A-to-D channel intensities are scaled from zero to 255. With no input on an A-to-D channel the display will read 000 if that channel is set to unipolar mode, or 128 if that channel is set to bipolar mode.

All observer commands, command responses, and the current status display are logged on the 20-character thermal printer whenever a command is issued. The printer can be disabled by an observer command.

<u>The video "strip chart recorder"</u>. Without a doubt, strip chart recorders are the bane of photoelectric observers worldwide. Renowned for clogging up, dripping inK, jamming paper or spewing it forth in voluminous quantities, these devices, while undoubtedly useful, often seem more troublesome than they are worth. Since the LODAS saves all observational data in a digital format on disk or tape there is no need for a printed chart record of the observations.

Yet, a chart recorder is an extraordinarily handy tool (when working) on the observing floor, even if used simply to visually monitor the ongoing photoelectric observations. An ideal chart recorder for use with SPICA-IV/LODAS would provide a graphic display of the event for the observer's immediate inspection, without the necessity of plotting it out on reams of paper. A further ideal would be such a device with no mechanical parts to fail, as they invariably do, while observing.

These ideals were transformed to reality with the introduction of what is now referred to as the video "strip chart recorder". Rather than having the photoelectric and/or timing signals drawn on a paper chart with an inking mechanism, data are displayed on a video monitor. The screen is divided into two halves. The left half, called the hi-res, or A-channel display, has a resolution of 0.25 percent (1 part in 256). The right half, called the lo-res, or B-channel display, has a resolution of roughly 1.5 percent (1 part in 64). Data input to any of the three A-to-D's can be selected to be displayed on either the A or B channels. While observing, typically the A-channel is used to display photometric data, and the B-channel is used to display the WWVB time code. The display rate is software selectable, and can be as fast as one-half the data taking rate, or as slow as 1/256 of the data taking rate. With a millisecond data acquisition rate it is possible to display 500 points per second.

The video "strip chart recorder" is best thought of as either a vertical two-channel programmable storage oscilloscope, or a chart recorder with a fixed paper and moving pen. When the pen hits the bottom of the page, the next point will be plotted on the top of the page as the point previously there is erased.

The LODAS activates the video "strip chart recorder" when an Integrate command is issued, and freezes the display when a Quit command times out. A sample video "strip chart recorder" display is seen in Figure 2-17. The number of short dashes at the top left and top right of the display indicate which A-to-D channels have been assigned to the hi-res and lo-res portions of the screen respectively. The short horizontal dashes (on the same line for both hi and lo-res) indicate where the cursor ("chart pen") was when data acquisition stopped. Tic marks on top and bottom indicate the 25, 50, and 75 percent signal levels for the hi-res display, and the 50 percent level for the lo-res display.

<u>Data archival</u>. Post-event observing data may be stored on either 8-inch floppy disks or on cassette tapes. To save data to disk the observer should command LODAS to enter the disk operating system (CODOS), save the observing data by issuing the command: SAVE filename 0700 4FFF, and then re-enter LODAS (if more observations are to be made) with a warm-start. To save data to cassette tape the LODAS Fc command is used. Approximately 10 minutes are required to save the observing data to cassette tape.

Figure 2-17. The SPICA-IV/LODAS and the video "strip chart recorder" display.



CHAPTER III NUMERICAL MODELING OF LUNAR OCCULTATIONS

The Physical Characterization of an Occultation Intensity Curve

When the moon occults a star the lunar limb acts as a diffracting aperture. As the moon's limb approaches the star (as seen topocentrically) the intensity of the starlight is seen to vary as a result of this diffraction. A topocentrically stationary observer sees the diffraction pattern projected onto the Earth's surface moving due to the orbital motion of the moon and rotation of the Earth. The time variation of the intensity of starlight diffracted around the lunar limb under ideal conditions would be characterized by the parameters listed in Table 3-1.

TABLE 3-1 PARAMETERS CHARACTERIZING AN OCCULTATION INTENSITY CURVE

1. The topocentric distance to the limb contact point.

- The angular velocity of the lunar limb, as measured topocentrically, projected onto a line joining the star and the point of contact on the limb.
- The peak wavelength, or spectral energy distribution of the star (a function of the photospheric temperature).
- 4. The unocculted intensity of the starlight.
- 5. The apparent angular diameter of the star.

6. The brightness distribution of the stellar disc.

The actual observed diffraction pattern, referred to as an occultation intensity curve, is affected by several non-idealized distorting effects. These effects are listed in Table 3-2.

TABLE 3-2 PARAMETERS AFFECTING THE OBSERVED INTENSITY CURVE

- The contribution of background skylight due to both atmospherically scattered moonlight, and Earthshine along the lunar limb.
- The local slope of the lunar limb due to surface irregularities.

- Noise in the observation due to both scintillation and photon statistics.
- The spectral response characteristics of the optical system (mirrors, windows, lenses, filter, and PMT).
- Instrumental effects such as PMT dark current and the electrical bandpass of the electrometer amplifier.

The observed occultation intensity curve results from a combination of the effects of the parameters given in both Tables 3-1 and 3-2. The process of "solving" a lunar occultation intensity curve involves isolating or removing the distorting effects and determining the individual parameters intrinsic to the star. The topocentric lunar distance and projected angular velocity (referred to as the R-Rate) can be computed from the moon's orbit. The instrumental parameters can be determined through calibration of the instrumental system. The other intrinsic parameters are determined by a process of fitting an intensity curve, computed from a set of model parameters, to the observed intensity curve. The generation of the computed intensity

curve and the fitting procedure will be addressed in detail in this chapter.

The Generation of a Model Occultation Intensity Curve Monochromatic Point Source Approximation

The process of the formation of a model occultation intensity curve adopted in this study essentially follows the method as outlined by Nather and McCants (1970). As a first approximation, the diffraction of light by a straight-edge from a monochromatic point source is considered. In this approximation the observed intensity of light seen by an observer at a given distance from the geometrical shadow may be expressed as

 $F(U) = (I/2)([X_i + S(U)]^2 + [X_i + C(U)]^2)$ (3-1) where I is the unocculted intensity of the point source. S(U) and C(U) are the familiar Fresnel integrals:

$$S(U) = \int_0^U \sin(\frac{\pi}{2}t^2) dt$$
 and $C(U) = \int_0^U \cos(\frac{\pi}{2}t^2) dt$ [3-2a,b]

The quantity, U, is a dimensionless number referred to as the Fresnel number. In performing diffraction calculations it is convenient to work in Fresnel space, where all linear measurements have been normalized to Fresnel numbers. The quantity U is defined as

 $U = X[2/(LD)]^{\frac{K}{2}}$ [3-3]

where X is the geometrical shadow to observer distance, L is the wavelength of the light from the source, and D is the

straight-edge to observer distance. A typical intensity curve resulting from the defining equations, (3-1, 3-2, and 3-3), as discussed by Borne and Wolf (1970) is shown in Figure 3-1, which also illustrates the geometry of the situation.

It should be noted that at the point of geometrical occultation, when the source, observer, and diffracting straight-edge are co-linear, the diffracted intensity is 0.25 of the intrinsic source intensity. In addition the zeroith order diffraction maximum is approximately 1.37 of the source intensity.

Lunar Limb Effects

While it is true that the lunar limb is not an ideal straight-edge the scale of the occultation phenomenon permits this assumption as a first approximation. The extensive high resolution photographic coverage of the lunar surface obtained by both the Lunar Orbiter probes (Bowker and Hughes, 1971) and subsequent Apollo missions indicate that while the lunar surface is quite irregular it tends to be smoothly undulating rather than possessed of sharp discontinuities. This is presumably due to ceaseless erosional processes caused by the solar wind.

One cannot, however, ignore the possibility of a discontinuity on the lunar surface due to an ill-placed boulder or other protrusional feature, particularly if an occultation is seen at near grazing incidence. The detection and handling of lunar limb irregularites in terms of their



Figure 3-1. Lunar occultation intensity curve for a monochromatic point source.



Figure 3-2. Example of a two dimensional grid used to discretely model a stellar quadrant.

distorting effects on the occultation intensity curve are discussed in detail by Evans (1970). Morbey (1972) has compiled an catalog showing the nature of these effects for a variety of surface discontinuities.

The effect upon occultation intensity curves by lunar surface structure along the line of sight was investigated by Murdin (1971) and found to be negligible.

A Monochromatic Extended Source and Limb Darkening

A star, however, is not a true point source, but has a finite angular diameter. If the star is assumed to be spherically symmetric then the visible disc projected onto the plane of the sky will have a circular cross-section. In most cases spherical symmetry is certainly a valid assumption to make. The vast majority of stars which can yield sensible diameters are of late spectral type and hence are not likely to be rotationally distorted out of spheroidicity (Tassoul, 1978). In the special case of occultations of close binaries the tidal distortions and effects of possibly rapid rotation acting on the component stars would cause the spherical symmetry argument to break down.

The computation of the diffraction intensity for an extended circular source is handled by dividing the disc into a series of discrete elements as shown in Figure 3-2. The two dimensional grid used to partition the model stellar disc is linear in one dimension and radial in the other. The partitioning of the model disc into radial zones allows the effects of limb darkening to be modeled discretely. The

linear strips are small regions containing surface elements (of variable radial intensity due to limb darkening) at evenly-spaced distances from the lunar limb.

The elemental contribution to the intensity of the entire non-limb darkened disc is simply proportional to the area of that particular surface element. Since spherical symmetry is invoked, only one quadrant of the model stellar disc need be computed. The area to a given radius R, bounded by the intercept of the radius with the lines defining the linear strip boundaries is

$$A = \int_{m}^{n} (R^{2} - X^{2})^{-X_{2}} dX = x_{1}[X(R^{2} - X^{2})^{-X_{2}}$$

$$+ R^{2} \sin^{-1}(X/R)] \Big|_{X=m}^{X=n}$$
(3-4)

where m is the X coordinate of the lower bound of the linear strip, and n is the upper bound. Hence, the area of a particular element of the grid bounded by rings of radius R_a and R_b , and strip boundries Sm and Sn (centered approximately at (R.S)) is

$$\begin{aligned} A(R,S) &= \Re[X(R_{a}^{2} - \chi^{2})^{-\chi_{b}} + R_{a}^{2} \sin^{-1}(X/R_{a})] \begin{vmatrix} X = Sn \\ X = Sm \end{vmatrix} \\ &= \Re[X(R_{b}^{2} - \chi^{2})^{-\chi_{b}} + R_{b}^{2} \sin^{-1}(X/R_{b})] \begin{vmatrix} X = Sn \\ X = Sm \end{vmatrix}$$
[3-5]
For convenience, the degree of the partitioning of the model stellar grid is referred to as the grid parameter. The grid parameter is the square root of the number of surface elements per quadrant on the model stellar disc. Thus a grid parameter of 6 would refer to a grid containing 36 elements. Since the grid elements are actually represented in a rectangular coordinate system, those elements which fall outside of the stellar disc have zero intensity. Table 3-3 shows the elemental contribution of each surface element as a percentage of the total intensity of the entire model disc. This is a non-limb darkened star with a grid parameter of 6. In this table element (Zone=6, S#=6) is closest to the center of the stellar disc. The total intensity of this quadrant has been normalized to 25 percent.

TABLE 3-3

BRIGHTNESS DISTRIBUTION FOR MODEL STAR WITH GRID PARMETER=6

S#	Zone 1	Zone 2	Zone 3	Zone 4	Zone 5	Zone 6	Sum
6	0.88918	0.89174	0.89697	0.91095	0.99722	0.69444	5.28050
5	0.00000	0.92110	0.94177	0.98906	1.19189	1.08612	5.12993
4	0.00000	0.00000	0.99725	1.07606	1.37182	1.36939	4.81451
з	0.00000	0.00000	0.00000	1.16089	1.53335	1.60326	4.29750
2	0.00000	0.00000	0.00000	0.00000	1.68034	1.80708	3.48743
1	0.00000	0.00000	0.00000	0.00000	0.00000	1.99012	1.99012

The relative attenuation of the zonal intensity due to limb darkening is modeled by a simple linear limb darkening function of the form

 $I(\theta)/I(0) = (1-\gamma) + \gamma_{cos}(\theta)$ [3-6] where γ is the limb darkening coefficient, and θ is the angle

between the star's radius vector and the line of sight. Table 3-4, which is similar to Table 3-3, shows the relative intensity of the surface elements of a fully limb darkened disc with a grid parameter of 6. The effects of limb darkening coefficients of values 0.0, 0.2, 0.4, 0.6, 0.8 and 1.0 are depicted in Figure 3-3. This figure shows the relative intensity of each linear strip on the model disc from the center of disc to the edge. A grid parameter of 100 was used in this illustration. The maximum variation found in the intensity distributions was for limb darkening coefficients of 0.0 and 1.0, and was only 13 percent.

TABLE 3-4

BRIGHTNESS DISTRIBUTION FOR MODEL STAR WITH GRID PAREMETER=6 AND A LIMB DARKENING COEFFICIENT OF 1.0

	the same and that work pair from the same of	the local data and the local data and			the state of the local data is not a local data in the local data is not a local data in the local data is not a local data in the local data is not a local data in the local data is not a local data in the local data is not a local data in the local data is not a local data in the local data is not a local data in the local data is not a local data in the local data is not a local data in the local data is not a local data in the local data is not a local data in the local data is not a local data in the local	A AMA WHEN HERE AND THE OWNER WHEN AND A	and the second s
S#	Zone 1	Zone 2	Zone 3	Zone 4	Zone 5	Zone 6	Sum
6	0.52257	0.86735	1.07134	1.21774	1.41985	1.01764	6.11648
5	0.00000	0.54132	0.91601	1.18132	1.59329	1.54643	5.77838
4	0.00000	0.00000	0.58608	1.04663	1.63850	1.83057	5.10177
з	0.00000	0.00000	0.00000	0.68225	1.49142	1.91493	4.08859
2	0.00000	0.00000	0.00000	0.00000	0.98753	1.75766	2.74519
1	0.00000	0.00000	0.00000	0.00000	0.00000	1.16958	1.16958
			and the second se	the same second s	the summaries of the su	in some work that white some root was been been	service and while ships ships ships and

The observed diffraction intensity of a monochromatic point source of intensity unity at a fixed distance from the observer is dependent only on the distance of the source from the diffracting aperture. In the limit, each surface element is essentially a point source. Those point sources equidistant from the lunar limb (i.e., along the same strip) will result in an observed diffraction intensity whose total intensity is the linear sum of the individual intensities. Hence, the strip itself can be treated as a point source.



of the stellar disc, strip 1 borders the surface.

Equation 3-1 may be computed for each linear strip, S. The Fresnel numbers, Ug, for the points of evaluation of each strip will be offset by an amount corresponding to the separation of the center of each strip and the occulting aperture. The intensities, I_e, are the fractional normalized intensities for the corresponding strips. Thus, the intensity curve resulting from a grid of individual elements is computed as the linear superposition of strip intensity curves, F_e(U_e). This is shown graphically in Figure 3-4. The uppermost curve is the linear sum of the individual curves shown. The sum curve is shown at a reduced scale so the details of the component curves can be seen on the same figure. In this example the resultant curve is that of a 10-millisecond-of-arc source, at a distance of 375,000 kilometers, and an angular velocity of 0.3 seconds-of-arc per second.

Polychromatic Intensity Curves

The monochromatic approximation to the occultation intensity curve is sufficient to model observations taken with a narrow bandwidth filter (less than 50 Angstrom FWHM). In order to model intensity curves over a broader bandwidth individual model curves, of the monochromatic type described above, must be computed for a number of discrete spectral regions, approximating the contiguous spectral energy distribution of the source. Each individual curve, computed for a specific wavelength, is intensity weighted by the



SCALED INTENSITY

relative total power of the source in that particular spectral region.

In the actual case of stellar modeling, a blackbody spectral energy distribution is assumed. The total power radiated by the star for a spectral region centered at a wavelength, λ_i , and of spectral bandwidth 2B is

$$E_{i}(\lambda_{i}) = \int_{\lambda_{i}=0}^{\lambda_{i}+B} [26729\lambda^{5}(e^{1+473887/\lambdaT}-i)]^{-1}d\lambda$$
 [3-7]

as given by Allen (1963). The quantity T, is the effective color temperature of the star, and could be sought as one of the parameters to be determined by the model fitting procedure. In practice, however, the effective color temperature is entered into the model based on a best guess predicated on the spectral type of the star under study. Numerical experimentation indicates that this procedure is justified. An error of 5 percent in the pre-determined color temperature of the star (quite a reasonable error estimate for stars of spectral type G and later) resulted in only very minor changes to model intensity curves.

Figure 3-5 illustrates an example of the linear superposition of intensity weighted monochromatic curves to generate a polychromatic intensity curve. As in Figure 3-4 the uppermost curve which is the summation of all other curves is shown in a reduced scale. In this example, the 10-millisecond-of-arc curve generated in Figure 3-3 was modeled as a 4500 degree Kelvin blackbody under the bandpass of a Johnson V filter. For computational purposes, only the



fractional contribution of $E_i(\lambda_i)$ is important. Thus, each $E_i(\lambda_i)$ is adjusted subject to the normalization that the sum of $E_i(\lambda_i)$ over all wavelengths considered is set equal to 1. The Effects of Discrete Modeling

The brightness distribution of the strips generated from the stellar grid and the contribution to the total source intensity from each spectral region considered are discrete or quantized representations of continuous functions. Two related questions arise which must be answered before a computational procedure can be developed from the numerical model. First, how fine must the stellar grid be to adequately represent the disc? Second, how small must each spectral region be in modeling the energy distribution of the source?

To answer these questions a number of numerical models were generated and compared. As to the partitioning of the stellar grid, it was found that as the source diameter was decreased, even a moderately coarse grid was sufficient to produce intensity curves indistinguishable from those generated with a fine grid. This may have been expected as the limiting case of a point source requires a grid with only one element. A source with an angular diameter of 20 milliseconds-of-arc required a grid parameter of 10 (400 points, or 20 strips) to achieve a variation of no more than 0.05 percent in the worst case, when compared with a model generated with a grid parameter of 32 (4096 points, or 64 strips).

The effect of choosing too few strips to represent the disc is easily seen in Figure 3-6. This figure shows a family of model occultation intensity curves for a 10-millisecond-of-arc star, similar to the model shown in Figure 3-4 for a variety of grid parameters (as indicated on the vertical axis). In this example the bandpass provided by the intermediate bandwidth "y" filter was used. Choosing grid parameters of 8 or larger caused no discernible difference in the intensity curves.

The model fitting procedure, which will be discussed, generally would use a grid parameter of 8, sufficient for all but the very largest sources. If the fitting procedure indicated a source larger than 5 milliseconds-of-arc, a re-fit was attempted using the roughly determined parameters and a finer grid parameter.

In order to determine the appropriate spectral width to use in discretely modeling an observed optical passband, model intensity curves were generated for sources of 0, 5, 10, and 20 milliseconds-of-arc with a grid parameter of 20. Each of these were "observed" with a Johnson V filter whose bandpass was modeled in spectral regions of 500, 200, 100, 50, and 10 Angstrom widths. In all cases 50 Angstrom steps across the bandpass were required to achieve a worst case precision of 0.05 percent.

The Effects of Instrumental Optical Response

The instrumental response of the optical system affects the measured spectral energy distribution of the source, and



Figure 3-6. The effect of discrete modeling on a 10-millisecond-of arc curve for the grid parameters indicated.

hence must be included in the formation of the model intensity curve. Thus, the relative power in each spectral region from the source, $E_i(\lambda_i)$, computed by Equation 3-7 must be modified by the instrumental response function. The intensity of light measured in a wavelength regime contributes fractionally to the light passing through the integrated bandpass for a given filter. Tables of instrumental response, reflecting the filter function for each of the filters actually used in the observations, were constructed for this purpose. The model un-occulted intensity of the blackbody source, I_i , in a given spectral region, λ_i , corrected for the instrumental response is simply

 $I_{i}(\lambda_{i}) = E_{i}(\lambda_{i})R_{i}(\lambda_{i}) \qquad [3-8]$ where $R_{i}(\lambda_{i})$ is the fractional wavelength-dependent instrumental response. The resulting intensity response function, $I_{i}(\lambda_{i})$, is then normalized such that $\Sigma I_{i}(\lambda_{i})=1$. It is these normalized intensity values which must be used in computing the wavelength dependent intensity curve.

The Polychromatic Extended Source Intensity Curve

The monochromatic point source intensity, F(U), given by Equation 3-1 is a one-dimensional, time-dependent function (since the value of X, the geometrical shadow-to-observer distance is time dependent). A polychromatic intensity curve for an extended source is a three dimensional analog of the one dimensional case, the resultant curve being the linear superposition of intensity weighted curves computed over three variables. The first variable is time, as in the

simple case of F(U). The second is spatially dependent upon the angular separation of the strips in the intensity weighted stellar grid. The third is wavelength dependent upon the spectral energy distribution of the adjusted spectral energy distribution.

Computing a net polychromatic curve for an extended source is quite computationally extensive. For example, consider the generation of an intensity curve over 200 milliseconds, with intensity points 1 millisecond apart (scaled in Fresnel space). If a grid parameter of 8 is employed then a grid of 256 points must first be computed. Intensity weighting this grid for limb darkening and summing along the linear axis after replicating guadrants (by applying spherical symmetry) reduces the grid to 16 points. These 16 points must be evaluated for each of the 200 milliseconds under consideration. If the polychromatic model is to snam the complete response bandpass of a Johnson V filter in 50 Angstrom steps, then an additional 53 values for each of the 3200 computed points must be determined. This amounts to 640.000 values to be calculated prior to weighting and summation in order to determine the net curve.

Models of Double Stars

Before considering the modeling of occultation intensity curves for double stars, a fundamental distinction must be drawn between binary systems which are widely spaced and those which have small angular separations of the components. Because of the millisecond of arc resolution achievable

through the analysis of lunar occultation observations, the more conventionally held concepts of "wide" and "close" binaries are not applicable. In the context of discussing occultation binaries the term "close" does not necessarily refer to contact, or semi-detatched systems. Rather, the reference here is to the apparent angular separation of the components. Occultation binaries which are considered "wide" underoo disappearances which are well separated in time. Specifically, if the intensity curve resulting from the the disappearance of the first star has essentially converged to its post-occultation level before any significant fringing effects are seen due to the disappearance of the second star, then the system is considered "wide". Typically, wide occultation binaries are separated by more than 50 milliseconds of arc. Widely separated double stars (in this sense of the definition) present no additional complications in modeling occultation intensity curves. The disappearance of each star (once detected and noted), can be modeled independently, as their respective intensity curves are not overlapping.

The component stars in close occultation binaries, on the other hand, undergo disappearances closely spaced in time. The observed intensity curve arises from the linear superposition of the diffracted intensity variation for both component stars. Hence the model curve for close double stars is formed by computing the individual curves for the

component stars and linearly superimposing them (along with a constant sky background).

A model intensity curve computed for close occultation binaries (referred to as the two-star case) uses the same filter and systemic response matrix, as well as background sky level for both stars. Individual stellar diameters. times of geometrical occultation, pre-occultation stellar intensities, and velocity parameters are employed. The velocity parameters, in principle, would be the same for both stars, if the lunar limb were smooth. However, since the lunar surface undulates, the contact points of geometrical occultation for the individual stars may have different slopes and therefore different projected angular velocities of lunar limb passage. If the individual spectral types of the component stars are known, then different effective photospheric color temperatures should be employed in computing the stellar spectral energy distributions.

The Differential Corrections (DC) Fitting Procedure A Note on Time

The process of solving a lunar occultation intensity curve requires referencing the observed data points to a zero point in time. Though this time is arbitrary, by convention the time of geometrical occultation has been adopted. In terms of Fresnel numbers, times after geometrical occultation for a disappearance (i.e. when the observer is within the geometrical shadow) have negative values. Unfortunately, the

time of geometrical occultation is unknown, and indeed is one of the parameters for which a solution is sought. Thus, throughout the fitting process times are referenced to the time of the first point in the observed data set. Units of time are referred to in milliseconds as this is most commensurable with the time-scale of the occultation phenomenon. All observations to which the fitting procedure has been applied have been made with 1-millisecond sampling. Thus, for convenience, the term "bin number" is often used interchangeably with "time with respect to data sample number zero". After the time of geometrical occultation has been found by the fitting procedure it is referenced to Coordinated Universal Time.

Choosing Initial Parameters for the DC Procedure

The fitting procedure is basically an iterative, non-linear least squares, differential corrections process following the method outlined by Nather and McCants (1970). Initially, a model curve is generated from physical quantities and model parameters listed in Table 3-5.

The choice of initial parameters in some cases is quite obvious. The number of points to be generated in the model must be the same as the length (i.e. number of data points) extracted from the 40% milliseconds of observational data which is to be fit to a model curve. The selection of this extracted data set must be done by visual inspection, and roughly centering the subset on the apparent time of occultation. The specification of the filter response table

TABLE 3-5

PARAMETERS FOR GENERATING THE INITIAL MODEL

Model Generating and Controlling Parameters

- The number of points to be generated in the model occultation intensity curve.
- 2. Specification of the filter/systemic response table.
- Maximum number of iterations to executed in the corrections procedure.
- Fraction of computed adjustments to be applied at each iterative adjustment step.

Fixed Physical Parameters

- 1. The topocentric distance to the lunar limb.
- The predicted R-rate, based on a smooth limb (seconds of arc per second).

Normally Fixed Physical Parameters

- The effective color temperature of the stellar photosphere.
- 2. The stellar limb darkening coefficient.

Adjustable Parameters

- 1. The stellar angular diameter.
- The pre-occultation intensity of the star plus the background skylight.
- 3. The velocity of lunar limb passage.
- 4. The time of geometrical occultation.
- The post-occultation intensity of the background skylight.

must obviously match the instrumental set-up used in obtaining the observational data.

The next two parameters are concerned not with the initial model generation, but with control of the adjustment procedure. The number of iterations required for a convergent solution depends strongly on the signal-to-noise ratio of the observations. Experimentation with the numerical and computational processes suggests that a maximum of 10 to 20 iterations is usually sufficient. The fraction of computed adjustments will be discussed in the section on the application of partial parametric adjustments.

The fixed parameters have already been addressed. In principle, there is no reason why the two parameters listed as "normally fixed" could not be treated as adjustable parameters as well. However, the model is rather insensitive to the effects of limb darkening, and as previously indicated stellar temperature (within the normal accuracy determined from the spectral type, and if available luminosity class). Unless there is an astrophysically compelling reason, a limb darkening coefficient of 0.5 is suitable in virtually all cases.

The stellar angular diameter, unless previously determined by another occultation, or in rare cases by interferometric methods, is unKnown. As already mentioned an initial guess can be made on the basis of the star's parallax and spectral type, or by application of the Barnes-Evans relation. However, an initial guess of a

"middle-of-the-road" value such as 5 milliseconds-of-arc will be adjusted rather quickly even if grossly wrong.

The initial quess for the velocity of lunar limb passage should be the R-Rate, as any variation from the predicted rate would be due to a local slope of the lunar limb. This slope, in fact is derived from the difference between the predicted and observed velocities. The numerical procedure actually works with a linear, rather than an angular rate. Thus the lunar limb distance is used in conjunction with the R-Rate to compute the L-Rate, the linear rate of the motion of the geometrical shadow across the telescope in units of meters per millisecond.

Initial values for the pre and post-occultation intensities can be determined by averaging the intensity of a few hundred milliseconds of data before and after the time of geometrical occultation. The time of geometrical occultation can be estimated from visual inspection of the observed intensity curve.

Parametric Adjustment

Once an initial model curve is computed, each of the adjustable parameters in turn is varied by a small percentage in order to numerically compute the partial derivatives of the intensity curve with respect to each of the parameters. Analytic evaluation of the partial derivatives is obviously impossible as the intensity curve itself is determined from discrete numerical functions. Due to different degrees of sensitivity of the intensity curve to numerical variation of

the parameters, and indeed to the values of the parameters themselves, experimentation was required to ascertain an optimal percentage variation to use. This would present no problem if an analytic form for the computation of the derivatives existed. Choosing a variation which is too small results in computational truncation errors and discontinuities in the partial derivatives. Too large a variation results in a loss of computation precision. It was found that for most ranges in the values of the parameters that a numerical variation (multiplicative factor) of 1.001 was applicable for the intensities and time of geometrical occultation. The partial derivative of the intensity curve proved better behaved with a variation of 1.005 for the velocity parameter, and 1.02 for the stellar diameter.

Since the diameter can go to zero for a point source, numerical limits must be put on the variation procedure to assure no computational singularites arise. This is also true, conceptually for the velocity parameter as well (and in the unrealistic case of no background sky contribution the post-occultation intensity). However, this would be important only for true grazing incidence which is virtually never seen. Hence, as the angular diameter or the projected velocity (or both) approach zero the variation of these parameters in computing the numerical partial derivatives should be increased accordingly.

Once the partial derivaties have been computed a matrix of residual equations of the form

$$O_j - C_j = \sum_{i=1}^{n} \frac{SC_j}{sP_i} dP_i$$

$$[3-9]$$

can be established. The P_j 's are the varied parameters, n is the number of parameters in the model (normally 5), the C_j 's are the computed intensity values, and the O_j 's are the observed values. The ΔP_j 's are the adjustments to be applied to each of the parameters. These adjustments are determined by solving the residual equations by the method of least squares and is discussed in a general formulation by Brown (1955).

Eichhorn and Clary (1974) have suggested that when the adjustment parameters derived from non-linear equations of condition are not significantly larger than the adjustment residuals, then the second order terms should be included in the adjustment residuals. However, as can be seen, Equation 3-9 has been linearized, and hence does not include higher order terms to represent the non-linear equations. This was done primarily for a practical reason. The numerical evaluation of the first derivatives of the non-linear, non-analytic equations alone is already computationally extensive. The actual computer time required to arrive at a solution is discussed in the section on computational procedures. The additional computations required for the inclusion of the higher order terms in the solution were computationally prohibitive with the computer resources available.

As a result of this linearization the derived adjustments are not completely correct. Thus, the adjustments must be applied to the parameters, and the process repeated until convergence is achieved. Fortunately, the least squares process is sufficiently robust that convergent solutions (with less than 1 percent change in successive adjustment attempts) were usually achieved for all parameters in less than a dozen iterations of the adjustment procedure.

The DC Fitting Procedure for Close Double Stars

A procedure analogous to the single star DC fitting process previously discussed is followed in the case of a close double star. In this case, refered to as the DC2 procedure, if the limb darkening coefficients of the stars, and photospheric temperatures are held constant, then the model is parameterized by nine adjustable quantities. These are listed in Table 3-5 as "Adjustable Parameters" numbers 1 through 4 (for each star) and the background skylight intensity (item 5). Hence, Equation 3-9 is the same as in the single star case, with n=9 instead of n=5.

The choice of initial parameters for the two-star fitting procedure is similar to that of the single star case, though "good guesses" for the starting values of some parameters are more difficult to select than for the fitting of a single star model. Since the observed curve is actually

the linear combination of the two single star curves, the approximate times of geometrical occultation and the individual pre-occultation stellar intensities are not easily determined by visual inspection. Though the DC procedure is quite forgiving for even wildly disparate initial parameters, the required computation time for convergence can be reduced by giving some thought to the starting parameter selection.

To serve as a guide to selecting the initial time and intensity parameters, a family of model curves was generated for a series of monochromatic double stars considered as point sources. The curves covered the range of stellar separation in Fresnel space from U = -2.5 to U = +2.5 in steps of 0.2, and a range of magnitude differences of 0.0 to 2.5 in steps 0.5 magnitudes. By comparing these against the observed intensity curve, a reasonable match could usually be found. The parameters used to generate the selected model curve were then applied to the particular two-star DC solution under consideration.

The task of generating the simple monochromatic two-star models from point sources was relegated to a desk-top microcomputer (Commodore SP9000). The computation of the Fresnel intensity curves (Equation 3-1), by the numerical integration of the Fresnel integrals (Equations 3-2a and b), and the production of the 306 graphs of the two-star intensity curves required 25 hours. However, this only had to be done once, and a compendium for future guidance in parameter selection was created. Figure 3-7 is a sample of a





two-star curve extracted from this compendium. In this case the star which disappeared first is the fainter star (by 0.5 magnitudes), and leads the brighter star by -0.96 Fresnel units. The individual curves, as well as the composite curve, are shown.

Validation of the Fitting Procedures

To validate the fitting procedures (and check on the function of the computer programs, which are discussed in Chapter 4), a number of model intensity curves were generated, treated as observations, and subjected to the fitting process. The model curves which served as synthetic observations were generated with a very fine grid of 4096 points. All the synthetic observations were created assuming a lunar distance of 375,000 kilometers, a stellar temperature of 4500 degrees Kelvin, and an R-Rate of 0.3 seconds of arc per second. Both monochromatic (5500 Angstoms) and polychromatic model curves were synthesized. The latter were modeled under the bandpass of a Johnson-V filter, with discrete spectral regions only 20 Angstroms wide.

For single stars, diameters of 0, 8 and 20 milliseconds of arc were considered. Double stars models used all combinations of 0, 8 and 20 milliseconds of arc with temporal separations of 2 and 20 milliseconds. The synthetic curves were subjected to the DC and DC2 procedures, respectively. The fitting programs employed a grid parameter of 8, and in the polychromatic cases each spectral region within the V filter bandpass was 50 Angstroms wide. The initial parametric guesses for the adjustable parameters were deliberately chosen to be far from the correct values. In every case the DC fitting process converged rapidly (in three to five iterations) to the correct values in all the adjustable parameters. The same, unfortunately, could not be said for the nine parameter DC2 fitting procedure. The DC2 process would usually recover the synthetic model parameters, but only after many (often greater than thirty) iterations. Worse, however, was that on several occasions the adjustment process would either fail to converge to a solution, or fall into a local minimum in parameter space, unreflective of the true parameter set. This was a problem that had to be addressed.

While these tests showed that the numerical (and computational) processes were indeed correct, they were overly optimistic as indicators of the the ease of recovering the parameters of the synthetic curves, both in terms of the formal errors attached to the solutions, and the small number of iterations required (for the DC procedure). This overoptimism resulted from the fact that observational noise (due to both seeing effects and scintillation) were not included in the initial synthetic models.

The effects on the DC and DC2 procedures from the introduction of noise to the synthetic observational curves were then considered. Synthetic noise could have been generated stochastically. However, this would have required making assumptions about both the distribution function (in

terms of relative amplitudes) of the noise, and the temporal variation in the stochastic noise on time-scales shorter than the length of the synthetic data set. Rather than make such assumptions, noise which was actually observed, as extracted from a pre-occultation observational record, was applied to the synthetic curves.

The character of the observational noise would be different for each observed occultation, and in fact is a function of the local seeing and scintillation effects and photometric signal-to-noise ratio (as will be discussed in Chapter 5). Thus, ideally, it would have been preferable to consider all of the above synthetic cases under a variety of observed noise conditions. However, due to the limited system throughput on the computer employed in this investigation, a compromise had to be made.

The noise characteristics of the occultation observation of ZC0835 were found to be nicely representative in terms of the noise figures and power spectra seen in most of the observational records. The distribution function of the observational noise obtained from this observation was applied to the synthetic models, and the noisy models then re-reduced by DC and DC2, with the same initial starting parameters used in the first trials.

Once again, the model parameters were essentially recovered: however, three effects were noticed. First, and somewhat expected, the formal errors of the solution were significantly worse than in the "clean" synthetic cases.

Second, the number of iterations required for the DC model to converge was increased by roughly a factor of four. And third, for the corresponding trials where DC2 solutions were previously obtained, the DC2 procedure would not asymptotically converge to recovering the solution parameters. Rather, after initial movement toward convergence, oscillation about some set of final values would set in. The amplitudes of these oscillations were often as great as 100 percent of the mean values about which the parameters were oscillating. This effect was also seen, but to a much smaller degree, in the DC solutions.

While this behavior (except for the recalcitrant DC2 solutions) was not unacceptable it lead to a number of numerical experiments in an attempt to improve the basic DC and DC2 fitting procedures.

Numerical Experiments to Improve the Fitting Procedure

Two types of improvements were sought to enhance both the utility of the fitting procedures and the statistical significance of the numerical results. First, from a practical standpoint, it was desired to reduce the amount of computation time required to achieve a convergent solution. Second, the DC procedure, (and usually DC2 as well), could successfully recover the physical parameters intrinsic to the intensity curve in question; it was of interest to see if the formal errors of the solution parameters could be reduced. Addressing the first point, it was noticed that the fitting process would approach convergent solutions rather rapidly in most cases. But, on occasion, rather than smoothly approaching the true parametric values (i.e., those used to generate the synthetic curves), the numerical values of the adjustable parameters would oscillate about some undetermined, final values. This effect was usualy small for the five parameter DC fitting process but often was quite pronounced for the nine parameter DC2 routine.

Since the residual equations had been linearized, as previously discussed, the adjustments applied at the end of each iteration step tended to overcompensate for the fact that the higher order terms had been omitted. Thus, as the solution was approached, the adjustments to the parameters computed from the residual equations would overcorrect the interim values of the adjustable parameters, and thereby lead to the noted oscillations.

The uppermost set of graphs in Figure 3-8 shows the progress of the DC fitting process as the adjustable parameters approached their final values through twenty iterations of parametric adjustments. From left to right the parameters presented are the angular diameter, pre-occultation intensity, post-occultation intensity, time of geometrical occultation and L-Rate. The correct values (those used in generating the synthetic curve) are indicated by the dashed lines. This example is from the DC fitting of an 8 millisecond of arc star, with observational noise added to the synthetic curve as specified in the previous section.

"noisy" data for twenty iterations. The set of graphs labeled "A" correspond to the solution obtained using full parametric adjustments. The "B" and "C" graphs which were Fourier smoothed. The dashed lines are the true intrinsic parametric Figure 3-8. An example of the convergence of the DC fitting process applied to values. The thin vertical lines represent a 1 percent variation from the true were obtained with a partial adjustment factor of 0.5. The "B" solutions were from a fit to the raw data, while the "C" solutions were from a fit to data values, except in the case of the angular diameter where these lines show 50 percent variations. The thick vertical lines represent the one sigma uncertainties in the solution parameters.



In this case, the noise was taken from the residual amplitude distribution function for the occultation observation of ZC1221 (see Figure 5-19). The mean value of this distribution function is skewed negatively by 1.706 percent. (The physical reasons for this skewing, and the Poisson character of some of the observed residual distributions are addressed in Chapter 5.)

The negative bias (or offset) in the noise was not corrected before application to the synthetic curve. Hence, there was an expectation that the pre-occultation and post-occultation intensities would be underestimated, and the L-Rate overestimated by this amount. This expectation was indeed borne out, within the one sigma uncertainties of the final parametric solution values.

The model generating parameters, and parameters for the best solution determined by the DC procedure, are presented in Table 3-6. The table indicates that the best solution, as determined by the smallest value of the sum of the squares of the residuals, was achieved on the fourteenth iteration. The parenthetical values are the one sigma uncertainties in the formal errors of the solutions. It should be noted that while the best solution required fourteen iterations, after only four iterations the interim values for all the adjustable parameters were not significantly different from the later determined best parametric solutions. Application of Partial Parametric Adjustments

To reduce the degree of oscillation about the finally determined values of the solution parameters, the DC process,

starting with the same initial values, was run again with one minor change. Rather than applying the full corrections derived from the linearized residual equations, only partial corrections were made at each iteration step. The initial tests of the partial corrections procedure applied a fractional correction factor, denoted §, of 0.5 (i.e. only half of the computed adjustments were applied to each of the interim values of the parameters at the end of each iteration step. If the adjustments at the end of each step had been absolutely correct (which they were not) then the solutions would be approached asymptotically. In that case, after n-iteration steps the current value of a given parameter, P_i , would be

 $P_{i,n} = P_{i,s} + [(1 - \xi^n) \cdot (P_{i,0} - P_{i,s})]$ [3-10] where $P_{i,0}$ is the initial guess for the ith parameter, and $P_{i,s}$ is the true value of the ith parameter.

Thus, after ten iterations the interim solution parameters would converge to better than 0.1 percent of their true values. Of course, § need not be given a value of 0.5. If a larger value is used the true parameter values are approached more quickly, but larger amplitude oscillations set in after initial convergence has begun.

The middle set of graphs in Figure 3-8 shows a typical historical run of the solution parameters with a partial adjustment of $\xi = 0.5$. As may have been expected, the best solution was obtained on the last iteration. By that time, however, the variations in the parameters from one iteration

to the next were insignificant. Table 3-6 shows that the parametric solutions obtained for the DC fitting with partial adjustments (ξ =0.5) were essentially identical to those obtained using the full parametric adjustments (ξ =1.0). This was found to be the case for the fitting of all other synthetic curves as well.

As a result of this numerical experiment a new procedure was adopted in the differential corrections process. Full parametric adjustments were made at each correction step until oscillation of the parameters began to show up. It was found that an increase in the sum of the squares of the residuals was a good indicator of the onset of oscillation. After that, only fractional adjustments (\$(1.0) were applied. Parametric Grouping into Computational Subsets

The application of partial parametric corrections to the DC2 procedure was quite successful in reducing the degree of oscillation of the solutions and obtaining convergent solutions for those cases where none were previously achieved. The number of iterations required to obtain acceptable convergence in the DC2 fitting process was greater than in the case of single star DC fitting, but not unacceptably so. This was expected, as the number of degrees of freedom had been increased by the inclusion of an additional four adjustable parameters. In principle, as long as the intrinsic parameters were recoverable, no further modifications to the DC2 process were required.

TABLE 3-6

COMPARATIVE	SAMPLE	0F	DC	FITTING	то	SYNTHETIC	CURVE
-------------	--------	----	----	---------	----	-----------	-------

PARAMETERS FOR SYNTHETIC INTENSITY CURVE							
Distance (km # of Points Temperature: Limb Darkeni Filter: Grid Paramet	Diame Time Pre-E Post- L-Rat	Diameter (ms. of arc): 8.0 Time (milliseconds): 101.0 Pre-Event Intensity: 1495.4 Post-Event Intensity: 498.5 L-Rate (meters/s.): 545.4					
		ETHOD OF	DC SOLUT	I ON			
	-						
Smoothing §	Raw 1.0		Raw 0.5		150 Her 0.5	tz	
Iterations	14		20			20	
Diameter Time (msec) Pre-Event Post-Event Velocity	8.10 101.2 1471.3 489.4 572.8	(1.28) (0.65) (10.3) (10.4) (18.6)	8.12 101.2 1471.2 489.2 571.4	(1.28) (0.65) (10.3) (10.4) (18.8)	8.10 101.2 1471.2 489.2 572.6	(0.75) (0.38) (6.3) (6.1) (10.7)	

It was of interest, however, to evaluate a computational procedure discussed by Wilson (1976), and to determine its applicability to fitting lunar occultation intensity curves by non-linear least squares, differential corrections. In determining his solution to the light curve of the eclipsing binary TX UMa, Wilson encountered similar difficulties. The approach he suggested, of dividing the adjustable parameters into two subsets which were adjusted independently, was tried. When coupled with the procedure of using only partial parametric adjustment, this parametric subgrouping method was found to yield convergent solutions, but not as rapidly as by adjusting all nine parameters simultaneously.

The variations in the methods of DC2 fitting were compared by attempting solutions to a variety of synthetic curves. Table 3-7 is an example of the comparative results obtained after twenty iterations of an illustrative case studied. The synthetic parameters for generating the two-star model curve are given in the table. This was a case for which full parametric adjustment (5=1.0) in the DC2 fitting process could yield no convergent solution. In both cases shown where the Wilson method was tried (for ξ=0.5 and \$=1.0) the process was still heading toward convergence where the DC2 process, with \$=0.5, had already obtained convergent solutions. As can be seen, after 20 iterations. the formal errors of the solution parameters are much better in the case of partial adjustments with the DC2 procedure. Figure 3-9 (similar to Figure 3-8) shows the history of the

TABLE 3-7 COMPARATIVE SAMPLE OF DC2 FITTING TO SYNTHETIC CURVE

PARAMETERS FOR SYNTHETIC INTENSITY CURVE								
Distance (km): 375000	Diameter Star	- 1	5.00				
# of Points	(msec): 201	Diameter Star	2:	10.00				
Temponature	Stap 1+ 4500	Intensity Sta	ar 1:	1000.00				
Temperature	Stan 2: 5000	Intensity Sta	ar 2:	800.00				
Limb Dankeni	no (1 & 2) * 0.	5 Time Star 1:		90.00				
Coid Papamot		Time Star 2:		101.00				
Stropping C	40	I-Rate Star 1	1 :	400.00				
Macabasa	- 5500 Aportoo	me L-Rate Star		545.41				
monochromati	c, JJ00 Higstro		sd.	400.00				
		SK7 Backgroui		400100				
		CURCERCE (CTARTIN						
INI 1	AL PARAMETRIC -	GUESSES (SIMKIING	5 VALUES/					
		2 Diseaster C	ton 24	0 204				
Diameter St	ar 1: 8.20	Z Diameter S	Car Zi	500				
Intensity S	tar 1: 1200	Intensity Star 2: 500						
Time Star 1	: 95	lime Star A	21	103				
L-Rate Star	1: 500	L-Rate Star	r 2:	500				
BacKground:	600							
	50	12011005						
M - 4	Course of	Conversion	Unon	nunad				
method	orouped	Grouped	01011	Dahen				
5	1.0	0.5	0.5					
a:	(40 (0 14)	(10 (1 00)	E 00	(0.00)				
Diameter I	6.40 (0.14)	6.48 (1.08)	10.00	(0.00)				
Diameter 2	8.98 (0.18)	0.8 / (0.21)	10.00	(0.00)				
Intensity 1	1166.5 (11.1)	11/8.8 (141./)	1000.00	(0.00)				
ntensity 2	631.6 (7.2)	629.1 (7.0)	800.00	(0.00)				
Time 1	91. (0.4)	91.3 (1.5)	90.00	(0.00)				
Time 2	103.7 (0.1)	103.9 (0.2)	101.00	(0.00)				
L-Rate 1	392.2 (2.7)	391.7 (10.1)	400.00	(0.00)				
L-Rate 2	494.0 (2.0)	489.4 (2.4)	545.41	(0.00)				
Background	394.5 (8.8)	394.1 (2.4)	400.00	(0.00)				
Comment:	Rejected	Rejected	Accepted					

Note: No convergent solution for ungrouped, ξ=1 case.
Figure 3-9. An example of the convergence of the DC2 fitting process applied to synthetic data for twenty iterations. The set of graphs labeled "A" correspond to the solutions obtained using full parametric adjustments. The "B" and "C" graphs were obtained with a partial adjustment factor of 0.5. The "B" solutions were from a fit via the parameter grouping method, while the "C" solutions were obtained by adjusting all parameters simultaneously. The dashed lines are the true intrinsic parametric values.





interim solutions for these three cases. Similar results were seen in all the synthetic trials considered.

Uniqueness of the Solution

The iterative differential correction process cannot absolutely guarantee, in a rigorous sense, that the set of solution parameters obtained is unique. To verify that the set of solution parameters is correct one must first confirm that they are physically meaningful, within the context of the observation. Clearly, solutions with astrophysically improbable implications (i.e. stars with diameters of hundreds of milliseconds of arc, or negative lunar L-rates), must be viewed with obvious suspicion and rejected. Fortunately, such occurences have been found to be very rare.

Meyer (1975) states that while a mathematically rigorous statement of the uniqueness of the solution is not possible, reasonable certainty in the validity of a set of solution parameters can still be obtained. If the parameters are indeed physically meaningful, the non-linear differential corrections procedure should be re-run several times with initial guesses of the adjustable parameters of widely differing values in parameter space. If the adjustable parameters converge to the same solution, then one can be quite confident that the solution has not been misguided by being "stuck" in a local minimum in parameter space.

This seems to exact the penalty of having to repeat the computational process several times. This however, is not necessarily true. Quite often, after only one or two

iterations of subsequent runs of the DC or DC2 procedures. with widely differing starting values, it is apparent that the parameters are being adjusted toward the same final values.

The basic DC method was adopted for the initial iterative adjustments in the case of single stars. The application of partial parametric corrections were employed in the two-star DC2 fitting process and in the the final iteration steps for the DC single star fitting procedure. The validity of the solutions was checked by multiple runs of the initial corrections steps with very different starting values.

Smoothing of the Observational Data

As previously mentioned, the second area where improvements to the basic fitting procedure seemed possible was in obtaining more realistic error estimates of the solution parameters. The formal errors of the solution parameters were derived from the variance-covariance matrix of the final set of residual equations. These errors reflect the magnitude of the squares of the residuals. If the residual amplitudes were smaller, then the error estimates would have been tighter as well.

Observationally, this problem is addressed by trying to improve the signal-to-noise ratio. This can be done by rejection of erroneous background light through a good optical baffle (see Chapter 1) and selection of a filter well suited to the spectral energy distribution of the star under study. Once data are acquired, the characteristics of the noise in the raw observational record are fixed: however, this does not preclude the possibility of the raw data being preprocessed before being submitted to the fitting procedure. The question which naturally arises is: can the data be preprocessed in such a manner as to effectively improve the signal-to-noise ratio, without degrading the underlying occultation record itself?

The time scale of variation of an occultation intensity curve is typically 10 to 50 milliseconds. This suggests that some type of time dependent smoothing might be applied to the data, which were acquired at a pair-averaged rate of one sample per millisecond. To see if this inference was indeed true the previously generated synthetic intensity curves, with noise added as before, were subjected to several different smoothing algorithms. The smoothed, synthetic observations were then re-fit by DC and DC2, as appropriate to the model in guestion.

<u>N-point unweighted smoothing</u>. Each of the synthetic curves previously considered was first smoothed by a simple N-point unweighted moving average. Values of N=3 and N=8 were considered, corresponding to sampling frequencies of 330 and 125 Hertz respectively. This method of smoothing has the advantage of reducing the effect of spurious single-sample values within the smoothing window, at the expense of losing information associated with level transitions which are fast compared to the smoothing length.

The application of N-point smoothing to the synthetic intensity curves resulted in solutions determined by the DC and DC2 procedures that were significantly worse in some of the parameters than those obtained in the unsmoothed trials. In every case the angular diameters of the stars were overestimated. In the case of 8-point smoothing, the determined angular diameters were typically too large by factors of 2 to 5. The times of geometrical occultations were shifted, by a smaller degree, to times earlier than used in generating the synthetic models. The L-rate was consistently low, but only to a small degree; however, the formal error attached to the determined L-rate in some cases was even worse than in the unsmoothed trials. Only the recovery of the true values of the pre-occultation and post-occultation signal levels were unhampered by N-point smoothing. Since the determination of the primary parameters of astrophysical interest (the stellar angular diameter, and time of geometrical occultation) were impeded by N-point unweighted smoothing, this approach was deemed untenable.

<u>N-point weighted, exponential smoothing</u>. At this point, the idea of applying a moving average to the observed (or synthetic) data was not yet abandoned. A weighted moving average was used to smooth the set of trial synthetic (noisy) curves. The weighting function chosen was an exponential in order to provide a numerical analog to observationally increasing the time constant of the instrumental system. The electrometer amplifier employed (see Chapter 1) uses a

passive filter network to set the 1/e folding time of the amplifier response. Hence, investigating the use of an exponentially decaying weighting function would mimic the effect of observing with various instrumental time constants.

To be commensurable with the trials attempted in the investigation of unweighted smoothing, decay constants were chosen to provide an effective FWHM in the smoothing function of 330 and 125 Hertz. The results of the DC and DC2 fitting were similar to, though not quite as bad, as fitting to the data which were smoothed with the unweighted functions. Here, too, the effect of increasing the smoothing width resulted in an overestimation of the stellar angular diameter.

Both N-point smoothing techniques proved unusable as methods to preprocess the data in an attempt to reduce the residual amplitudes. These smoothing methods, however, were useful in the visual inspection of the observed intensity curves. It is difficult for the eye to see structure in noisy data, and N-point smoothing helped bring out details which would otherwise would have been visually obscured. Figure 3-10 is a 200 millisecond extract from the observing record of the occultation of ZC1222, centered roughly on the time of geometrical occultation. Eight-point (125 Hertz) unweighted smoothing was applied to these data, and the smoothed curve is superimposed on the raw data.

<u>Smoothing by forward and inverse Fourier transformation</u>. Since the background noise was superimposed on the synthetic (and observational) data, the simple N-point smoothing





algorithms acted with equal effect in not only smoothing out the noise, but the underlying occultation intensity curve as well. What was needed was a method of smoothing only the background noise, while leaving the character of the intensity curve itself effectively undistorted.

Examination of the raw occultation data records seemed to indicate that the temporal characteristics of the variation in background noise were, generally, quite a bit faster than the 10 to 50 millisecond time-scale of variation typical of the intensity curves. This implied that a numerical analog to active low-pass filtering of the data might improve the effective signal-to-noise ratio and subsequently the statistical certainty of the recovered solution parameters.

To see if this implication was true, the power spectra of synthetic "clean" intensity curves had to be compared with the power spectra of actual observed background noise. Several additional synthetic curves were generated and Fourier transformed into the frequency domain. In addition to the synthetic curves already available, point source and 10 milliseconds of arc sources were modeled applying the spectral response of the intermediate bandwidth "y" (as well as the Johnson-V) filter which would actually be used in the observations. The synthetic curves consisted of 1024 data points, evenly spaced in time by one-millisecond. Actually, 2048 data points, spaced one-half millisecond apart were generated and pair averaged to model more appropriately the

implemented data acquistion scheme. The length of these synthetic data sets allowed the examination of their power spectra up to a frequency of 512 Hertz.

Figure 3-11 shows three representative power spectra. Only the first 150 integral frequency components (roughly up to 150 Hertz) are shown. At frequencies higher than this, the power spectra continue to smoothly decay. The amplitude of the power contributions for each of the frequency components has been normalized so that the total power associated with all frequency components up to 512 Hertz is unity.

It was found that as either the source diameter or the spectral bandwidth of the optical filters was increased, the power contributions of the higher frequency components diminished. More importantly, in every case, the power associated with the frequency components higher than 150 Hertz contributed less than 0.001 percent to the total power of the intensity curve. The power spectra scale in frequency linearly with the R-rate. The power spectra shown were for R-rates of 0.3 seconds of arc per second, which is a typical value for most occultations.

Examination of the power spectra of the pre-occultation star+sky intensity (shown for each event in Chapter 5) show that, in general, at frequencies above about 150 Hertz the background noise dominates the observational occultation records. But, in almost all cases, at low frequencies it is the signature of the occultation curve itself which is



Figure 3-11. Power spectra of three representitive synthetic occultation intensity curves.

dominant. This is exemplified quite nicely in the comparative power spectra shown for the occultation of ZC0835 (Figure 5-124). These findings were significant. They implied that better DC solutions might be obtained by subjecting the observational data to Fourier smoothing. Here, Fourier smoothing means transforming the raw data into the frequency domain, removing the high frequency components characteristic of only the background noise (i.e. greater than 150 Hertz), and inversely transforming the frequency-truncated data back into the spatial domain.

Whenever such transformations are applied, careful consideration must be given to the inverse transformation process. Since these numerical transformations do not operate on infinite data sets, spurious results can be obtained due to edge effects. The endpoints of the transformed functions (i.e. ends of the observational data sets) appear as discontinuites in the numerical functions, and often cause "ringing", or high frequency oscillation, to be seen in the re-transformed data. This effect can be seen in Figure 3-12. The upper curve shows the synthetic curve of a 10 millisecond of arc star observed with a Johnson-V filter. This is the curve whose power spectra is shown as the bottom graph in Figure 3-11. The "ringing" effects of inverse transformation, after the high frequency components have been removed from the data in the frequency domain, are shown in Figure 3-12.





The "ringing" effects due to data windowing are often effectively suppressed by the judicious application of an apodizing function, as discussed by Bracewell (1965). In practice, however, the intensity curves fit by the DC and DC2 processes are small subsets of the observational data (typically 200 milliseconds out of 4096). The ringing effects, which are pronounced near the ends of a transformation window become unimportant near the middle of the data window (as exemplified by Figure 3-12). Hence, if 200 data points are to be fit by the DC or DC2 procedure, 1024 data points, roughly centered on the time of geometrical occultation are Fourier smoothed. The data to be fit are then extracted from the Fourier smoothed intensity curve.

This process of Fourier smoothing was applied to all of the aformentioned synthetic curves, for which DC and DC2 solutions had already been obtained. In every case the parameters intrinsic to the synthetic curves were recovered with virtually identical values, but with significantly tighter error estimates. The last set of entries in Table 3-6 reflects this for the DC solution of a Fourier smoothed "noisy" synthetic curve of an 8 millisecond of arc star.

The technique of Fourier smoothing was adopted as a viable data preprocessing method to improve the statistical significance of the solution parameters. This method was not applied blindly, but rather was used in conjunction with the more conventional technique of fitting to the raw data. If

cases had arisen where fitting to smoothed and unsmoothed data from the same observation yielded significantly different results, further investigation would have been needed. Fortunately, in all the observations reduced and analyzed this problem never arose.

CHAPTER IV COMPUTATIONAL DATA REDUCTION PROCEDURES

A Choice of Programming Languages: The APL Decision

The data uploading and preprocessing programs, data reduction algorithms, graphics display software, and all numerical experimentation employed in this investigation were implemented and carried out in APL. APL. an understatedly modest acronym for A Programming Language, is a powerful algebraic notation which was invented by Iverson (1962). An interpreter for this notation was developed by International Business Machines and first released as programming product XM/6 (Falkoff and Iverson, 1968 and 1970), more commonly referred to as APL\360. Since that time, numerous releases of APL interpreters and translators have become available for virtually every mainframe and mini-computer as well as for quite a few microcomputers. Brenner (1982) has recently developed an operating philosophy for a version of APL which would be executable in the parallel processing environment realized in the new generation of "supercomputers".

APL is a concise and totally self-consistent implementation of an array algebra, and is capable of dealing with multidimensional arrays with far greater ease than most other notations and computer languages can deal with simple scalar quantities. A rich variety of intrinsic (referred to

as primitive) functions enable easy array manipulation. A number of operators which act on the primitive functions create, in effect, new functions and enhance the power of this notation dramatically. In addition, APL allows for the creation of "user defined functions" extending the scope of its inherent capabilities only to the limits of the imagination of the user. Excellent treatments of APL as a computer language, including the practical limitations of the notation as implemented on finite machines, are addressed by Pakin and Polvka (1975) and by Gillman and Rose (1974).

It has been found by Schneider and Brown (1976) that the time required for program development and debugging in APL is typically ten to twenty-five times less than in FORTRAN or similar languages. This same study notes that the number of actual lines of "program code" required to carry out equivalent tasks in APL and FORTRAN is reduced by at least the same amount, and in exceptional cases by as much as one hundred times. Furthermore, the transportability of APL programs from one computer to another is accomplished with a decree of success far in excess of any other computer language. A pro-forma APL standard presented by Falkoff and Orth (1979) has been almost universally agreed upon and was adhered to assiduously in the course of this investigation. The most novel aspect of APL is perhaps the notion that the expression of a problem in APL notation is its solution. No extraneous statements heavily laden with arcane syntax as

required by virtually all other computer languages are needed. Given the above, the decision to employ APL in the numerical computations and associated data processing in this project was an easy one to make. It remains a mystery to this investigator why any researcher involved in data processing would make any other choice.

The discussion of the algorithms developed for the solution of lunar occultation intensity curves refers both to the APL functions which have implemented these algorithms (Appendices C, D, and E) and to the underlying equations which have already been presented. This duality in presentation is made to guide the reader unfamiliar with APL notation to an understanding of both the overall data reduction scheme and the computer programs.

Downloading and Uploading of Observational Data

Once the observational data are saved as a contiguous memory image on a CODOS disk file, these data must be transferred to the computer which will be used to perform the computational data reduction. The numerical computations performed on the observations during the course of this investigation were carried out, for the most part, on four mainframe computers. Warner Computer Systems, Inc. (New York) provided access to their Xerox Sigma-9 computer where the initial computational algorithms and basic reduction procedures were developed and tested. A small amount of numerical experimentation was carried out on the Northeast

Regional Data Center's Amdahl 470/V6 and IBM 4341. The latter was done when the MVS operating system was first established and computing was free for a limited time. The majority of computing and the final reduction of all observations was done on the Harris-500 computer operated by the University of Florida's Center for Intelligent Machines and Robotics.

The process of moving data from CODOS disk files onto any of the above mainframes was accomplished by a data transfer program called OCCTRANS. This program (listed in Appendix B) is a 6502 machine language program written to run on a SPICA-IV system equipped with an RS-232 serial port. A SPICA-IV computer, similar to the SPICA-IV/LODAS with the addition of a 6850 asynchronous communications adapter interface (ACIA), which also serves as a backup system for the RHO SPICA-IV/LODAS, was used for this purpose. The communication protocols in terms of baud rates, number of start and stop bits, parity, character prompting, and handshaking are different for each of these machines. As a result OCCTRANS must be modified slightly to establish the proper communications interface for different mainframe computers. The version of the OCCTRANS program listed in Appendix B is for communication with the HARRIS-500 computer.

This program converts the CODOS file which contains the file header and three channel data buffers to ASCII before it is sent to the mainframe computer. Data are transferred in blocks whose size is appropriate to the length of the input

buffer on the receiving computer. Data are received directly into APL via a very simple receiving program called READ, given in the listing of the APL workspace OCCPREP (Appendix C). The OCCTRANS and READ programs "talk" to each other, handshaking and transferring the data. The data, once transferred into the APL OCCPREP workspace, exist as a text vector of length 37376.

This vector is then submitted as the right argument to the function TRANSLATE. TRANSLATE unpacks the ASCII representation of the 12-bit observing data and header information and stores this information as global variables in the workspace. The variables CH1, CH2, and CH3, each numeric vectors of length 4096, hold the contiguous intensity readings taken on each of the three data channels. These data have been "unfolded" from each channel's circulating buffer so that the first element in each vector is the first sample taken in the wraparound data windows. The time and date of the last data sample, the "foldpoint" (sample pointer to the position of the last sample taken in the circulating buffer), and the last comment entered by the observer before saving the data to CODOS disk are also saved in the workspace. The function TRANSLATE and subordinate functions are shown in the OCCPREP workspace listing.

The Occultation Reduction Workspace (OCCRED)

The APL workspace OCCRED (OCCultation REDuction) listed in Appendix D contains the complete set of functions and

global variables needed to reduce a lunar occultation observation. This workspace listing should be consulted in order to elucidate the discussion of the computational procedures which follow.

Global Variables Used by OCCRED

Parameters for the reduction run of the occultation solution are established interactively by the function INPUT. INPUT is monadic and takes as a right argument the global vector of raw observational data (i.e. CH1, CH2, or CH3) created by the function TRANSLATE. The entry of each parameter is prompted for conversationally, and on completion INPUT leaves the established parameters as global variables in the active workspace. A convention of using underscored variable names for parameters passed globally between APL functions is used throughout this and other supporting APL workspaces. Table 4-1 lists the global variables created by INPUT and briefly describes their content. The listing of the function INPUT should be consulted for the physical units of the numerical values. The use of these variables is addressed in the explanation of the differential corrections procedure.

Typically, a maximum of about 250 points (except in the case of very low R-Rates which occur for occultations of near-grazing incidence) is appropriate. The global variables listed in Table 4-2 also reside in the OCCRED workspace and are required by the differential corrections procedure. The names of the filter response matrices are the only global

variables employed whose names are not underscored. The <u>FREN</u> vector is explained in the discussion of the FRESNEL function.

TABLE 4-1 GLOBAL VARIABLES CREATED BY THE APL FUNCTION INPUT

Name	Shape	Туре	Description		
BIN	Scalar	NUM	Estimated Bin Number of Geometrical Occultation.		
<u>COMN</u>	VAR	TEXT	Any Applicable Comments of Special Note		
DATE	VAR	TEXT	U. T. Date of the Observation		
FILT	VAR	TEXT	Name of Filter Matrix to Use		
<u>LIMS</u>	26		Parametric Limitation Matrix, First Row Are Upper Limits, Last Row are Lower Limits, Values are Limits for <u>PV</u> Elements [6] to [11]		
NAME	VAR	TEXT	Name or Catalog Number of the Star		
<u>085</u>	VAR	NUM	Subset of Observational Data to be Used		
PV	14	NUM	The Parameter Vector Indexed as follows:		
			 Central Wavelength of Passband Lunar Limb Distance Number of Observation Points Extracted for Solution. Effective Stellar Temperature Square Root of the Number of GRID Points per Stellar Quadrant Stellar Angular Diameter Pre-Event Signal Level Post-Event Signal Level Ime of Geometrical Occultation L-Rate Maximum Number of Iterations for DC Spectral Width of Filter Matrix Fraction of Adjustments to Apply 		
Note	s: An I VAR	ndex O indica	rigin of 1 is used in this table. tes a variable length vector.		

TABLE 4-2 GLOBAL VARIABLES RESIDENT IN THE OCCRED WORKSPACE

Name	Shape	Туре	Description
NARROWV	5,2	NUM	"y" filter fractional response, normalized, 100 Angstrom steps
VFILTER	53,2	NUM	Johnson V filter fractional response, normalized, 50 Angstrom steps
VFILT	6,2	NUM	Johnson V filter fractional response, normalized, 500 Angstrom steps
NARROWB	5,2	NUM	"b" filter fractional response, normalized, 100 Angstrom steps
BFILTER	42,2	NUM	Johnson B filter fractional response, normalized, 50 Angstrom steps
BFILT	5,2	NUM	Johnson B filter fractional response, normalized, 500 Angstrom steps
FREN	4001	NUM	Diffraction intensity values ordered in decreasing Fresnel units

The Computational Differential Corrections Procedure

The primary function, which sets up and controls the flow of program logic for the solution of occultation intensity curves, is the differential corrections routine called DC. Minor modifications to DC can be made prior to its execution depending upon the nature of the parametric model. For example, the limb darkening coefficient of the stellar atmosphere can be treated as a free parameter. As noted earlier, however, in most cases the model is rather insensitive to changes in limb darkening; hence, it is usually held fixed. Because of the relatively long amounts of CPU time required for computation, "fine-tuning" of the DC procedure to fit the needs of a given observation is felt warranted. A general computational approach may seem more aesthetically pleasing from a programming point of view but is somewhat less efficient computationally. Fortunately, the interactive nature of APL allows variations in the basic DC procedure to be implemented without any special effort. Indeed, alternate execution paths through DC and several subordinate functions are realized by simply turning certain executable statements into comments and vice-versa.

Several global variables are created in the process of determining the occultation solution. In most cases these globals provide information on the solution after completion of the reduction run. A few computational parameters are passed between functions by global variables (though in most cases variables are passed as functional arguments).

The DC solution, in its basic form, is passed the parameter vector as an argument called P. DC[3] begins by establishing two global variables and one local variable. As execution proceeds, the variable <u>SOLS</u> accumulates the adjusted parameter vectors after each iteration step. When execution of DC is completed, <u>SOLS</u> provides an historical record of the various parameter combinations used at each step on the way to arriving at the final solution. The variable <u>SSE</u>, similarly, is an historical record of the sum-of-squares of the residuals from the O-C's after each correction step. DI is a vector containing the fractional variation to be applied to each of the parameters in

numerically computing the partial derivatives of the intensity curve. The index position of each entry in this vector corresponds to the index position of the associated parameters in \underline{PV} . As an example, the intensity curve is fairly insensitive to a small variation in stellar diameter for small diameter sources, hence, DI[6] is typically assigned the moderately large value of 1.02.

DC[4] sets up the iteration counter ITER and the parameter variation control vector CVEC. The first element in CVEC, at any time, determines which parameter is to be varied next in the corrections process. DC[5] computes the spectral response function, R. This response function, corresponding to Equation 3-8, is dependent upon both the spectral energy distribution of the source and the instrumental spectral response function. In order to compute the spectral energy distribution of the source Equation 3-7 is evaluated by the function BBDY.

BBDY takes two arguments. The right argument. L, is a two element vector. L[1] specifies the blackbody temperature, and L[2] specifies the number of terms (i.e. width of each spectral region in Angstroms) to be used in the numerical integration of Equation 3-7. The left argument, W, is a vector giving a list of wavelengths for which the blackbody function is to be computed. BBDY normalizes the integrated blackbody function and returns a vector conformable in length to W, whose elements give the fractional contribution to the

total blackbody power distribution function for the wavelength regions evaluated.

The arguments passed to BBDY are computed by DC[5]. The temperature is taken from the local parameter vector P[4] and the spectral width computed from the first difference of the wavelength specifications of the first two entries in the filter matrix specified by FILT. The wavelengths for evaluation are taken directly from the same filter matrix. After the blackbody distribution function is computed, each wavelength regime is attenuated by an amount representing the filter function also given by the <u>FILT</u> matrix. The modified spectral energy distribution function is the spectral response function, R.

The model strip brightness distribution of the projected quadrant of the stellar disc, as previously discussed. is computed in DC[6]. The computation of the non-limb-darkened grid, which represents surface elements on the stellar disc, is handled by the function GRID. GRID takes a scalar right argument specifying the square root of the number of grid points to be computed for each quadrant on the model stellar disc. DC[6] is passed this value from the parameter vector. GRID carries out the analytic solution (Equation 3-5) for the integration of Equation 3-4 for each projected surface element. The result is a normalized matrix where each element gives the fractional brightness contribution of a projected surface element in the grid coordinate system whose origin is centered at the upper right corner of the matrix,

as shown in Figure 3-2. Only one quarter of the stellar disc is computed, as spherical symmetry is assumed.

The stellar grid is then limb darkened by the function LDARKEN. LDARKEN takes two arguments. The right argument is the stellar grid quadrant as computed by GRID. The left argument is the limb darkening coefficient to be applied. The linear limb darkening law given in Equation 3-6 is applied. The limb darkened quadrant is renormalized and passed as the resulting matrix.

Since spherical symmetry is assumed, the hemispherical strip brightness distribution is computed by DC[6] by doubling the values and summing the matrix along the last coordinate axis. Again applying symmetry, the vector resulting from the previous step when rotated and concatenated onto itself yields the total strip brightness distribution. The normalized limb-darkened strip brightness distribution is stored in the variable §.

DC[7], labeled L0, is the entry point for the outer iterative loop of the differential corrections procedure. This point is re-entered after all parameters have been varied and parametric corrections have been applied to all parameters on the basis of the partial derivatives of the residual matrix. DC[5] establishes a matrix, X, which will hold the numerical partial derivatives of the intensity curve.

DC[8] calculates the computed intensity curve on the basis of the current parameters in P. To do this, the

normalized polychromatic intensity curve must be computed from the linear superposition of monochromatic curves. This is carried out by the function WIDE.

WIDE takes two arguments. The parameter vector is passed as the right argument. The left argument, R, is the previously computed spectral response function. WIDE first calculates the individual monochromatic curve (as an approximation to each spectral region), then sums these curves. The normalized monochromatic curves, which are dependent upon the stellar diameter, velocity parameter and limb darkening coefficient, are computed prior to summation by WIDE[6]. In the case where the limb darkening coefficient is not held constant, the stellar grid brightness distribution must be recomputed at this step. Then WIDE[5] would be "uncommented" and the variable <u>S</u> recomputed in the same manner as it had been in DC[6].

The computation of the series of normalized monochromatic curves which eventually forms the polychromatic curve is carried out in Fresnel space. Thus, the parameters given in physical units of Kilometers, milliseconds-of-arc, Angstroms and so on, are converted by the function FNOS into Fresnel numbers. At each step through the iterative WIDE function, one monochromatic curve is computed. FNOS takes as its right argument the parameter vector with the first element replaced by the wavelength, in Angstroms, of the particular monochromatic curve sought in the current iteration step. FNOS returns a matrix of Fresnel numbers

corresponding to the points of evaluation in time and space for the wavelength and parameter vector specified. The columns of this matrix correspond to the strip segments, \underline{S} , uniformly separated in space in accordance with the stellar diameter and number of strips. The rows of the matrix correspond to the temporally separated points determined from the data acquisition rate (i.e. one millisecond per point) and the velocity parameter. The zero point references are the center of the stellar disc for the columns, and the time of geometrical occultation for the rows.

It is a simple matter to compute the normalized diffraction intensity for a given Fresnel number by numerically integrating Equations 3-2a and 3-2b and applying Equation 3-1. If this were done for each pass through the WIDE function, both the Fresnel sine and cosine integrals would each have to be numerically integrated approximately 50 million times for a typical DC solution. This is, needless to say, prohibitive. Rather, the diffraction intensities, as a function of Fresnel numbers, in the range of Fresnel numbers from -20 to +20 were computed by the function FRESNEL once and stored as a global vector called FREN in the workspace. FRESNEL performed the numerical integration by rectangular approximation to a precision of approximately one part in 5000. Diffraction intensities are stored in FREN in serial order starting with the intensity for a Frenel number of -20 and incremented by 0.01.

Once the Fresnel numbers are computed by FNOS, the corresponding intensities are found by linear interpolation of the <u>FREN</u> vector by the function NPOL. The right argument of NPOL is the array of Fresnel numbers and the result, conformable in shape to the argument array, is an array of corresponding intensities. It was found that simple linear interpolation would produce values correct to one part in 2000 in the worst case.

The matrix product of the intensity array with the strip brightness distribution yields the normalized monochromatic curve. WIDE[6] then applies the response function for the wavelength evaluated to each point in the curve and accumulates a sum of successive curves. The iterative process is completed when all spectral regions have been computed. The execution of WIDE results in a normalized intensity curve with the spectral response function included.

It should be noted that, in principle, WIDE could have been defined non-iteratively; however, the amount of memory space required was far in excess of that available.

DC[8] scales the polychromatic wideband curve for the parametric value of the signal level (in counts) and adds the parametric value of the sky background. The model curve is then stored in the global variable <u>CDMP</u> and subtracted from the selected subset of the raw observations, <u>OBS</u>, to obtain the residuals which are stored in the local variable Y. The sum-of-squares of these residuals for the iteration just completed (or from the initial parametric guesses in the case

of iteration zero) is accumulated in the variable <u>SSE</u> as part of the iterative solution history. The current local parameter vector P is saved in a vector of tentative solutions called <u>SOL</u>.

It may be noted that DC[8] displays both the current parametric values and the sum of the squares of the errors of the just completed iteration step. This output, in general, is somewhat superfluous. However, the Harris-500 computer on which most computations were performed was unreliable and often would crash before a DC run would go to completion. Having this information displayed periodically allowed the procedure to be restarted after the computer was brought back up. At that time, the last iteration values which had been displayed were used as new starting parameters.

The inner iterative loop labeled L1 is entered at DC[9]. During the process of computing numerical partial derivatives of the intensity curve, numerical singularities may occur for domains in which the intensity curve is insensitive to variation of a particular parameter. This problem is particularly acute as the angular diameter of the star approaches zero. Hence DC[9] determines what value of dD (variation in diameter) should be used in computing the numerical partial derivative dI/dD. The values were chosen after some computational experimentation. For source diameters less than 1-millisecond of arc, a variation factor of 1.4 is used in the source diameter. Sources larger than 1-millisecond of arc but less than 2-milliseconds of arc have

their diameters varied by a factor of 1.2. Sources larger than 2-milliseconds of arc but less than 3-milliseconds of arc are varied by a factor of 1.08, and sources with any angular diameter larger than this are varied by 2-percent.

DC[10] begins the process of parametric variation and adjustment. CVEC is rotated by one to point to the successive parameters in each pass through the L1 procedure. The parameter is then varied by the amount specified by DI and resaved back in the P vector. In order to prevent further numerical singularities due to the finite computational precision of the computer, the varied parameters are never allowed to be fuzzed to zero but are limited to a minimum value of 10⁻¹².

DC[11] recomputes the intensity curve in the same manner as DC[8] using the varied parametric value. The initial model curve, <u>COMP</u>, is subtracted from this resultant curve and the residuals stored in the vector D.

DC[13] computes the partial derivative of the intensity curve as a function of the just-varied parameter. Each partial derivative, as it is computed, is saved as a column in the matrix X. The P vector is restored to its previous value which was tentatively retained in <u>SOL</u>. The variation procedure and computation of partial derivatives is continued until all the parameters whose adjustments are sought have been handled. In the case where a previous DC run resulted in a particularly recalcitrant solution (i.e. divergent, or non-convergent through oscillation) the parameter grouping method previously discussed may be employed. In this case partitioning of the adjustment set, as suggested by the commented DC[12], may be executed in place of DC[13] (with appropriate alterations to DC[15]). Here only a subset of the parameters (two or three in a five parameter model) are varied at a time before the L1 process is completed.

The adjustment procedure is terminated by DC[14] when the number of iterations reaches the maximum set in the parameter vector. When this occurs the termination procedure, labeled L2, is executed.

Once all the partial derivatives under consideration in the current iteration of the L1 procedure have been computed, the adjustments to the parameters are found by DC[15]. A least squares fit of the form given by the Taylor series in Equation 3-9 is performed, and the adjustments are applied within the constraints allowed by the limitations matrix, <u>LIMS</u>, to the parameters in the P vector. The amount of adjustment to the parametric values is determined by P[14]. An initial run would have P[14]=1, while a final solution would probably have a smaller value of P[14], (e.g., 0.5).

The iterative solution history accumulated in the variable <u>SOLS</u> is updated by DC[16] and the next iteration of the LO procedure entered.

The L2 procedure at DC(18] is entered on completion of all iterations through the L0 loop. The solution history giving the iteration number, ranking of the solution in terms of the sums-of-squares of the residuals and the successively tried parameter vectors are displayed.

The variance/covariance array of the final adjustments, computed from the residual vector Y, and the partial derivative matrix X are computed by DC[19]. The standard errors of the estimates are found from the variance/covariance array, and displayed. The algorithms for these computations are taken from Smillie (1976). The DC function concludes by assigning the last partial derivative matrix to the global variable <u>PDER</u>. On completion DC leaves a number of global variables in the workspace reflecting the final solution and the history of the DC run. A summary of these variables is given in Table 4-3.

In most cases the last iteration has the lowest ranking, that is, it is the solution which has the smallest sum-of-squares of the residuals. If the variations in successive solutions were small, so that near-asymptotic convergence had been reached, then no further execution is required.

The last iteration, however, need not produce the lowest ranked solution. Meyer (1975) points out that non-linear least squares solutions are not necessarily monotonic. In this case the best solution should be extracted and DC re-executed with an iteration maximum of zero specified in

TABLE 4-3

GLOBAL VARIABLES CREATED BY THE APL FUNCTION DC

Name	Shape	Туре	Description
COMP	<u>PV</u> [3]	NUM	Computed Intensity Curve
<u>707</u>	55	NUM	Variance/Covariance Matrix for the Last Adjustments Made by DC
<u>ER</u>	5	NUM	Standard Errors of the Estimates
PDER	<u>PV</u> [12],5	NUM	Partial Derivatives of the Computed Intensity Curve
<u>SOL</u>	14	NUM	Last Adjusted Parameter Vector Used by DC
SOLS	1+ <u>PV</u> [12],14	NUM	Solutions Matrix, a Chronological History of Successively Adjusted Parameter Vectors
<u>SSE</u>	1+ <u>PV</u> [12]	NUM	Chronological History of the Sum- Square-Errors of <u>OBS-COMP</u> for Successive <u>COMP</u> 's Generated.

the parameter vector. If the DC procedure is not invoked directly, but controlled by the function START, then this is done automatically. This is normally done in all reduction runs.

The solution history must be examined to determine if the last computed set of parameters is sufficiently near convergence as to not require any further adjustment. This evaluation is left to human judgement rather than to a preprogrammed algorithmic decision. If it is desired to continue the DC process then the best solution may be extracted from <u>SQLS</u> and these parametric values used in establishing a new <u>PV</u>. DC may then be re-executed with the adjusted parameters as the new starting values. It may be desirable to reduce the iteration maximum on subsequent runs of DC.

The DC procedure will attempt to fit any data presented to a model occultation curve. Thus the raw observed data may be preprocessed before being submitted to DC without any variation in the operation of the computational DC procedure. This was done in the numerical experimentation performed to evaluate the effects of removing high frequency components in the observed data by foward and inverse Fourier transformation, as well as applying simple N-point and exponential and smoothing to the observational data. The Two-Star Differential Corrections Procedure (DC2)

<u>Global parameters for DC2</u>. The APL function DC2 is a two star analog to DC. The numerical procedures discussed in Chapter 3 have been implemented and are carried out by DC2. Where DC was defined monadicaly, DC2 is defined dyadically. The right argument is a parameter vector, identical in shape and analagous in content to that which would be passed to DC. This parameter vector, <u>PV</u>, would hold the initial parametric guesses for one of the component stars (arbitrarily referred to as star 1) in the two-star system. Here, the pre-event signal level <u>PV</u>[7], is the out-of-occultation signal level of the first star only. The post-event signal level, <u>PV</u>[8], refers to the background sky level, not the star 2 plus sky level.

The left argument passed to DC2 is a similar parameter vector, which will be denoted <u>PV2</u>, and holds the parameter
specifications for the second star. The shape and content are similar to <u>PU</u> with a few minor exceptions. The pre-event signal level, <u>PU2</u>[7] refers to the out-of-occultation signal level of star 2 only. <u>PU</u>[8] is an unused element in the parameter vector, as the background sky level is assumed constant throughout the event. Therefore, the background sky level specified for star 1 is used for star 2 as well.

Some of the elements in <u>PV</u>2 are redundant, as they are identical to those in <u>PV</u> (i.e., the lunar limb distance, number of data points extracted for solution, etc.). Structuring <u>PV</u>2 in a manner similar to <u>PV</u>, however, allowed the same previously defined subordinate functions to be employed with no modifications.

The APL function INPUT2, listed in Appendix D, is used to interactively establish the two parameter vectors used by DC2. In addition to the global variables created by the APL function INPUT (see Table 4-1), INPUT2 creates the variables \underline{PV} [2] and \underline{LIMS} 2. The latter of these is the parametric adjustment limitations matrix for star 2, and has the same arrangement of elements as the \underline{LIMS} matrix.

The APL function DC2. The function DC2, given in Appendix-D, was developed from DC and has a great resemblance to the one-star function. The variable names and algorithmic procedures, where applicable, were retained from the DC function and include the second star. Variables related to star 2 have names which are postfixed with a "2": P2, D12,

R2, X2. <u>SOLS</u>2, and <u>SOL</u>2. They have the same meaning in algorithmic context as the related variables for star 1.

As can be seen, the structure and content of DC2 resembles DC closely. The version of the DC2 function shown retains the capability of either using the Wilson parameter grouping method or solving for all nine parameters simultaneously. The function as shown uses the latter approach. To implement the parameter grouping method those function lines which are preceded by double comment symbols should be uncommented and those lines which are followed by double comments should be turned into comments.

The function shown here defines the same strip brightness vector, \underline{S} (on DC2[7]), to be used for both stars. If different limb darkening coefficients or grid parameter specifications are desired for the two stars, a separate strip brightness distribution, $\underline{S}2$, for star 2 should be computed as well. This information must be passed to WIDE which currently uses \underline{S} for computing the model intensity curve. The simplest approach is to create a function identical to WIDE, called WIDE2, which refers to $\underline{S}2$ instead of \underline{S} . WIDE2 would then be called, instead of WIDE, in the computation of the polychromatic model curve for star 2.

The interim model for the two-star (with sky background) curve, and the two individual one-star curves (under adjustment) which comprise it are computed by function lines DC2[9], DC2[14], and DC2[20], respectively. These are equivalent to DC[8] and DC[14] in the one-star case.

The adjustment procedure has been partitioned into two sections labeled L1 and L1B. L1 handles the variation of the model parameters of star 1 for the numerical computation of the partial derivatives. L1B is the same for star 2. In L1B, PV2[8] is skipped as it does not enter into the model. In the usual fitting process, the labeling of L1B has no significance as it is not a line that is entered conditionally. If the parametric grouping code is activated, however, this becomes an entry point for the parametric adjustment of the subset of parameters associated only with star 2. In that case, the computation of the adjustments after the computation of the partial derivatives for each star is handled by DC2[16] and DC2[17], respectively.

After the iterative adjustment process is completed, DC2 concludes in the same manner as DC. Four additional columns are added to the historical synopsis matrix <u>SOLS</u> (defined on DC2[29]) which are not present on the <u>SOLS</u> matrix created by DC. These contain, for each iteration, the interim values of the adjustable parameters for star 2. The variance/ covariance array, <u>COV</u>, is a 9-by-9 matrix, rather than a 5-by-5. Similarly, the vector of error estimates of the solution parameters, <u>ER</u>, is of length 9. These two variables are computed on DC[30]. The 9 column matrix, <u>PDER</u>, contains the numerical partial derivatives of each of the parameters evaluated for every point of observation.

Preprocessing of the Observational Data

Before submitting the raw observational data to either DC or DC2, Fourier smoothing (as discussed in Chapter 3) may be applied. The dyadic APL function FTSM00TH carries out the Fourier smoothing. The raw observational data, stored clobally in the workspace under the name CH1, must be resident. The right argument is a two element vector. The first element is the bin number of the time of geometrical occultation. In the case of a one star event the variable BIN, created by INPUT, should be used. For a two-star event the mean of BIN and BIN2 is more appropriate. The second element is the number of points centered on the specified time to be included in the Fourier smoothed data set. The left aroument, a scalar, is the cutoff frequency (in Hertz) used in the Fourier smoothing. FTSMOOTH will store the result in the global variable OBS since the smoothed data is to be treated by DC or DC2 as if it were raw observational data.

FTSMOOTH uses the subordinate functions FFT and FFTI. These functions, listed in Appendix E, are also resident in the APL workspace OCCPLOTS and are discussed in Chapter 5. Presentation of the Results of the DC Run

The APL function OUTPUT, coresident in this workspace, produces a formatted report detailing computational results of the occultation solution. The information presented by OUTPUT was transcribed for presentation in the section describing the observational results. OUTPUT makes use of

the global variables created by both INPUT and DC. The information presented by the OUTPUT function is listed in Table 4-4.

The APL function OUTPUT2 produces a similar report for the results obtained from a DC2 fit. In addition to the information presented for DC fitting, the parameters and derived quantities salient to the two-star model are given. These additional items are listed in Table 4-5.

A number of APL functions have been assembled to present a graphical depiction of the solution and other salient information relating to the fitting process. These functions are discussed in detail in Chapter 5. TABLE 4-4

INFORMATION PRESENTED BY THE APL FUNCTION OUTPUT

General Information

1. Name of the star

2. U. T. Date of the occultation

3. Any significant comments to head the output

Input Parameters

1. Central wavelength of the passband (Angstroms)

Topocentic lunar limb distance (kilometers)

3. Effective photospheric color temperature (Kelvins)

4. Limb darkening coefficient

Model Parameters

1. Number of data points used in the solution

2. Number of grid points on the model stellar grid

3. Number of discrete spectral regions modeled

4. Width of each of the spectral regions (Angstoms)

5. Number of iterations of the DC adjustment procedure

Solutions

1. Stellar angular diameter (milliseconds-of-arc)

2. Bin number of geometrical occultation

3. Pre-event (star-plus-sky) signal level (counts)

4. Post-event (limb-plus-sky) signal level (counts)

5. Observed lunar shadow velocity (kilometers/second)

6. Predicted lunar shadow velocity (kilometers/second)

7. Local slope of the lunar limb (degrees)

Statistics of the Solution

- 1. The variance/covariance matrix
- 2. The correlation matrix

3. The sum of the squares of the residuals

4. Sigma (one standard error)

5. The normalized standard error

6. The photometric (Signal-plus-Noise)/Noise ratio

7. The (Change in Intensity)/Background Intensity

8. The change in magnitude

Tabulation of the Observation and the Solution

1. Time from start of solution subset (milliseconds)

2. Observed Intensity (counts)

3. Computed Intensity (counts)

Observed-Computed Intensity (counts)

TABLE 4-5

TWO-STAR QUANTITIES PRESENTED BY THE APL FUNCTION DC2

Solutions 1. Stellar angular diameters for each of the components 2. Bin numbers of geometrical occultation for each star 3. Intensity (in counts) of each star 4. L-Rate determined for each component 5. Local limb slopes at the contact points 6. Background sky intensity Derived Quantities _____ 1. Temporal separation of the stars 2. Intensity-weighted mean L-Rate 3. Projected spatial separation based on: a) Predicted R-Rate b) Determined weighted-mean R-Rate 4. Brightness ratio of the components 5. Magnitude difference of the stars

CHAPTER V THE OCCULTATION OBSERVATIONS AND RESULTS OF THEIR ANALYSIS

Presentation Format

To facilitate the discussion of the occultation observations and the results of their subsequent analysis, the observations, their solutions, and related information are presented in both tabular and graphical form. The order of presentation is chronological, based on the time and date of the observations, beginning with the occultation of ZC0916 (1 Geminorum).

Format and Content of the Tables

The occultation summary table. A synopsis of each event is presented on a table referred to as an occultation summary table. Table 5-9 is an example of an occultation summary table. Each of these tables, headed by the name of the occulted star is divided into six sections. The first of these sections, labeled "Stellar and Observing Information", conveys both the primary characteristics for classifying the star and information with regard to the instrumental configuration employed while making the observation. The primary reference given for the star is either its Robertsons' (1940) Zodiacal Catalog (ZC) number or United States Naval Observatory (USNO) Extended Catalog (X) number. The Smithsonian Astrophysical Observatory (SAO), Bonner

Durchmusterung (DM), and other pertinent catalog numbers or star names are also given. Following these are the Right Ascension and Declination of the star, precessed to the equinox and equator of the date of observation. Also listed are the star's apparent V magnitude, spectral type, and luminosity class (if Known). The instrumental specifications include the filter (designated V, B, "y", or "b"), the diaphragm (designated by the letter codes given in Table 2-3), the amplifier gain setting, and the PMT voltage.

The second section, labeled "Lunar Information", presents the lunar geometry and characteristics which are unique to each occultation event. These data refer to the time of geometrical occultation. The percentage of the lunar disc illuminated (by area) as seen at RHO is given first. This is followed by the moon's elongation from the sun and altitude above the horizon. The topocentric distance to the contact point of the lunar limb is given next. The predicted topocentric linear rate at which the lunar shadow crossed the observatory, and the apparent angular rate of the moon's limb relative to the occulted star are the last entries in this section.

The third section is labeled "Event Information". The U.T. date and predicted time of geometrical occultation are given first. This is followed by the USNO "Value" and "Observability" codes (V/O), as defined by Van Flandern (1973). The Hour Angle (HA) of the star at the time of the event, in units of degrees, minutes, and seconds is given

next. A negative hour angle indicates a pre-transit event. The selenocentric geometry of the occultation, defined by the position angle, cusp angle, contact angle, and Watts angle, (Van Flandern, 1973) of the star concludes this section. Both the lunar and event information presented are derived from data supplied by LuKac (1983, and 1984).

The fourth section, labeled "Model Parameters", presents the fixed computational model parameters and assumed stellar characteristics used in the final differential corrections solution. Listed first is the number of data points (milliseconds of data) extracted from the observation data buffer that were used in the reduction. The number of grid points used to represent the stellar disc of the star, the number of discrete spectral regions employed to model the stellar blackbody function, the instrumental spectral response, and the width (in Angstroms) of those spectral regions are listed next. The instrumental spectral response function used in all cases corresponds to the filter employed. These response functions are given in Appendix D. The assumed stellar limb darkening coefficient and effective photospheric temperature (based on the star's spectral and luminosity classes) are the last entries in this section.

The fifth section, labeled "Solutions", shows the principal computational results obtained from the model fitting procedure. Each solution parameter also has its formal standard error presented alongside in parentheses. The stellar diameter (in milliseconds of arc), and time of

geometrical occultation (in milliseconds) relative to the sample number of the first intensity point taken in the extracted observational subset used in the reduction, are given first. Listed next are the out-of-occultation intensity levels, in raw "counts" (with a full scale value of 4095), before and after the disappearance. The observed topocentric linear velocity of the lunar shadow and the lunar limb slope derived from this velocity are given next. If a decodable WWVB time signal was obtained for the occultation event, the Universal Time of geometrical occultation is presented and is the last entry in this section.

The sixth, and final section of the Occultation Summary tables, labeled "Photometric Noise Information", contains information useful in assessing the photometric quality of the sky during the event. The sum of the square of the residuals is given in raw counts, as is the standard error (Sigma) of the residuals. The absolute unscaled values of these numbers depend both on the number of samples used in the reduction and on the arbitrary full scale intensity value used. To aid in the intercomparison of different observations, a normalized standard error, adjusted for both the size of the data sample and full scale intensity, is also given. The signal-plus-noise to noise (S+N/N) ratio during the time of the observation subset is presented. Out-of-occultation, the S+N/N would be essentially constant

over time intervals short enough not to affect the integrated transparency function of the sky.

The amplitude of scintillation noise, which is the dominant noise source for bright stars, varies linearly with the signal level. Scintillation noise affects only the intensity of the source and not that of the background sky. Photon shot noise, which varies as the square root of the intensity, is normally much less important in the observation of bright stars. Yet during an occultation disappearance the intensity of the source itself is diminished rapidly. The photon arrival statistics are even worse for fast photometry when the arrival time between photons can approach the sample interval. Fortunately, this is in the region of the occultation curve which is well into the asymptotic falloff and is insensitive to these effects.

The sky, however, is typically bright due to scattered moonlight. The noise statistics of this background tend to have a Poisson distribution. Henden and Kaitchuck (1982) discuss a method of observationally determining the noise associated with this background. Due to the stellar disappearance at the time of occultation the dominant noise source changes rapidly from being scintillation dominated Gaussian noise, to background dominated Poisson noise. Hence, the S+N/N ratio would also change. A proper computation of the S+N/N ratio must allow for the switchover in dominant noise sources, and in fact must consider both of them. The scintillation noise is easily measured, and the noise due to the background is determined as suggested by Henden and Kaitchuck (1982). The noise sources are then added in quadrature to determine the final figure presented as the S+N/N ratio.

The time-averaged sky background is, for the purpose of considering the noise sources in an occultation event, simply a zero offset. This offset level varies enormously for different events, depending on the lunar phase and the position angle of the event geometry. Thus, an indication of the quality of the degree of detection is the ratio of the change in intensity (the pre-occultation and post-occultation signal levels) compared to the background sky level. This ratio, listed in this section, is independent of the superimposed scintillation noise. This change in intensity, expressed as a stellar magnitude difference, is the last entry in an Occultation Summary table.

An additional section is inserted into the occultation summary tables, in the case of solutions of "close" binary stars. The derived quantities which are listed in Table 4-5 are presented in this additional section.

<u>Variance-covariance, correlation, and dI/dP_i </u>. The formal errors of the solution parameters, as listed in the occultation summary tables, are derived from the variance-covariance matrix of the residual matrix. In the case where high correlations exist between the model parameters, the formal errors of these parameters are often insufficient for judging the quality of the solution

(Bevington, 1969). This is the case for occultation observations where some of the physical parameters describing the occultation intensity curve have a high degree of coupling. The degree of coupling is a function of the scintillation noise, expressed as the photometric S+N/N, the dynamic range of the signal decrease (change in magnitude), and the timescale of the event (projected lunar velocity) in comparison to the data acquisition rate. Thus, the correlation between the parameters is of interest. The correlation coefficients indicate the degree of sensitivity of a solution parameter to a change in the other parameters. The variance-covariance and correlation matrices for the determined solution parameters are presented on tables of numerical and statistical significance. Table 5-6 is an example of this type of table.

As previously discussed, the regions of sensitivity to variation of the intensity curve with respect to changes in each of the parameters are reflected by the partial derivatives of the intensity curve (dI/dP_i) . It is noted here, however, that while a list of the partial derivatives tabulated at each point of evaluation is too lengthy to present for each observation, the range of values for each of the partial derivatives are given at the bottom of these tables.

<u>Observed and computed intensity values</u>. A list of the actual observations, computed intensities, and their residuals, are also tabulated. The data presented on these

tables are for the subset of observations used in the DC solution. As an example of this type of table see Table 5-6. The first entry corresponds to the first data point in the DC solution subset. These tables are organized in a vertical fashion, and are read from top to bottom within each column of figures.

Format and Content of the Graphs

The graphical representation of the occultations are also of a standard form and are produced by APL functions in the workspace OCCPLOTS listed in Appendix E. As may be seen, the generating functions employ subordinate functions to draw axes, labels, plot points, and perform similar tasks. These functions are not shown here but are discussed by Selfridge (1983). For clarity, each graph will be referred to by the name of its APL generating function. The information presented on each type of graph is discussed here.

<u>Graph of the entire event, RAWPLOT</u>. The first type of graph shows the entire 4096 milliseconds of acquired data for each event. These graphs are produced by the APL function RAWPLOT, an example of which is Figure 5-1. On each RAWPLOT the horizontal axis gives the time of each acquired data point, in milliseconds, relative to the time of the first intensity value (time 0000). The right hand vertical axis gives the actual number of recorded counts scaled from zero to 4096 for each millisecond of observed data. The left hand vertical axis indicates the counts normalized on a scale from zero to one. The plotting device used to draw these graphs

(a Houston Instruments digital incremental plotter) has a resolution of 0.01 inches. Hence, trying to depict all 4096 points on the RAWPLOT would result in several retraces along each step in the horizontal direction. Thus, RAWPLOT displays only every fourth point (1024 in all) of the acquired data.

A cursory examination of a RAWPLOT will give a preliminary indication of the quality of the photometric data. The scatter in the data from point to point results from high frequency scintillation noise. Longer period variation in overall signal intensity comes about from the slower variation in the atmospheric transparency. The approximate time of geometrical occultation relative to time zero can be seen where the intensity data takes a sudden downward step. The relative amount of background light from sky and lunar limb brightness can be seen by noting the post-event intensity above the zero count level. The contribution to the total intensity due to the starlight is noted by comparing the intensity of the signal before and after geometrical occultation.

The integration plot, INTPLOT. The second type of graph, produced by the APL function INTPLOT, is an integration plot of the entire 4096 milliseconds of observed data. These graphs are typified by Figure 5-3. This type of plot was first suggested by Dunham et al. (1973), and may be used to examine the occultation record for evidence of stellar duplicity. The mean intensity taken over all

observed data points is computed and this mean is then subtracted from each observed point. A running sum is then produced from this difference. When plotted against time, a single star disappearance will ascend to a maximum and descend back to zero. A double star, however, will give itself away by the presence of a change in slope in the ascending or descending branches of the curve. Some such detections, as in the case of 1 Geminorum, are obvious. If, however, the second star is of considerably lesser brightness, the detection might be quite subtle, as is evidenced in the INTPLOT of ZC0126. In these cases the change in inflection can be enhanced by choosing a subset of the data, centered on the time of disappearance of the suspected second star to be processed by the INTPLOT function over a shorter time interval. Several of these shorter INTPLOT figures have been generated and included.

<u>Graphic depiction of the best fit, FITPLOT</u>. Graphs such as Figure 5-13, generated by the APL function FITPLOT, each show the best theoretical fit to the observation for the subset of contiguous data points used in the reduction. This fit is superimposed on the observed data. The bottom horizontal axis, as in the case of RAWPLOT, gives the time in milliseconds relative to time zero. The recorded counts are once again presented on the right hand vertical axis. The left hand vertical axis has been normalized such that zero corresponds to the smallest intensity in the data subset, and one to the largest. The relative time of geometrical

occultation is indicated by a dashed vertical line. This time corresponds to an intensity level which is 25 percent above the post-occultation intensity with respect to the pre-occultation intensity. A second dashed line appears in the case of two-star solutions. The computed pre-occultation intensity level (due to starlight, skylight, and lunar limb brightness) is indicated by a short horizontal line extending to the right from the left hand vertical axis. The computed post-occultation intensity level (due only to skylight and lunar limb brightness) is indicated by a short horizontal line extending to the left from the right hand vertical axis.

Across the top of the graph is an additional horizontal axis, scaling the event linearly. This axis gives the distance, in meters, of the projection (along the Earth's surface) of the geometrical shadow from the telescope. The zero reference is taken at the point of geometrical occultation. Negative values along this axis correspond to pre-occultation intensities.

Noise statistics of the observation, NOISEPLOT. The fourth type of graph (of which Figure 5-6 is an example) is generated by the APL function NOISEPLOT and shows the noise figure of the event throughout the entire 4096 milliseconds of data. The observed intensity values are subtracted from the computed values and the residuals binned into fifty classes. Each class width is two percent of the range of the residuals. The binned residuals are plotted against the number of residuals per bin. The horizontal axis indicates the value of the residuals as a percentage of the mean value. The vertical axis indicates the number of residuals per bin. A dashed vertical line shows where the mean of the distribution falls. The one sigma width of the distribution is shown as a horizontal error bar centered on the line representing the mean value.

The non-linear least squares, differential correction procedure, employed to obtain the occultation solution parameters, assumes that the residuals are randomly distributed. Thus, the noise figure should be essentially Gaussian. Often, the distribution function of the residuals, taken across the entire data set, seems to have a small negative-going Poisson tail. This results from an increase in background light during the 4096 milliseconds of data acquisition. The two major sources contributing to a time varying background are: Earthshine along the lunar limb (dominant for events of small solar elongation), and the radial brightness distribution of the lunar aureole (dominant for events occuring at large solar elongations). The background sky intensity determined as one of the solution parameters reflects the light-level at the time of the event. and typically, is somewhat higher than the background observed earlier in the observation. The Poisson tail arising in the residual distribution functions due to these effects is usually inconsequential, as the effect is very small over the short timescale of the occultation itself.

<u>Power spectra</u>, <u>POWERPLOT</u>. While the character of the distribution function of the residual amplitudes is important in confirming that the fundamental assumption of stochastic noise implicit in the reduction process is true, so too is the distribution of the power components associated with the observed data. The question of an optimum data acquisition rate, effective system time constant for the analog electronics, and their relation to scintillation noise and the power spectrum of an occultation observation is discussed in Chapter 3. Figures generated by the APL function POWERPLOT contain three power spectra for each occultation, and show the relative power contributions for all frequency components up to 500 Hertz.

Figure 5-7 is an example of a POWERPLOT. The lower power spectrum in a POWERPLOT results from a discrete Fourier transformation of 1000 millisecond observations centered on the time of geometrical occultation. These may be compared with the middle spectrum, which for the same event shows the pre-occultation power spectrum derived from 1024 millisecond observations typically centered 1500 milliseconds from the time of geometrical occultation. The 512 millisecond separation in time between the endpoints of the pre-occultation and occultation power spectra are sufficiently separated in time so diffraction fringing effects will not corrupt the pre-occultation data. Also, in this short time interval the atmospheric conditions cannot appreciably change and thereby affect the character of the

seeing or scintillation. The upper power spectrum is that of the computed occultation curve, also centered on the time of geometrical occultation. In a few cases the data acquisition process was halted too late to to preserve 1024 milliseconds of pre-occultation data. In these cases the out-of-occultation data subset is shorter, and hence, the power spectrum is generated to correspondingly lower frequencies.

The horizontal axes give the frequency components in Hertz. The power spectra are plotted on a logarithmic (base 10) scale. Each decade is marked along the vertical axes. All three power spectra, on each POWERPLOT, are normalized to the same logarithmic power scale. Most of the Fourier transformations were carried out using the APL functions FFT and FFTI (listed in Appendix E), supplied by Selfridge (1984). These are "fast" Fourier transformations and require data to be transformed to have a number of elements evenly divisible by a power of two. In cases where this was not possible or practical, the discrete Fourier transformation functions FT and INVFT were used. These functions are discussed in detail by Schneider (1981).

Examination of the three power spectra for each event will reveal which frequency components are important to the occultation intensity curve and which are primarily associated with the background (i.e. scintillation) noise. It can be seen in all cases that the frequency components in excess of about 150 Hertz are generally unimportant in the

occultation curve itself. These graphs also allow one to quickly spot any induced artificial noise, such as a 60Hz component (which could result from improper electrical shielding or loss of a ground or other source of AC pickup in the analog electronics).

<u>Sensitivity of solution to variation of parameters.</u> <u>PDPLOT</u>. Of interest in the final solution to an occultation observation is the sensitivity of the computed curve to variations in the solution parameters. In the computational process of determining differential corrections, numerical partial derivatives of the intensity curve at each observed point are computed for each parameter considered in the model. These derivatives, dI/dP₁, are retained for the best determined theoretical fit. Where each of these derivatives is at a maximum corresponds to an area of the computed intensity curve of high sensitivity to perturbations in the corresponding parameter.

Graphs, such as Figure 5-14, generated by the APL function PDPLOT, show dI/dP; for the variation of the solution parameters. The partial derivative curves are labeled as follows: PREI = Pre-event Intensity ; POST = Post-event Intensity ; TIME = Time of Geometrical Occultation ; DIAM = Angular Diameter of Stellar Disk ; and VELO = Velocity of Lunar Shadow Passage. In the case of two-star solutions the labeling of the derivative curves are postfixed with a "1" or "2" as is appropriate. Each curve has been normalized from -1 to 1, with the line corresponding to $dI/dP_i = 0$ drawn in. The actual values of dI/dP_i differ by orders of magnitude for the five parameters. Thus, the relative degree of sensitivity of the intensity curve is not easily seen. The maximum and minimum values for dI/dP_i are given in the tables of supplementary statistical information.

Discussion of Individual Occultation Events 200916 (1 Geminorum)

<u>Historical notes</u>. Few binary stars have a detectable duplicity both visually and spectroscopically due to observational selection effects. Fortunately, though only a handful of such stars are known, they do exist and 1 Geminorum is among them. 1 Gem is an interesting star, with perhaps an equally interesting history. It was first discovered to be a spectroscopic binary by Campbell and Moore (1907) from radial velocity data obtained from plates taken at Lick Observatory. A period of 9.590 days was first given to the spectroscopic variation by Young (1919) from 77 plates taken in Victoria and Ottawa, Canada.

Though there had never been any previous visual detection of duplicity, Kuiper (1948) found 1 Gem to be double while using the star to focus the 82-inch telescope at McDonald Observatory. Using a Cassegrain spectrograph, he found the two visual components to be late giants of similar spectral type. Since that time many observations of the visual pair have been carried out with 58 listed by Worley (1984) through 1981.12. Most of these, until 1976, are

visual micrometer observations. More recent observations by McAlister (Worley, 1984) have been via speckle interferometry.

Heintz found a period for the visual pair of 13.17 years, as reported by Muller (1961). The orbital elements he determined, along with the associated Thiele-Innes constants, were listed by Finsen and Worley (1970).

Abt and Kallarakal (1963) first elucidated the nature of the 1 Gem system. They determined that the fainter visual component, 1 Gem B, was the spectroscopic binary. They further determined the spectroscopic period to be 9.6595 days, with a mass function of 0.139 solar masses. In addition, the AB visual system was given a period of 11.00 years, with an eccentricity of 0.255 and a total mass of 3.16 solar masses.

Griffin and Radford (1976) rexamined the 1 Gem B spectroscopic binary in the light of photoelectric radial velocity measurements. Their findings, while not substantially different from Abt and Kallarakal, do indicate a period of 9.59659 days (closer to Young's value), and a circular orbit as well. Additionally, they indicate that the dips exhibited in their photoelectric radial velocity data are of "very unequal depth", from which they suggest that the B component is of "substantially earlier spectral type than A" (1976, p. 191).

If a fourth component in this system does indeed exist, as the occultation observation seems to indicate, then the

question of final determination of spectral types may still be open.

<u>The observation</u>. On 22 March 1983, a lunar occultation of the KOIIII/G8III-VII (Abt and Kallarakal, 1963) visual-spectroscopic binary star 1 Geminorum (BS 2134, HR 2134, HD 41116, KUI 23) was observed under favorable circumstances from RHO. The specifics of the observation are given on Table 5-1. The RAWPLOT of the occultation

TABLE 5-1 THE OBSERVATION OF THE LUNAR OCCULTATION OF ZC0916

STELLAR AND OBSERVING INFORMATION

Star: ZC0916 (1 Geminorum, SAO 077915, DM +23 1170) RA: 060305 DEC: +231556 mV: 4.30 Sp: G5 Filter: V Diaphragm: Gain: C11+ Voltage: 1000

LUNAR INFORMATION

Surface Illumination: 50 percent Elongation from Sun: 90 degrees Altitude Above Horizon: 48 degrees Lunar Limb Distance: 367212 Kilometers Predicted Shadow Velocity: 716.3 meters/sec. Predicted Angular Rate: 0.4024 ancsec/sec.

EVENT INFORMATION

Date: March 22, 1983 IIT of Event: 02:44:34	
01 01 Eventte 02111101	
USNO V/O Code: 19 HA of Event: +464840	
Position Angle: 114.0 Cusp Angle: 66S	
Contact Angle: -21.8 Watts Angle: 112.5	

observation is shown Figure 5-1, and a subset of these data (renormalized to full scale) showing the occultation in detail is presented in Figure 5-2. The stepped disappearance, well separated in time, is quite obvious. The brighter star (the A component of the visual system) underwent disappearance first. Examination of either the







Figure 5-2. Detail of the ZC0916 double disappearance.

complete or detailed integration plots. Figures 5-3 and 5-4, shows no evidence of any disappearances other than those of the A and B visual components.

<u>Reduction and analysis of the 1 Geminorum A observation</u>. The diffraction curves resulting from the occultations of the A and B components were sufficiently separated to allow treating each of these independently. It can be seen from the extracted data set that the B star is so far into the asymptotic falloff of the A star curve that the A star contribution is negligable.

Inspection of the A star curve immediately reveals irregularites unanticipated in the disapearance of a single star, but of the sort commonly seen in the occultation of a close double star. Attributing this to rapid atmospheric variations is difficult, as the sky was well behaved through the other 3900 milliseconds of data taken.

Despite the fact that there was an appearance of duplicity, an attempt was first made to fit the A curve to a single star model with the single-star DC program. Though liberal constraints were applied to the differential corrections, no solution could be obtained. An attempt was then made to fit the data to a two-star model, using the DC2 program.

The initial run of the two-star model allowed for the variation of nine parameters: the diameters of the two suspected stars, their times of geometrical occultation, their unocculted intensity, the local slopes of the lunar







Figure 5-4. Integration plot detailing the double disappearance.

limb (actually the projected component of the lunar velocity vector) at the contact points of geometrical occultation, and the background skylight. The slopes were presumed to be independent as the to-be-determined separation might be sufficiently large to allow this.

One hundred milliseconds of data were used in the reduction of 1 Gem A, as indicated in Table 5-2.

TABLE 5-2 TWO-STAR SOLUTION OF THE OCCULTATION OF 1 GEMINORUM A

MODEL PARAMETERS				
Number of Data Points: 101 Number of Grid Points: 4096 Number of Spectral Regions: 5 Width of Spectral Regions: 50 Angstroms Limb Darkening Coefficient: 0.5 Effective Stellar Temperature: 5400 K				
SOLUTIONS				
Brighter Star Fainter Star				
Stellar Diameter (ams): 0.73 (2.42) 0.00 (FIXED) Time: (Relative): 54.82 (0.37) 34.02 (2.46) Pre-Event Signal: 372.60 (5.6) 120.30 (12.1) Velocity (m/sec): 752.3 (9.0) 621.6 (76.1) Lunar Limb Slope (deg): -8.89 (0.57) +14.9 (6.3) Background Sky Level: 429.8 (6.8) Separation in Time (milliseconds): 20.80 (2.49) Projected Spatial Separation (ams): 8.41 (1.01)				
Sum-of-Squares of Residuals: 116250 Sigma (Standard Error): 36.54 Normalized Standard Error: 0.5797 Photometric (StN)/N Ratio: 2.725 Total Intensity/Background: 1.147 Change in Magnitude: 0.149				

The data set initially was chosen to include sufficient diffraction fringing effects on the leading edge, and

truncated so as not to include significant effects from the 1 Gem B curve. Additional reduction runs containing earlier data points in the sample did not significantly alter the solutions. The solution obtained by the DC2 procedure indicated that both stars were essentially point sources, but a large formal error was attached to the diameter of the fainter star. This is understandable when one notes that the fainter star contributes only 1/3 of the total light to the A curve. Hence, the solution is, unfortunately, insensitive to the diameter of the fainter star. With this in mind the program was re-run, but this time holding the diameter of the second star fixed as a point source. It was found, as might have been anticipated, that the resulting solutions were changed insignificantly, though the formal errors were widened slightly in most of the parameters.

The parametric solutions are also presented in Table 5-2. The difference in time between the geometrical occultations for the two A stars was 20.8 (+/- 2.5) milliseconds. This corresponds to a separation of 8.37 (+/- 1.00) milliseconds of arc projected along a line in the direction of the moon's topocentric motion. The observational evidence is quite strong that there almost certainly is a detection of a fainter star in the 1 Gem A occultation record.

The best two-star solution obtained for the 1 Gem A observation is depicted graphically on the FITPLOT, Figure 5-5. The sensitivity of this solution to variations





of each of the parameters (i.e., the first derivatives of the intensity curve) is shown on the PDPLOT, Figure 5-7.

Reduction and analysis of the 1 Geminorum B observation. A similar approach was taken for the solution of 1 Gem B. Here, however, DC2 was modified to determine dynamically if a one, or two-star solution would represent a better choice. This decision was made at the end of each differential correction step. It was felt, in the light of the fact that the fainter component of 1 Gem B had never been detected. that a one star solution would be more appropriate. Nevertheless, the model would invariably return to re-entering the additional parameters necessary for a two-star solution. Forcing a one star model resulted in a guasi-convergent solution: that is, the parameters exhibited large amplitude oscillations about some mean values, rather than asymptotically approaching final values. The formal errors and the sum square error of this "solution", and a visual inspection of the best obtainable fit were significantly worse than the two-star model. The two-star solutions for 1 Gem B is presented in Table 5-3, and graphically in Figure 5-6.

Note here that the fainter, previously undetected, spectroscopic component of 1 Gem B is almost ten times fainter than the brighter star. This led to a complete insensitivity of the solution to the diameter of the fainter star. Its detection, however, is quite evident, and the fit is made possible only with the change in level provided by







the occultation of the fainter B star component. The PDPLOT for the 1 Gem B solution is presented as figure 5-8.

TABLE 5-3. TWO-STAR SOLUTION OF THE OCCULTATION OF 1 GEMINORUM B

MDDEL PARAMETERS Number of Data Points: 101 Number of Grid Points: 4096 Number of Spectral Regions: 5 Width of Spectral Regions: 50 Limb Darkening Coefficient: 0.5 Effective Stellar Temperature: 5400 K

SOLUTIONS

Brighter Star Fainter Star

Stellar Diameter (ams):	5.89	(0.84)	Point Source
Time: (Relative):	61.00	(0.45)	52.69 (3.05)
Pre-Event Signal:	301.30	(15.1)	33.50 (15.0)
Velocity (m/sec):	773.5	(9.0)	787.5 (72.6)
Lunar Limb Slope (deg):	-11.1	(1.1)	+12.3 (5.8)

Background Sky Level: 115.9 (3.3) Separation in Time (milliseconds): 8.31 (3.08) Projected Spatial Separation (ams): 3.34 (1.24)

PHOTOMETRIC NOISE INFORMATION

Sum-of-Squares of Residuals:	348100
Sigma (Standard Error):	18.754
Normalized Standard Error:	0.0857
Photometric (S+N)/N Ratio:	12.68
Total Intensity/Background:	2.89
Change in Magnitude:	1.15

A discussion of the 1 Geminorum results. The

distribution function of the residuals of the four-star fit to the observed data is shown in the NOISEPLOT, Figure 5-9. The power spectra of the observed and computed four-star fit, as well as the power spectrum of the pre-occultation signal, can be seen on the POWERPLOT, Figure 5-10.




For purposes of clarity in this discussion the brighter component of 1 Gem A is referred to as A1 and the fainter as A2. Similarly denoted is the brighter member of 1 Gem B as B1 and its unseen spectroscopic companion as B2. The intention in doing so is not necessarily to imply that the A stars are actually a physical pair (i.e., that they orbit about one another, as the B stars do), but the question does of necessity arise.

There are, then, three possibilities if the existence of 1 Gem A2 is accepted: first, that it is a member of a physical pair with A1; second, that it is not, but rather is in an independent orbit about the barycenter of the 1 Gem system; or third, that it is simply, by chance, a bright star in the same field. The third possibility is highly unlikely, as the projected separation of the A components is only 8.4 milliseconds of arc.

One must ask, before addressing the question as to whether the first or second case is the correct one, why the A2 star has not made its presence known through its spectral lines or its possible effect on the radial velocity curve of 1 Gem A1. Table 5-4 shows the relative brightnesses and magnitudes of the four component stars in the 1 Gem system, on the basis of the occultation observation.

The derived individual magnitudes come from a total systemic V magnitude of 4.18 (Eggen, 1965). With the A2 star 1.23 (+/- 0.02) magnitudes fainter than the B1 star, and contributing only 15 percent of the total light of the

system, it is quite understandable that its spectral signature does not make itself readily apparent.

Star	Intensity	%Total	mV	⊿imV
A1	372.6	45.02	5.05(+/-0.02)	
A2	120.3	14.53	6.27 (+0.12,-0.10)	1.22
A1+A2	492.9	59.55	4.74	
B1	301.3	36.40	5.28 (+0.06,-0.05)	
B2	33.5	4.05	7.66 (+0.64,-0.40)	2.40
B1+B2	334.8	40.45	5.19	0.43

TABLE 5-4 1 GEM: COMPONENT CONTRIBUTION TO TOTAL SYSTEM INTENSITY

As to its effect on the long period radial velocity curve of 1 Gem A, it need have none if A1 and A2 are a physical pair with an inclination near zero degrees. It would, however, be presumptuous to assume, a priori, that this is the case. If A2 has an independent orbit about the barycenter of the system, one still might see an effect in the long period radial velocity curve of the AB system. To resolve this question one will have to look carefully at the available radial velocity data, with an eye to the idea that a fourth component most probably exists. This is a matter for future investigation. It is of interest to note that in their exposition of 1 Geminorum. Abt and Kallarakal conclude that "the astrometric and spectroscopic observations are consistent with a multiple system in which the primary is a KOIII star that may be single" (1963, p. 144), while Griffin and Radford state that the system is "at least triple" (1976. p. 189). In either case the previous observations do not rule out the existance of this newely discovered component.

In principle an occultation observation itself can provide a constraint on the orbit of a visual binary. The position angle, ρ , and angular separation θ , are inseparable from a single occultation observation alone. Only a projected separation, along a line in space can be found. However, given an orbit with its associated errors, the intersection of that line with the orbit provides an independent, highly accurate (though only solitary) observation. Unfortunately, with the addition of a fourth member in the 1 Gem system the situation is not quite as clear cut.

If 1 Gem A is assumed to be a physical pair, that is to say a hierarchy 2 quadruple system a la Evans (1968), then the visual observations of the AB pair would have been made not with respect to the AB positions, but with respect to the centers-of-light (CL's) of the A and B component systems respectively. In the case of 1 Gem B the difference is negligible, but not necessarily so for 1 Gem A.

Figure 5-11 shows the chords both in time and space containing the A and B stars. The figure is oriented so the topocentric radial direction of approach of each star to the smooth lunar limb is vertical. Thus, the stars lie somewhere along the horizontal chords. The intersection of the horizontal chords with the time scale on the left is simply the relative times of geometrical occultation. The bottom



which is tangent to the lunar limb. Their separation, both in time and The angles shown, and the line labeled Pspace are given along the vertical axes. The centers-of-light of the A1-A2 and B1-B2 chords are shown as dashed lines, and their measured These chords are tangent to a line chords labeled A1, A2, B1, B2. separation is indicated by D. are explained in the text. Figure 5-11.

line is both the chord of A2 (which disappeared first) and the smooth lunar limb at the time of A2's geometrical occultation. The top line corresponds to B1's disappearance, which happened last. From the measured intensities of each of the four stars, the CL's of the A1,A2 and B1,B2 pairs were computed as the intensity and time-weighted-means of the respective components. Hence, the time of geometrical occultation of the CL of the A system is

CL(A)=[(54.82x372.6)+(34.02x120.3)1/(372.6+120.3) and is similarly found for the B1,B2 stars. The absolute times in terms of the total data window and reduced to Universal Time are given in Table 5-5. The dashed lines shown on Figure 5-11 are the chords of the CL's of the A and B pairs. The direction to North is 24 degrees westward of the horizontal (which is tangent to the smooth lunar limb).

		TABLE 5-5	5
TIMES	0F	GEOMETRICAL	OCCULTATIONS

	A-VISUAL COMPONENT		B VISU	B VISUAL COMPONENT		
Relative	Times:	3382 msec	+	3450 msec =		
Star 2 CL Star 1	34.02 49.87 54.82	3416.02 3431.87 (0.77) 3436.82	52.69 60.16 61.00	3502.69 3510.16 (0.61) 3511.00		
U.T.'s: Star 2 CL Star 1	02:44:3 02:44:3 02:44:3	3.3258 (0.0027) 3.3417 (0.0013) 3.3466 (0.0012)	02:44:3 02:44:3 02:44:3	3.4125 (0.0032) 3.4200 (0.0013) 3.4208 (0.0013)		

The projected separation in time between CL(A) and CL(B) was 78.29 (+/- 0.98) milliseconds. Since the projected radial approach rate was 0.4024 seconds of arc per second, then the projected angular separation. D. is 31.5 (+/- 0.39) milliseconds of arc. Next. it is logical to inquire if this is consistent with the 1 Gem orbit derived from visual-micrometer measurements and with recent speckle interferometric observations.

Using the Thiele-Innes constants for Heintz's orbit, as catalogued by Finsen and Worley (1970), the predicted P and θ for the date of the occultation observation. 1983.219, were found to be 0:0855 and 68.1436 degrees respectively. Since this orbit was derived from observations predating 1962, it was of interest to determine how well more recent observations of P and θ compared to those predicted using Heintz's orbit. For this purpose a compendium of fifty-eight observations up to and including 1981.120 were obtained from Worley (1984). Prior to 1976, most observations were micrometer measurements. Since that time the majority reported were speckle interferometric observations made by McAlister. It was immediately noticed that over the last decade the θ residuals were running negative and increasingly so with time.

In order to define any secular changes quantitatively in both the P and θ residuals, the O-C's for all observations were fit via least squares, along with the dates of the corresponding observations, to seven different curves. The curve which fit best, in terms of the sum of the squares of the residuals as well as the Pearson product-moment correlation coefficient, was a third order polynomial.

Figure 5-12 shows the observed and computed residuals for both P and Θ .

There is no intention to suggest that there is any physical meaning to any of the particular coefficients. Rather, the fitting was done to provide a reasonable method of extrapolation, on the basis of the O-C's for the orbit, for corrections to be applied to predicted values of P and θ for the date of the occultation observation. In doing so corrections of -22.099 for θ , and +0.004547 for P were found. The corrected quantities are denoted $\theta' = 4\delta_0^0 04$ and P' = 0.090047.

The position angle of the occultation along the lunar limb was at a PA = 114 $^{\circ}$ 0. Thus, the angle measured eastward between the line defining the smooth lunar limb (tangent to PA = 114 $^{\circ}$) and the θ' direction joining the A and B CL's is $\phi = \theta' - (PA - 90^{\circ})$, or in this case 22 $^{\circ}$ 04. The projection of θ' along the line defining the topocentric radial direction of approach is simply $D' = P' \sin \phi$, or 33.8 milliseconds of arc. This is in good physical agreement with the observed projected separation of 31.5 milliseconds of arc, the difference being only 7.3 percent.

With regard to 1 Gem B, the positional information derived from this observation will probably not contribute much to the Knowledge of the short period spectroscopic binary. Here, as in the case of 1 Gem A, the derived separation (8.31 (+/- 3.08) milliseconds, or 3.32 (+/- 1.24) milliseconds of arc) is a projection along the



Figure 5-12. The 0 and P residuals based on Heintz's determination of the 1 Geminorum visual orbit are shown for 58 observations as listed by Worley. The best fit to the residuals (a third order polynomial) is also shown. The numbers or letters indicate the following observers:

0-Kuiper 1-Wilson 2-Muller 3-Van Beisbroeck 4-Finsen 5-Heintz &-Couteau 7-Worley 8-Van Den Bos 9-Baize A-McAlister 8-Tokovinin C-Morel D-Extrapolation to 1983.219

radial direction of lunar approach. Since the inclination of the orbit (as well as the orientation of the major axis of the projected orbital ellipse) is unknown, the relative positions of the individual components cannot be determined. Since, however, the fainter star disappeared first, it must have been essentially to the west of B1. Knowing the lunar PA of the event, one can then say that the B1 star was at a position angle greater than 24° , but less than 204° , measured eastward from North of the B2 star. Finally, it is noted from Griffin's period and (TO) = MJD 2440443.129 that this observation occured at phase 0.15 (+/- 0.03) of the short period radial velocity curve.

<u>Concluding remarks on 1 Geminorum</u>. The complete solution to the occultation of all four stellar components is presented graphically in Figure 5-13. Tables 5-7 and 5-8 give the supplemental statistics of the A and B component solutions, including the variance/co-variance and correlation matrices and the numerical ranges of the partial derivatives.

As a result of this observation there is now strong evidence pointing to the existence of a fourth component in the 1 Geminorum system. While this previously undetected component is more than three times fainter than 1 Gem A1, and almost six times fainter than the combined light of 1 Gem A1 and 1 Gem A2, supporting evidence might be found by a careful examination of speckle interferometric and radial velocity observations.



TABLE 5-6 ZCO916: OBSERVATIONS, COMPUTED VALUES, AND RESIDUALS FROM BIN 3362

								=====			
NUM	OBS	COMP	RESID	NUM	OBS	COMP	RESID	NUM	OBS	COMP	RESID
	0BS 905 928 974 978 978 987 978 9867 978 9867 9366 1004 10060 9601 1004 10663 9916 8990 8916 8990	COMP 920.3 937.5 937.5 937.5 927.7 921.2 917.8 910.2 918.2 939.9 918.2 939.9 939.9 939.9 938.8 939.9 938.8 939.5 937.5	RESID -29.3 -20.5 36.5 87.3 55.6 57.4 82.8 73.6 120.7 -24.2 24.8 82.8 73.6 120.7 -48.8 -48.8 -48.8 -68.3	===== NUM 644 656 667 689 700 71 722 774 75 76 777 78 90 81	OBS 9088750 98087501 660385562 50043660 55043660 44418452 465494 44524452 46497	COMP 826.2 789.1 752.2 716.5 588.0 575.5 555.6 523.7 510.9 490.5 449.5 449.5 445.6 449.4 848.4 469.8	RESID 82.8 18.9 -2.2 14.5 -25.4 -33.55 -55.8 -34.6 -20.7 -55.8 -34.6 -20.7 -56.9 -20.7 -58.9 -42.5 -58.9 -42.5 -11.6 -10.6 -10.6 -22.2	NUM 1227 128 127 1311 1323 1334 1355 1367 138 1390 141 1423 1441	0BS 401 480 507 5404 4937 5507 5524 477 4291 33324 33324 33324	COMP 420.7 435.8 457.4 481.8 505.5 524.6 537.3 526.5 538.2 508.0 484.2 456.9 3366.1 337.1 336.1 337.1 326.4 9 28.4 9 28.4 9	RESID -19.7 25.2 22.6 11.2 -17.6 2.7 -17.6 2.8 2.3 2.8 2.3 2.8 2.3 2.4 6.0 -7.2 -9.6 .2 .1 1.4 .7 -7.2 -10.1 1.4 .7 -10.1 2.9 2.8 2.3 2.4 .5 2.4 .5 2.4 .5 2.4 .5 2.4 .5 2.5 2.5 2.5 2.5 2.5 2.5 2.5 2.5 2.5
10012234552722222222222222222222222222222222	825 8811 803 8966 8888 8967 8888 8967 931 11112 9924 9991 1081 1112 9924 9950 9951 9950 9950 9950 9950 9950 9950	3269.6.2 9094.9.9 9094.9.9 9145.5 9269.2 9269.2 9269.2 9265.0 9265.0 9269.2 920		3233456788999923345687889999999999999999999999999999999999	462 4877 4660 4404 4900 44037 4261 4099 40537 4099 40537 4099 40537 4099 40537 4099 40537 4099 40537 4099 40537 4054 4054 4057 4054 4054	$\begin{array}{c} -2.6\\ -2.6\\ +2.5\\$	101298195287361562477591363287306 1028887444298825271441200314416588 11028887444298825271441200314416588	144789012334567890123444789012315567890123444489012553456789012345567890123455567890123455567890123455567890123555555555555555555555555555555555555	273 2771 196 181 1799 1688 1799 1682 159 159 159 151 151 151 151 151 151 151	$\begin{array}{c} 2622 \cdot 4 \\ 222096 \cdot 0.6 \\ 844 \cdot 8.5 \\ 1659 \cdot 4.4 \\ 1659 \cdot 4.4 \\ 1659 \cdot 4.4 \\ 1437 \cdot 1.5 \\ 1320 \cdot 6.9 \\ 1233 \cdot 4.4 \\ 1444 \cdot 4.1 \\ 1337 \cdot 1.5 \\ 1320 \cdot 6.9 \\ 1233 \cdot 4.4 \\ 1337 \cdot 1.5 \\ 1320 \cdot 6.9 \\ 1233 \cdot 4.5 \\ 1234 \cdot 5.5 \\ 1232 \cdot 6.9 \\ 1233 \cdot 6.9 \\ 1234 \cdot 5.5 \\$	1081-80.062.554.66.94.37.5.61.553.23.52.83.82.41 -280-24.17.66.69.4.37.5.61.553.23.52.83.82.41 -1.52.50.23.52.83.82.41 -1.15.25.50.153.23.52.83.82.41 -1.15.25.50.153.23.52.83.12.20.24 -1.12.52.50.153.23.52.83.12.20.24 -1.12.52.50.153.23.52.83.12.20.24 -1.12.52.52.53.52.12.53.12.20.24 -1.12.52.53.52.23.52.52.52.52.52.52.52.52.52.52.52.52.52.
490123345567 555555555567 556789012 =	723 716 810 840 768 935 932 928 936 958 958 958 959 959 937	774.95 800.42 830.42 864.63 926.64 955.83 926.64 955.83 942.60 922.04 894.48 924.48 924.48	-53.09 -53.99 -299.62 -299.62 -299.62 -29.62 -18.64 -278.34 -278.34 -18.34 -278.34 -18.34 -278.34 -18.34 -275.22	111 112 113 114 115 116 117 118 120 121 122 123 124 124	4199 402 449 449 443 443 485 432 457 431 409 436	464.4 448.4 427.4 432.10 463.6 478.4 840.6 467.6 480.6 467.6 480.6 467.6 480.6 490.6 400.6 40.6 4	34.6 -14.8 -21.4 -30.1 15.4 -35.4 -35.6 -14.8 -35.6 -14.8 -35.6 -14.8 -35.6 -14.8 -35.6 -14.8 -35.6 -14.8 -35.6 -14.8 -30.0 -14.8 -30.0 -14.8 -30.0 -14.8 -30.0 -15.4 -30.0 -15.4 -30.6 -15.4 -30.6 -15.4 -30.6 -15.4 -30.6 -15.4 -30.6 -15.4 -30.6 -15.4 -30.6 -15.4 -30.6 -15.4 -30.6 -15.4 -30.6 -15.4 -30.6 -15.4 -30.6 -15.8 -10.1 -15.4 -30.6 -15.8 -10.1 -15.4 -30.6 -10.1 -15.4 -30.6 -10.1 -15.6 -10.1 -15.4 -30.6 -10.1 -15.4 -30.6 -10.1 -15.4 -30.6 -10.1 -15.4 -30.6 -10.1 -15.4 -10.1 -10.1 -10.1 -10.6 -10.1 -1	174 175 175 177 178 179 180 181 182 183 184 185 186 187 188	111 102 117 132 118 129 115 122 107 113 136 123 114 106 106	119.5 119.5 119.5 119.5 119.5 118.5 118.3 118.3 118.3 118.3 118.3 118.3 118.3 117.8 117.8 117.8	-17.5 12.7 -12.5 12.7 -10.2 -3.7 -11.3 -5.2 17.9 -3.8 -16.7 -11.7

IN FORMATION SUPPLEMENTAL STATISTICAL 5-7 TABLE C0916-A: MATRIX ł CΕ VARIAN CE/ CO-VARIAN

VEL2
TIM2 -4.0042820 -1.042820 -1.042820 -1.5608203 1.5608203 5.887820 5.887820 5.887820 5.847800 -6.2708202
IN T2 -7.542E0 -7.542E0 -7.542E0 -7.622E0 -2.340E02 1.647E00 -2.340E02 1.447E00 -2.340E02 1.487E00 -2.340E02 1.487E00 -2.340E02 -2.45000 -2.450000 -2.450000 -2.45000 -2.45000 -2.45000 -2.4500000000000000000000000000000000000
DIA2
VEL1 1.165-11 1.065E-02 1.065E-02 4.535E02 -2.535E02 -2.53408E-05 -2.3408E-01 1.798E-03 -1.560E-03 -1.798E-05
TIM1 -7.128-10 -7.128-10 -7.3078-01 1.6538200 -2.5398203 1.6538200 1.6472800 1.659201 -1.4268-03
SKY -7.112E08 -7.112E08 -7.112E01 -1.065E00 1.625E002 3.3522E009 -7.6222E009 -8.2032E001 -8.2032E001 -8.2032E001
INT1 1.089E07 2.168E07 2.115E01 -7.315E01 -7.305E02 -1.382E02 -1.382E02 -1.042E01 -1.721E01
DIA1 7.252 7.252 1.089E 1.089E 1.089E 1.089E 1.145E 1.

CORRELATION MATRIX

VEL2 7.781E02 8.901E02 VEL2 0.912597 0.912597 0.912597 0.722321 0.722321 0.904073 -0.904073 -0.904073 -0.904002 TIM2 1.061E01 1.073E01 TIM2 -0.916912 -0.916912 0.442498 -0.7338199 -0.7338199 -0.916134 1.0010134 -0.999412 IN T2 1.368EQ0 5.914E 03 DERIVATIVES E07 DIA2 7.970E07 8.731E07 E PARTIAL VEL1 3.461E03 3.280E03 VEL1 0.557477 0.5574776 0.557867 0.235867 0.235867 0.235867 0.833559 0.704793 E 0Ł TIM1 3.944E01 4.040E01 TIMI -0.554781 -0.554781 -0.204687 -0.204687 -0.204687 -0.204687 -0.204687 -0.204687 -0.25478153 -0.252153 -0.722321 RANGES UMERICAL SKY 1.000E00 1.000E00 SKY -0.659525 -0.659525 1.000000 -0.204687 0.327547 0.327547 0.451776 Ξi IN T1 IN T1 0.999072 0.999725 0.9554781 0.5554781 0.8944776 0.912597 0.912597 DIA1 3.873E08 4.265E08 DIA1 1.00000 0.999072 0.585661 0.585661 0.585661 0.912874 0.915836 0.915395 MUMI XAM MUMI NIM DIAI INTI SKY VELI VELI DIA2 INT2 VEL2

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ZC0916-B: SUPPLEMENTAL STATISTICAL INFORMATION

VARIAN CE/ CO-VARIAN CE MATRIX

VEL2 5.96750 5.96750 5.96750 5.9166202 -0.
TIM2
IN T2
DIA2
VEL1 3.152E10 3.152FE01 -1.2290E-02 -1.27290E-02 -3.125E003 -3.125E000 -3.125E00 -3.125E0000 -3.125E0000 -3.125E00000 -3.125E000000000000000000000000000000000000
TIM1
SKY
INT1
DIA1

CORRELATION MATRIX

	DIAI	IT NI	SKY	TIMI	VEL1	DIA2	IN T2	TIM2	VEL2
DIAI	1.000000	066666.0	0.933433	-0.999828	0.999732	-0.999963	-0.999952	-0.999879	0.999831
ILNI	0666666°0	1,000000	0.932503	-0.9999806	0.999719	-0.999955	-0.999943	-0.999810	0.999758
SKY	0.933433	0.932503	1.000000	-0.939145	0.940034	-0.935804	-0.936137	-0.933788	0.932952
TIMI -	-0.999828	- 0.999806 -	0.939145	1.000000	-0.9999987	0.9999946	0.999952	0.999594	-0.999486
VELI	0.999732	0.999719	0.940034	-0.999987	1.000000	-0.999892	-0.999902	-0.999434	0.999307
DIA2	-0.999963	-0.999955 -	0.935804	0.999946	-0.999892	1.000000	0,999999	0 * 9976767	-0.999689
LI IZ	-0.999952	-0.999943 -	0.936137	0.999952	-0.999902	0.999999	1.000000	0.999754	-0.999674
TIM2 -	-0.9999879	-0.999810 -	0.933788	0.999594	-0.999434	0.999767	0.999754	1.000000	-0.999993
V ELZ	0.999831	0.999738	0.932952	-0.999486	0.999307	-0.999689	-0.999674	-0.999993	1.000000
			NUMERIC	AL RANGES (OF THE PART	IAL DERIVAT	IVES		
	DIA1	II NI	SKY	TIMI	VEL1	DIA2	IN T2	TIM2	VEL2
MAX IMUM	2.949E0	3 1.369EQ0	1.000E00	3 263E0	1 2.357E0	3 2.997E0	7 1.369EQ0	3.716E00	2.180E02
MUMI NIM	-3 .150E0	5.562E 03	1.000E00	-3.334E0	1 -2.073E0	3 -2.797E0	7 3.607E 0	13 -3.788E00	-2.009E02

The contention that the lack of any detectable spectral features from the spectroscopic component of 1 Gem B is due to a considerably lower intrinsic luminosity seems consistent with the V magnitude found for this component.

ZC1221 (9 Cancri)

The moderately bright (mV=6.24) M3-III star ZC1221 (9 Cancri) was occulted under favorable conditions on March 24, 1983. Table 5-9 gives a synopsis of the occultation observation. The RAWPLOT, Figure 5-14, shows the raw observational data. Data acquisition was halted approximately 800 milliseconds before the end of the data window.

A careful examination of the integration plot, Figure 5-15, indicates a slight downward deflection at 2300 milliseconds. This change in integrated level, due possibly to the presence of a second star, is rather subtle. In order to more easily see this level change, an integration plot approximately 1200 milliseconds in length, centered roughly at the suspected time of secondary disappearance and excluding the primary occultation, was plotted and is shown in Figure 5-16. From the cusp seen at 2280 milliseconds, it is evident that a secondary event did occur.

The best fit to the primary occultation event, shown graphically in Figure 5-17, places the relative time of geometrical occultation close to 3426 milliseconds. Thus, the fainter star disappeared 1146 milliseconds before the brighter one. The angular velocity of lunar limb passage was

TABLE 5-9 ZC1221: LUNAR OCCULTATION SUMMARY

STELLAR AND OBSERVING INFORMATION

Star: ZC1221, (9 Cancri, SAO 079940, DM +23 1887) RA: 080518 DEC: +224106 mV: 6.24 Sp: M3-III Filter: V Diaphragm: I Gain: C7+ Voltage: 1000

LUNAR INFORMATION

Surface Illumination:	72	percent
Elongation from Sun:	116	degrees
Altitude Above Horizon:	78	degrees
Lunar Limb Distance:	361038	kilometers
Predicted Shadow Velocity:	657	2 meters/sec.
Predicted Angular Rate:	0	.3756 arcsec/sec.

EVENT INFORMATION

Date: March 24,	1983	UT of Event:	00:44:56
USNO V/O Code:	17	HA of Event:	-114426
Position Angle:	110.4	Cusp Angle:	835
Contact Angle:	-11.0	Watts Angle:	97.4

MODEL PARAMETERS

Number of Data Points:	201
Number of Grid Points:	1600
Number of Spectral Regions:	53
Width of Spectral Regions:	50 Angstroms
Limb Darkening Coefficient:	0.5
Effective Stellar Temperature:	3300K

SOLUTIONS

Stellar Diameter (ams): 2.70 (2.41) Time: (relative to Bin 0): 3425.8 (0.6) Pre-Event Signal: 2931.6 (16.7) Background Sky Level: 1578.0 (18.5) Velocity (meters/sec.): 691.0 (16.3) Lunar Limb Slope (degrees): -0.47 (1.35) U.T. of Occultation: 00:44:55.877 (0.012)

PHOTOMETRIC NOISE INFORMATION

Sum-of-Squares of Residuals:	5663500
Sigma (Standard Error):	168.28
Normalized Standard Error:	0.12432
Photometric (S+N)/N Ratio:	9.044
(Change in Intensity)/Background:	0.85782
Change in Magnitude:	0.67251







Figure 5-16. Detailed integration plot of ZC1221 showing the secondary event.



Figure 5-17. FITPLOT of the occultation of 2C1221.

found to be 0.3881 seconds of arc per second. The projected separation between the two stars, therefore, was 0.44 seconds of arc.

The actual change in mean signal level due to the fainter star's disappearance, as determined from the 500-millisecond samples before and after the secondary event, was 29.5 counts. This is only 0.022 of the intensity of the brighter star. Hence, the fainter star is approximately 4 magnitudes fainter in V than the brighter star. While there is always the possibility of a second star coincidentally being in the field of view under study, in an 15.2 arcsecond diaphragm (such as was employed for this observation) the likelihood of such a coincidence is only one part in 10000. Thus, it is extremely likely that ZC1221 is a double star. with a previously undetected component having a V magnitude of approximately 10.4.

Since the secondary event was well separated in time from the primary (i.e., "wide" in the sense of an occultation binary), the disappearance of the fainter star did not enter into the solution of the primary event. The solution to the occultation intensity curve of ZC1221 yielded a sensible angular diameter of 2.70 milliseconds of arc. Figure 5-18 shows the sensitivity of the solution intensity curve to variations in the solution parameters. The diameter determination is most sensitive at the time of the first order fringe maximum, though still highly sensitive at the zeroith and first order minima. As can be seen in





Figure 5-18, the best fit to the observed data, in this region, is quite good.

Figure 5-19 shows the residuals of the fit to the 4096 milliseconds of the observation. The Gaussian nature of the residuals (with a slight Poisson tail) is typical of a good fit and a well behaved sky background. The mean deviation from the observed intensity (a measure of the overall observational noise) is only 5.8 percent.

The contiguous subset of observed intensity data, the computed intensity, and the resulting residuals are given in Table 5-10. The variance-covariance and correlation matrices of the adjusted solution parameters, and the dynamic ranges of the partial derivatives of the intensity curve are presented in Table 5-11.

The POWERPLOT, Figure 5-20, of this event clearly shows that the high frequency components of the noise in the raw data arise from scintillation and photon statistics and are not inherent in the occultation signal.

It is of interest to note if the diameter determined by the occultation observation is reasonable in light of what is known about this and similar stars. The distance of ZC1221, as taken from its parallax given by Schlesinger and Jenkins (1940), is roughly 250 parsecs. This is also indicated by an absolute magnitude of -0.7, assumed from its spectral type and luminosity class. Typical characteristics of giant stars are tabulated by Mihalas and Binney (1981). Working from these tables one would expect the linear diameter of an

TABLE 5-10 ZC1221: OBSERVATIONS, COMPUTED VALUES, AND RESIDUALS FROM BIN 3315 NUM OBS COMP NUM OBS COMP RESID NUM OBS COMP RESID RESID ---- -----134 1441 1601.6 -160 .6 135 136 137 1145 1599.9 -454.9 1598.4 -195.4 1599.9 -454.9 1598.4 -195.4 1597.0 -192.0 1595.8 -241.8 1594.7 -63.7 1403 1405 137 1405 138 1354 139 1531 140 1548 141 1645 142 1685 1593.7 -45.7 93.1 219.9 -84.3 $\begin{array}{c} 143 \\ 143 \\ 144 \\ 1506 \\ 145 \\ 1491 \\ 145 \\ 1491 \\ 147 \\ 1485 \\ 1485 \\ 148 \\ 1615 \\ 1574 \\ 151 \\ 152 \\ 1574 \\ 151 \\ 152 \\ 1574 \\ 151 \\ 152 \\ 1574 \\ 151 \\ 152 \\ 1574 \\ 151 \\ 152 \\ 1574 \\ 151 \\ 152 \\ 1574 \\ 152 \\ 1574 \\ 152 \\ 1574 \\ 151 \\ 152 \\ 1574 \\ 151 \\ 152 \\ 1574 \\ 153 \\ 1574 \\ 155 \\ 1574 \\ 155 \\ 1574 \\ 155 \\ 1574 \\ 155 \\ 1574$ 10 - 97.6 -198.0 -103.4 -2 -8 -12.9 159 -8 .8 .8 20 21 22 23 110.5 53.2 54.5 13.8 -197.0 -23.8 -86.6 $\begin{array}{c} 3162 \\ 2396 \\ 2395 \\ 23$ 24 25 26 27 .6 -80.0 19.6 52.8 -175.0 -70.8 28 162 163 1603 29 1636 97 3172 98 2833 ΞÓ 3101.6 70.4 2987.6 -154.6 164 1408 165 1512 166 1423 99 100 101 102 103 -159.6 167 1502 168 1508 -80.4 34 -74 -74.3 35 169 1830 36 1475 -107 ۰Ó $\begin{array}{c} 104 \\ 105 \\ 213 \\ 2105 \\ 2107 \\ 2105 \\ 2107 \\ 2108 \\ 2107 \\ 2108 \\ 2107 \\ 2108 \\ 2107 \\ 2108 \\ 2107 \\ 2108 \\ 2107 \\ 2108 \\ 2107 \\ 2108 \\ 2107 \\ 2108 \\ 2107 \\ 2108 \\ 2107 \\ 2108 \\ 2107 \\ 2108 \\ 2107 \\ 2108 \\ 2107 \\ 2108$ 1ŏ4 171 1728 1581.9 172 1587 1581.8 173 1755 1581.8 174 1610 1581.5 175 1735 1581.4 176 1751 1581.4 177 1577 1581.2 178 1655 1581.1 1728 146 . žá 5.2 39 173.4 40 28.5 41 •6 169.7 42 43 -4. 2834 2937.5 2723 2940.3 2927 2934.8 2989 2929.2 2631 2931.0 44 73.9 179 1736 1581.0 45 155.0 46 1446 180 181 1581.0 -135.0 1580.9 -124.9 -300.0 -122.1 52.0 113.0 421.5 -263.2 1456 2941.1 2951.0 2950.0 48 2819 182 1627 1580.8 46 49 221...0 90..7 -58..9 -133..6 -232.0 126..9 228..6 -69.1 -128.28 -69.1 -128.28 -30.88 -103.1 -79.7 -36.5 1580.7 1580.6 1580.6 51.3 183 1632 50 3063 184 1488 •6 29305.55 2917.2 2917.2 2953.2 2977.6 2953.3 2911.3 2971.3 2873.55 2873.55 2961.9 3015.55 3039.9 3029.9 18.4 51 1599 1658 185 1867 1887 190 191 192 193 194 195 196 197 198 2654 2967 3039 2862 3007 1580.5 -263.2 56.3 114.3 -91.2 29.4 241.4 -430.3 -48.3 1658 1580.5 1629 1580.4 1728 1580.4 1660 1580.3 1498 1580.3 1548 1580.2 1551 1580.2 1676 1580.1 1807 1580.1 534 555 56 48.6 147.6 79.7 -82.3 3007 3221 2523 2863 3093 -32.2 58 -29.2 59 -48.3 -48.3 215.5 419.5 268.1 105.1 227.5 164.1 60
 1676
 1380.1
 226.9

 1455
 1580.1
 226.9

 1455
 1580.0
 -125.0

 1455
 1579.9
 114.9

 1465
 1579.9
 35.1

 1465
 1579.9
 35.1

 1719
 1579.8
 139.2

 1775
 1579.8
 195.2
 1807 1455 1359 1465 1615 1719 32 93 129 130 131 132 133 62 3174 1602 1613.8 1580 1505 1526 1640 1613.8 1610.8 1608.1 1605.7 1603.5 3067 3243 3204 2892 64 65 199 66 -128.9

VARIANCE/CO-VARIANCE MATRIX -------

DIAM	PREI	POST	TIME	VELO
2.463E ⁻ 16	3.103E 08	-2.809E 08	-7.701E ⁻ 10	2.575E ⁻ 11
3.103E 08	2.781E02	8.816E00	-1.967E00	4.594E ⁰ 2
-2.809E 08	8.816E00	3.423E02	-2.411E00	5.674E ⁻ 02
-7.701E ⁻ 10	-1.967E00	-2.411E00	3.342E ⁻ 01	-7.973E 03
2.575E ⁻ 11	4.594E ⁻ 02	5.674E 02	-7.973E 03	2.643E 04

CORRELATION MATRIX ------

	DIAM	PREI	POST	TIME	VELO
DIAM	1.000000	0.797793	-0.744069	0.077028	-0.081285
PREI	0.797793	1.000000	-0.191055	-0.535414	0.531644
POST	-0.744069	-0.191055	1.000000	-0.721129	0.723939
TIME	0.077028	-0.535414	-0.721129	1.000000	-0.999988
VELO	-0.081285	0.531644	0.723939	-0.999988	1.000000

NUMERICAL RANGES OF THE PARTIAL DERIVATIVES

DIAM	2.730E09	-2.742E09	
PREI	1.358E00	1.324E_03	
POST	9.987E 01	-3.578E 01	
TIME	1.200E02	-1.243EU2	
VELO	4.841 E03	- 5.026E03	



M3-III star to be on the order of 50 solar diameters. At a distance of 250 parsecs this corresponds to an angular diameter of 1.9 milliseconds of arc. Hence, the observationally determined diameter seems entirely reasonable and in accord with the previously known stellar data.

A rather good internal precision of the time of geometrical occultation, with a formal error of only 0.58 milliseconds, was achieved in the reduction process. However, the radio signal carrying the WWUB time code was unusually noisy, leading to a large one sigma error in the Coordinated Universal Time of 0.012 seconds. The Coordinated Universal Time of geometrical occultation was determined to be 00:44:55.877 (+/- 0.012 seconds).

ZC1222

As indicated on the occultation summary for the 60 star ZC1222 (Table 5-12), this event occured only one half hour after the previously discussed occultation of 9 Cancri. The photometric conditions for these two events were nearly identical. This is evidenced by an intercomparison of the normalized standard errors, and the photometric (S+N)/N ratios of the two events (noting that ZC1222 is both a magnitude fainter and considerably bluer than ZC1221). The event was extremely well placed, being at a zenith distance of only 8 degrees.

The RAWPLOT of the event is shown in Figure 5-21. The integration plot over the entire data window (Figure 5-22), as well as an examination of other detailed integration plots

TABLE 5-12 ZC1222: LUNAR OCCULTATION SUMMARY

STELLAR AND OBSERVING INFORMATION -------Star: ZC1222, (SAO 079948, DM +22 1854) RA: 080535 DEC: +223024 mV: 7.22 Sp: G0 Filter: V Diaphragm: I Gain: C10 Voltage: 1200 LUNAR INFORMATION _____ Surface Illumination: 72 percent Elongation from Sun: 116 degrees Altitude Above Horizon: 82 degrees Lunar Limb Distance: 360913 kilometers Predicted Shadow Velocity: 336.9 meters/sec. Predicted Angular Rate: 0.2467 arcsec/sec. EVENT INFORMATION _____ Date: March 24, 1983 UT of Event: 01:14:46 USNO V/O Code: 26 HA of Event: -042030 Position Angle: 150.8 Cusp Angle: 425 Contact Angle: -49.6 Watts Angle: 137.8 MODEL PARAMETERS -----201 Number of Data Points: Number of Grid Points: 256 Number of Spectral Regions: 53 Width of Spectral Regions: 50 Angstroms Limb Darkening Coefficient: 0.5 Effective Stellar Temperature: 5900 K SOLUTIONS _____ Stellar Diameter (ams): Point Source Time: (relative to Bin 0): 2923.3 (2.1) Pre-Event Signal: 2117.4 (20.7) Background Sky Level: 1305.0 (44.8) Velocity (meters/sec.): 431.7 (12.5) Lunar Limb Slope (degrees): +19.34 (1.67) U.T. of Occultation: 01:14:46.177 (0.004) PHOTOMETRIC NOISE INFORMATION -------Sum-of-Squares of Residuals: 13498450 Sigma (Standard Error): 259.7927 Normalized Standard Error: 0.31979 Photometric (S+N)/N Ratio: 4.12699 (Change in Intensity)/Background: 0.62251 Change in Magnitude: 0.52547





of data subsets, show no compelling evidence of "wide" stellar duplicity.

The best model fit to the observations indicated that the star was a point source (or at least below the detection threshold of approximately 1 millisecond of arc). This was as expected, considering the spectral type and apparent V magnitude of the star. The graphic depiction of the fit is shown in Figure 5-23. The observational data, the solution curve intensities and the residuals are given in Table 5-13. The formal statistics of the solution parameters (variance-covariance, correlation matrices, and range of partial derivatives) can be found in Table 5-14. The PDPLOT, showing the sensitivity of the model intensity curve to the variation of the solution parameters, is presented as Figure 5-24.

The internal formal error of the time of geometrical occultation was not quite as good as one might have expected, if that expectation were based solely on the noise statistics of the observation. It should be noted that the determined L-rate for this event (431.7 meters/second) is 63 percent slower than the L-rate determined for the ZC1221 event (691.0 meters/second). Under identical photometric conditions this would widen the uncertainty in the time of geometrical occultation by the same amount. Fortunately, in the intervening half hour between these two events, the signal strength of the WWVB time code improved dramatically. The one sigma error of the WWVB time reference was a reasonable



TABLE 5-13 ZC1222: OBSERVATIONS, COMPUTED VALUES, AND RESIDUALS FROM BIN 2760

TABLE 5-14 ZC1222: SUPPLEMENTAL STATISTICAL INFORMATION

VARIAN CE/CO-VARIAN CE MATRIX

DIAM	PREI	POST	TIME	VELO
2.849E ⁻ 15	1.373E 07	-5.788E07	7.215E ⁻ 10	1.217E ⁻ 11
1.373E 07	4.349E02	7.193E01	-1.093E01	5.701E_02
-5.788E07	7.193E01	2.154E03	-4.181 E01	2.091E_01
7.215E ⁻ 10	-1.093E01	-4.181 E01	4.480E00	-2.315E_02
1.217E ⁻ 11	5.701E ⁰²	2.091E ⁻ 01	-2.315E ⁻ 02	1.587E 04

CORRELATION MATRIX

	DTAM	PREI	POST	TIME	VELO
DIAM	1.000000	0.288964	-0.968314	0.871718	-0.865589
PREI	0.288964	1.000000	-0.040942	-0.210200	0.221926
POST	-0.968314	-0.040942	1.000000	-0.965427	0.962118
TIME	0.871718	-0.210200	-0.965427	1.000000	-0.999923
VELO	-0.865589	0.221926	0.962118	-0.999923	1.000000

NUMERICAL RANGES OF THE PARTIAL DERIVATIVES

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	MAXIMUM	M IN IMUM
DIAM	8.475E08	-8.332E08
PREI	1.366E00	2.803E 02
POST	9.720E 01	-3.664E 01
TIME	3.576E01	-3.819E01
VELO	6.809E03	-7.093E03







Figure 5-25. NOISEPLOT of the occultation of ZC1222.

3.1 milliseconds. The Universal Time of geometrical occultation was found to be 01:14:50.125 (+/- 0.004 seconds).

The distribution function of the residuals over the 4096 milliseconds of the observation can be seen in Figure 5-25. As expected, this distribution function is quite similar in character to that seen for the ZC1221 occultation. The POWERPLOT for this event is presented as Figure 5-26.

X07589

The occultation of X07589 on April 18, 1983, was observed with the moon only 25 percent illuminated, at a solar elongation of 59 degrees (as noted in Table 5-15). The Earthshine on the shadowed portion of the lunar disk was quite bright, easily detectable to the unaided eye and, of course, photometrically. Examination of the 4096 millisecond plot of the raw data, presented in Figure 5-27, reveals unquestionably that the moon's limb entered the photometeric field of view approximately 2200 milliseconds from the beginning of the data window. For most occultation observations the increase in background signal level due to sunlight reflected off the lunar limb is such a small contributing source that it can be ignored. However, this is not the case for small lunar elongations, when Earthshine can be a major variable source of background light. This is evidenced, in this case, by the inverted nature of the integration plot for X07589 (Figure 5-27) prior to the time of occultation. Stellar duplicity would produce a cusp,



TABLE 5-15 X07589: LUNAR OCCULTATION SUMMARY

STELLAR AND OBSERVING INFORMATION _____ Star: X07589 (SAD 077552, DM +23 1042) RA: 054609 DEC: +230959 mV: 8.6 Sp: M0 Filter: V Diabhraom: I Gain: C10+ Voltage: 1200 LUNAR INFORMATION Surface Illumination: 25 percent Elonoation from Sun: 59 degrees 36 degrees Altitude Above Horizon: 36 degrees Lunar Limb Distance: 368613 kilometers Predicted Shadow Velocity: 843.5 meters/sec. Predicted Angular Rate: 0.4720 accsec/sec. EVENT INFORMATION USNO V/O Code: 26 HA of Event: 01:35:12 USNO V/O Code: 26 HA of Event: +601759 Position Angle: 95.7 Cusp Angle: 835 Contact Angle: -4.3 Watte April MODEL PARAMETERS _____ Number of Data Points: 201 Number of Grid Points: 256 53 Number of Spectral Regions: Width of Spectral Regions: 50 A Limb Darkening Coefficient: 0.5 50 Anostroms Effective Stellar Temperature: 5900K SOLUTIONS

502011040

Stellar Diameter (ams):	Point Source	
Time: (relative to Bin 0):	3983.3 (2.1)	
Pre-Event Signal:	1977.9 (34.8)	
Background Sky Level:	1129.0 (36.1)	
Velocity (meters/sec.):	1341.3 (95.8)	
Lunar Limb Slope (degrees):	-25.5 (6.5)	
U.T. of Occultation: 01	:35:12.172 (0.051)	

PHOTOMETRIC NOISE INFORMATION




deflecting the curve in a downward direction. The intrusion of increasing Earthshine midway through the data window caused a decrease in the integrated intensity (subtracted from the mean) for data taken before the moon's limb entered the diaphragm.

Fortunately, since the star was held close to the center of the field-of-view (as called for by standard observing procedures), the increase in background illumination was very close to linear over the short timescale of the occultation event. The lunar angular velocity for this event was 0.4720 seconds of arc per second. Thus, in the 200 milliseconds of data extracted for the DC solution, the lunar limb traversed an angular distance of only about 100 milliseconds of arc. This is only 1/150 of the diameter of the field-of-view given by the diaphragm employed in the observation. (See Table 2-3).

The assumption of linearity in the increase in background light is indeed proven correct by examination of the integration plot. The non-linear change in intensity (due to Earthshine) begins at roughly 2200 milliseconds. The slope of the descending branch of the integration plot changes here. The time at which the change in intensity due to contributed Earthlight is greatest, is at the minumum of the integration plot, at approximately 2700 milliseconds. While still increasing, the change in Earthlight is essentially linear by millisecond 3600, where the slope of

the integration plot remains constant until the decrease caused by the occultation itself sets in.

The result of a linear increase in background light causes an effective tilt in the observed occultation intensity curve. Rather than modeling the slope of this tilt as an additional free parameter in the DC solution, a simpler approach was taken. A linear least squares fit was done to both the pre-occultation and post-occultation observations, each containing 100 samples of data bordering the solution data subset. The characteristic slopes obtained were found to be in good agreement and were averaged. The occultation data were then "detilted" by the appropriate amount before being submitted to the DC solution process.

The only time this method of removing increasing moonlight from the occultation data set would not be valid are in the cases of near-grazing incidence, or if the star were near the eastern or western extreme of the field-of-view. In either of these cases, the increase of moonlight at the time of occultation would be non-linear.

In this case, within the linear portions of the integration plot examined in detail, there was no indication of a secondary stellar component.

The solution to the intensity curve is presented graphically on Figure 5-29. The geometry of this event led to a rather rapid L-rate, and as a result the intensity variation of the solution curve is a bit obscured in Figure 5-29. A detailed depiction of the region of interest





Figure 5-30. 100 millisecond detailed FITPLOT for X07589.

(covering only 100 milliseconds of the observation) is presented in Figure 5-30.

The diameter of X07589 was below the detection threshold, and hence, was indistinguishable from a point source. The PDPLOT for the solution curve is presented as Figure 5-31. The observational data subset used in the solution determination and the computed intensity values, and their residuals are listed in Table 5-16. The variance-covariance and correlation matrices, and the ranges of the partial derivatives are given on Table 5-17.

No attempt was made to fit the non-linear increase in the sky background as the moon entered the diaphragm. Hence, the noise statistics for this event, as well as the distribution function (Figure 5-32) of the noise, reflect only the first 1600 milliseconds of pre-occultation data and detilted occultation data. Even the early part of the pre-occultation record contains an increase in the background level. This is evidenced by the skewed nature of the NOISEPLOT, Figure 5-32.

The power spectra of the event, the star-plus-sky signal, and the solution intensity curve are shown on Figure 5-33. The observed power spectra (occultation and star-plus-sky) reflect the actual observation with the increasing sky background included.

A formal internal error for the time of geometrical occultation of only 1.1 milliseconds was obtained for this event. As had been seen from previous experience, the signal



Figure 5-31. PDPLOT of the occultation of X07589.



Figure 5-32. NOISEPLOT of the first 1600 milliseconds of data taken for the occultation observation of X07589.

TABLE 5-16 X07589: OBSERVATIONS, COMPUTED VALUES, AND RESIDUALS FROM BIN 3878

NTH OBS COMP RESID NUM OBS COMP RESID Add Add <th>===:</th> <th></th> <th></th> <th></th> <th></th> <th></th> <th></th> <th></th> <th>======</th> <th></th> <th></th> <th>DECID</th>	===:								======			DECID
$\begin{array}{cccccccccccccccccccccccccccccccccccc$	NUM	OBS	COMP	RESID	NUM	OBS	COMP	RESID	N UM	OBS	COMP	RESID
55 1944 1980.8 - 36.4 122 691 1136.7 - 446.2 189 775 1129.4 - 354.8	=N1-012345678901234567890123456789012345678901234567890012345678900123456789000000000000000000000000000000000000	= 0 = 0 = 0 = 0 = 0 = 0 = 0 = 0 = 0 = 0	COMP 	$\begin{array}{c} \text{Resc} 1 \\ \text{Resc} 1 \\$	$ \begin{array}{l} = = = = = \\ & - u - u - u \\ = u - u - u - u \\ = u - u - u \\ = u - u - u \\ = $	$\begin{array}{c} = & = & = & = & = & = & = & = & = & = $	$\begin{array}{c} \hline & \hline \\ \hline \\$	$\begin{array}{c} \qquad \qquad$	$ \begin{array}{c} = = = \\ = = \\ = = \\ = \\ 1 \\ 1 \\ 1 \\ 3 \\ 5 \\ 1 \\ 1 \\ 3 \\ 5 \\ 1 \\ 1 \\ 1 \\ 3 \\ 5 \\ 1 \\ 1 \\ 1 \\ 1 \\ 4 \\ 3 \\ 1 \\ 1 \\ 4 \\ 1 \\ 4 \\ 4 \\ 3 \\ 1 \\ 1 \\ 1 \\ 1 \\ 1 \\ 1 \\ 1 \\ 1 \\ 1$	$\begin{array}{c} = & = & = & = \\ = & = & = & = & \\ = & = &$	COMP COMP	$\begin{array}{c} \textbf{RESID} \\ \textbf{RESID } \\ RES$
J4 ZJ17 17/7*/ JJ7*Z 1Z1 11J4 113/*0 10*J 100 73J 1127*4 -174*3	489 512 523 54	1729 2007 1365 2034 2103 1633 2319	1980.5 1980.2 1980.1 1980.6 1980.0 1980.8 1979.7	-251.8 27.0 -615.4 53.7 122.9 -347.5 339.2	115 116 117 118 119 120 121	860 824 1400 1234 1520 1078 1154	1148.6 1145.7 1143.3 1141.4 1139.9 1138.7 1137.6	-288.8 -321.4 256.5 92.9 380.0 -60.2 16.3	182 183 184 185 186 187 188	923 809 465 1592 1127 1254 935	1129.4 1129.4 1129.4 1129.4 1129.4 1129.4 1129.4 1129.4 1129.4	-206.6 -320.0 -664.5 463.0 -2.4 124.1 -194.3
	63 64 65 66	1505 2050 2108 2759	1976.8 1982.4 1982.4 1982.5	-472.0 67.9 125.4 783.9	130 131 132 133	947 2210 791 1038	1132.7 1132.4 1132.2 1131.9	-185.8 1077.9 -341.3 -93.5	197 198 199 200	1221 1578 1451 923	1129.3 1129.3 1129.3 1129.3	91.6 449.1 321.7 -206.8

TABLE 5-17

X07589: SUPPLEMENTAL STATISTICAL INFORMATION

VARIANCE/CO-VARIANCE MATRIX

DIAM	PREI	POST	TIME	VELO
1.143E ⁻ 16	-4.465E 08	-5.040E ⁻ 09	1.273E 09	-1.521E ⁻ 10
-4.465E ⁻ 08	1.212E03	2.450E01	-5.787E00	4.580E ⁻ 01
-5.040E_09	2.450E01	1.303E03	-5.392E00	3.745E ⁻ 01
1.273E_09	-5.787E00	-5.392E00	1.196E00	-8.278E 02
-1.521E ⁻ 10	4.580E ⁻ 01	3.745E ⁻ 01	-8.278E 02	9.180E 03

CORRELATION MATRIX

	DIAM	PREI	POST	TIME	VELO
	~ ~				
DIAM	1.000000	-0.994887	0.113292	0.727451	-0.778838
PREI	-0.994887	1.000000	-0.210486	-0.654490	0.711548
POST	0.113292	-0.210486	1.000000	-0.587384	0.523549
TIME	0.727451	-0.654490	-0.587384	1.000000	-0.996928
VELO	-0.778838	0.711548	0.523549	-0.996928	1.000000

NUMERICAL RANGES OF THE PARTIAL DERIVATIVES

	MAXIMUM	MIN IMUM
DIAM	8.781E08	-8.462E08
PREI	1.366E00	3.206E 04
POST	9.997E ⁰ 1	-3.656E 01
TIME	1.469E02	-1.496E02
VELO	1.601E03	-1.799E03

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strength of the WWVB radio signal during and shortly after evening twilight tends to be quite low. As a result, no WWVB time code was available at the time of the event. Fortunately, the signal strength had improved sufficiently thirty seven minutes later to reliably detect the digital time code. The crystal controlled clock circuit of the SPICA-IV/LODAS system has been found to have a cumulative timing error of at most two seconds per day. Taking this as a worst case condition, the time of the event reckoned from the SPICA-IV/LODAS clock and synchronized to WWVB a half hour after occultation would have a one sigma error of approximately 0.025 seconds. The Universal Time of geometrical occultation was 01:35:12.17 (+/- 0.03 seconds). X07598

Thirty five minutes after the previously discussed occultation of X07589, the star X07598 was occulted by the moon. The RAWPLOT of the digital photoelectric record obtained is presented in Figure 5-34. The integration plot, Figure 5-35 leads to no indication of "wide" stellar duplicity.

The near grazing geometry of this event (see Table 5-18) led to a predicted R-rate of only 0.0855 seconds of arc per second. Hence, the time-scale of the diffraction phenomena for this event was roughly four times slower than a more typical event. While this was known beforehand, a solution was attempted by fitting only 200 milliseconds of data. Even though a solution was found, the resulting model curve did

TABLE 5-18 X07598: LUNAR OCCULTATION SUMMARY

STELLAR AND OBSERVING INFORMATION Star: X07598 (SAO 077559, DM +22 1032) RA: 054627 DEC: +225501 mV: 7.5 Sp: K0 Filter: V Diaphram: I Gain: C8+ Voltage: 1200 LUNAR INFORMATION Surface Illumination: 25 percent Elongation from Sun: 59 degrees Altitude Above Horizon: 28 degrees Lunar Limb Distance: 369285 kilometers 59 degrees 28 degrees Predicted Shadow Velocity: 153.1 meters/sec. Predicted Angular Rate: 0.0855 arcsec/s 0.0855 arcsec/sec. EVENT INFORMATION _____ Date: April 18, 1983 UT of Event: 02:12:03 USN0 U/O Code: 47 HA of Event: +672638 Position Angle: 172.1 Cusp Angle: 75 Contact Angle: -80.3 Watts Angle: 172.3 MODEL PARAMETERS _____ Number of Data Points: 401 Number of Grid Points: 256 Number of Spectral Regions: 53 Width of Spectral Regions: 50 Angstroms Limb Darkening Coefficient: 0.5 Effective Stellar Temperature: 5100 K SOLUTIONS _____ Stellar Diameter (ams): 5.45 (2.04) Time: (relative to Bin 0): 2645.1 (3.8) Pre-Event Signal: 2253.0 (21.9) Background Sky Level: 1139.3 (38.9) Velocity (meters/sec.) 136.3 (4.4) Lunar Limb Slope (degrees): +13.5 (1.6) U.T. of Occultation: 02:12:00.519 (0.004) PHOTOMETRIC NOISE INFORMATION ------Sum-of-Squares of Residuals: 47842640 Sigma (Standard Error): 345.8419 Normalized Standard Error: 0.31053 Photometric (S+N)/N Ratio: 4.22027 (Change in Intensity)/Background: 0.97752 Change in Magnitude: Change in Magnitude: 0.74030





not span a large enough portion of the intensity curve to be warranted as significant. The observation was re-reduced considering 400 milliseconds of data centered on the estimated time of geometrical occultation.

The solution obtained from the 400 millisecond data set is presented graphically in the FITPLOT, Figure 5-36. Except for the anomalous depression in the zero order diffraction maximum, the fit is quite good. This depression is not too troublesome, remembering that the time-scale is elongated by a factor of roughly four. It is unfortunate that variation in atmospheric transparency took its toll at this point, but it was not unduly critical in the determination of the solution.

The observational data used in the solution, the computed intensities and their residuals are given in Table 5-19. The usual supplementary statistical information is listed in Table 5-20.

An angular diameter for this K0 (mV=7.5) star of 5.45 (+/- 2.04) milliseconds of arc was determined. There is no apparent reason to doubt the validity of the solution in terms of numerical or computational problems as evidenced by either the PDPLOT (Figure 5-37), the distribution function of the residuals as given in the NOISEPLOT (Figure 5-38), or the comparative power spectra (Figure 5-39). A further attempt was made to check the validity of the numerical solution in this case by attempting a refit to the data by a four parameter model, holding the stellar diameter fixed as a



TABLE 5-19 X07598: OBSERVATIONS, COMPUTED VALUES, AND RESIDUALS FROM BIN 2350 NUM OBS COMP RESID NUM OBS COMP RESID NUM OBS COMP RESID ------- -------- ---- --------- ---- -----0 2240 2449 2383 2355 2305 2066 2356 2143 2263 215 2649 1885 1885 56 89 12 3 1885 2146 2159 2632 îй 15 2575 2451 18 19 20 99 2281.6 -182.6 -182.0 739.5 286.2 1137.5 709.6 3023 2283.5 2571 3423 2284.8 2285.5 2995 24 25 26 20 95 2107 2284.7 -189.7 2283.3 -176.3 2281.2 453.8 2281.2 453.8 2278.3 -766.3 2735 1512 28 2919 2772 2280 2274.9 644.1 501.1 2266.5 2280 2407 2240 2825 2407 1920 1584 1743 2261.6 2256.5 2251.2 2245.9 145.4 -16.5 573.8 161.1 34 35 2243.9 101.1 2240.6 -320.6 2235.6 -651.6 2230.8 -487.8 2226.6 -7.6 36 2230.86 22226.66 22222.88 2217.55 2216.00 22215.88 2217.02 2215.88 2217.02 22215.88 22231.88 22231.88 22231.88 22231.88 22231.88 22231.88 22231.88 22231.88 22231.88 22231.88 22231.88 22231.88 22231.88 22231.88 22231.88 22231.88 22233.88 22333.888 22333.88 22333.88 22333.88 22333.88 22333.88 22333.88 22333.88 38 2219 1506 1833 2015 2445 2195 2195 2655 2105 1864 -716.8 -386.8 **4**0 41 -202.5 -202.5 229.0 -20.4 -20.8 438.0 -114.2 -358.4 -562.4 43 44 45 46 4Ť -358.4 -562.4 -681.2 234.2 287.0 -213.7 -33.9 -73.4 48 1664 49 1550 2471 2530 5Ó 51 2243.0 185 2605 186 2680 187 2475 188 2567 189 2352 190 2528 191 2421 192 2512 193 2459 194 2298 195 2032 196 190 196 190 197 2368 198 1963 $\begin{array}{c} 2235018 \\ 224762.67 \\ 224762.67 \\ 22502.47 \\ 22502.47 \\ 22502.47 \\ 1-171.1 \\ 22502.47 \\ 1-171.1 \\ 22502.47 \\ 1-171.1 \\ 22502.47 \\ 1-171.1 \\ 1-17.1 \\$ 52 53 54 55 56 57 58 59 60 64 -655.4 -286.2 -178.0 65 66 199 2338 2624.2 200 2451 2629.0

TABLE 5-19. CONTINUED.

NUM	OBS	COMP	RESTD	NUM	OBS	COMP	RESID	NUM	OBS	COMP	RESID
201		2622 0	-200 8	260	1771	1700 5	70 5		1206	1226 0	50 1
202	2079	2635.6	-556.6	269	1523	1686.6	-163.6	336	1225	1234.5	-9.5
203	2116	2637.4	-521.4	270	1641	1673.0	-32.0	337	1297	1232.3	64.7
205	2017	2637.9	-620.9	272	1239	1646.5	-407.5	339	1310	1227.9	82.1
206	2465	2636.7	262.6	273	1574	1621.1	-47.1	340	886	1223.8	-337.8
208	3099	2631.3	467.7	275	1073	1608.8	-535.8	342	983	1221.8	-238.8
210	2739	2622.0	117.0	277	1551	1585.0	-34.0	344	1860	1218.1	-358.1
212	2650	2616.1	-205.2	278	1519	15/3.5	-388.5	345	1095	1216.2	-119.5
213	2555	2601.5	-46.5	280 281	1359	1551.3	-192.3	347 348	1087	1212.8	-125.8
215	3363	2583.6	779.4	282	1263	1530.1	-267.1	349	1121	1209.5	-88.5
217	3221	2562.6	658.4	284	1464	1509.9	-45.9	351	1451	1206.4	244.6
218	2885	2538.8	333.9 757.2	285	1695	1490.6	204.4	352	1367	1204.9	163.5
220	2889	2525.9	363.1	287	1450	1481.4	-31.4	354	1294	1202.1	91.9
222	2619	2498.3	120.7	289	1380	1463.6	-83.6	356	1768	1199.4	568.6
224	3487	2463.7	1018.4	290	1243	1455.0	-203.6	358	1534	1196.8	337.2
225 226	3583 3267	2452.9	1130.1 830.1	292 293	1420	1438.5	-18.5	359 360	1440 1207	1195.6	244.4
227	2691	2420.4	270.6	294	1603	1422.8	180.2	361	1405	1193.2	211.8
229	2531	2386.3	144.7	296	1151	1408.0	-257.0	363	1127	1191.0	-64.0
230	2821	2350.0	470.0	297	995	1393.9	-398.9	364	1251	1189.9	62.1
232 233	2802	2332.9	469.1	299 300	1022	1387.1	-365.1	366	1457	1187.9	269.1
234	23 93	22 96 .2	96.8	301	1353	1374.2	-21.2	368	1269	1185.9	83.1
236	2407	2258.8	148.2	303	1476	1361.9	114.1	370	1224	1184.1	39.9
237	2361 2151	2239.9	121.1	304 305	1103 1208	1356.1	-253.1	371 372	1548 1258	1183.2	364.8
239	2595	2201.8	393.2	306	1949	1344.8	604.2	373	1002	1181.5	-179.5
241	2153	2163.7	-10.7	308	1767	1334.1	432.9	375	1348	1179.9	168.1
243	1680	2125.5	-445.5	310	930	1324.0	-394.0	377	1645	1178.3	466.7
244	1903	2087.6	-184.6	311	1082	1319.2	-237.2	378	1152 817	1177.6	-25.6
246 247	1938	2068.7	-130.7	313	1243	1310.0	-67.0	380	1111	1176.2	-65.2
248	2104	2031.3	72.7	315	1224	1301.3	-77.3	382	1431	1174.8	256.2
250	1783	1994.5	-211.5	317	1061	1293.0	-232.0	384	1471	1173.5	297.5
251	1015	1976.3	-361.3	318	1178	1289.1	-111.1	385	1277	1172.9	104.1
253	1087	1940.4	-853.4	320	1487	1281.5	205.5	387	1759	1171.7	587.3
255	1404	1905.3	-501.3	322	1169	1274.4	-105.4	389	1600	1170.5	429.5
257	1797	1871.1	-74.1	324	1334	1267.7	66.3	390	1415	1169.4	245.6
258	1955	1854.3	-260.8	325 326	1112	1264.4	-152.4	3 92 3 93	1336 1124	1168.9	167.1
260	1737	1821.5	-84.5	327	1007	1258.3	-251.3	394	973	1167.9	-194.9
262	1857	1789.7	67.3	329	1920	1252.5	667.5	3 96	1591	1166.9	-575.9
264	1222	1758.9	-536.9	331	1458	1247.0	211.0	3 98	1102	1166.0	-64.0
265	1935	1729.2	205.8	332 333	15/6	1244.3	331.7 169.2	399 400	1392 1157	1165.5	226.5
267	1705	1714.7	-9.7	334	1307	1239.3	67.7				

TABLE 5-20 X07598: SUPPLEMENTAL STATISTICAL INFORMATION

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VARIANCE/CO-VARIANCE MATRIX

DIAM	PREI	POST	TIME	VELO
1.447E ⁻ 16	3.947E 08	-9.496E ⁻ 08	-1.637E ⁻ 09	3.319E ⁻ 12
3.947E 08	4.785E02	8.940 E01	-2.714E01	2.714E 02
-9.496E_08	8.940E01	1.517E03	-6.788E01	6.483E-02
-1.637E_09	-2.714E01	-6.788E01	1.463E01	-1.444E02
3.319E ⁻ 12	2.714E ⁰ 2	6.483E ⁻ 02	-1.444E ⁰ 2	1.891E-05

CORRELATION MATRIX

	DIAM	PREI	POST	TIME	VELO
DIAM	1.000000	0.394094	-0.908842	0.681148	-0.667250
PREI	0.394094	1.000000	0.024095	-0.389071	0.405659
POST	-0.908842	0.024095	1.000000	-0.921565	0.914093
TIME	0.681148	-0.389071	-0.921565	1.000000	-0.999822
VELO	-0.667250	0.405659	0.914093	-0.999822	1.000000

NUMERICAL RANGES OF THE PARTIAL DERIVATIVES

	MAXIMUM	MIN IMUM
DIAM	3.271E09	-3.408E09
PREI	1.346E00	2.314E 02
POST	9.769E ⁰¹	-3,458E01
TIME	1.910E01	-1.897E01
VELO	1.805E04	-1.822E04



Figure 5-37. PDPLOT of the occultation of X07598.



Figure 5-38. NOISEPLOT of the occultation of X07598.



point source. No convergent solution could be forced on a point source model. The determined angular diameter does indeed seem to be real.

Because of the relatively slow disappearance (due to the small R-rate) the internal formal error in the time of geometrical occultation was not very well determined (+/- 4.4 milliseconds). The Coordinated Universal Time of geometrical occultation was found to be 02:12:04.713 (+/- 0.0049 seconds).

<u>X13534</u>

The occultation of the KO star X13534 on April 21, 1983 U. T., (as noted in the occultation summary, Table 5-21) was the first of two events observed on that night. Interestingly enough, these two events both gave rise to the discovery of previously unknown "close" companions. Other than the recognition of 1 Geminorum A1 (2C0916-A1), these two were the only such discoveries made in the reduction and analysis of the twenty-two observations where "close" duplicity could have been uncovered.

Even from the plot of the raw occultation data (Figure 5-40), one can note a diminution of the positive-going signal spiking (due to scintillation noise) immediately before occultation. While this is only marginally noticeable on the presented RAWPLOT, this was seen quite easily when examined at the four-times greater resolution obtained by using a high resolution graphics display terminal. This diminution in the noise spiking (as

TABLE 5-21 X13534: LUNAR OCCULTATION SUMMARY

_____ STELLAR AND OBSERVING INFORMATION Star: X13534. (SAD 080497, DM +21 1939) RA: 085323 DEC: +203905 mV: 8.4 Sp: K0 Filter: V Diaphragm: I Gain: C8 Voltage: 1200 LUNAR INFORMATION _____ Surface Illumination: 58 percent 100 degrees Elongation from Sun: Altitude Above Horizon: 58 degrees Lunar Limb Distance: 364349 kilometers Predicted Shadow Velocity: 661.7 meters/sec. Predicted Angular Rate: 0.3746 arcsec/sec. EVENT INFORMATION ______ UT of Event: 02:46:31 Date: April 21, 1983 USNO V/O Code: 15 HA of Event: +342042 Cusp Angle: 53S Position Angle: 143.4 Contact Angle: -30.0 Watts Angle: 126.7 TWO-STAR MODEL PARAMETERS Number of Data Points: 201 256, 256 Number of Grid Points: Number of Spectral Regions: 53 Width of Spectral Regions: 50 Angstroms Limb Darkening Coefficients: 0.5, 0.05 Effective Stellar Temperatures: 5100 K, 5100 K SOLUTIONS ______ Sky Level: 1336.5 (39.2) STAR 2 STAR 1 Stellar Diameter (ams): Point Source Point Source Time (relative to Bin 0): 2389.2 (2.2) 2372.3 (7.2) Stellar Intensity: 649.5 (100.2) 229.9 (98.7) Velocity (meters/sec.): 666.9 (49.4) 559.0 (116.3) Lunar Limb Slope (degrees) -3.58 (4.28) -16.3 (10.1)

U.T. of Occultation: 02:46:29.632 (0.003) :29.594 (0.008)

Table 5-21. Continued.

TWO-STAR DERIVED QUANTITIES

Temporal Separation (milliseconds): 16.9 (7.6) Intensity Weighted Mean L-Rate (meters/sec): 634.7 (48.4) Intensity Weighted Mean R-Rate (arcsec/sec): 0.359 (0.027) Projected Spatial Separation. Based on Predicted R-Rate: 6.32 arc-ms Based on Weighted Mean R-Rate: 6.07 (2.74) arc-ms. Brightness Ratio (Brighter/Fainter): 2.82 (1.289) Magnitude Difference: 1.13 (+0.66, -0.41) mV of Star 1: 8.73 (0.14) mV of Star 2: 9.86 (0.37)

PHOTOMETRIC NOISE INFORMATION

Sum-of-Squares of Residuals:	17673786
Sigma (Standard Error):	297.269
Normalized Standard Error:	0.33803
Photometric (S+N)/N Ratio:	3.9583
(Change in Intensity)/Background:	0.6580
Change in Magnitude:	0.5490

-



it so appears when plotted at this scale) was actually due to a reduction in the zero order diffraction fringe maximum brought about by a second star.

The presence of a second star was uncovered by the DC2 fitting process. No "wide" stellar components are indicated by the integration plot, Figure 5-41. An initial attempt to fit the observational data to a one star model failed, and the divergence of the fitting procedure indicated that a two-star model was the appropriate one to use. The FITPLOT of the two-star solution is presented as Figure 5-42. The times of geometrical occultation of the individual components are noted by the dashed vertical lines. The arbitrary zero point reference on the linear distance scale coincides with the point of geometrical occultation of the brighter star.

Neither of the two stars had resolvable stellar diameters, and hence, led to point source solutions. The projected spatial separation of the components was found to be 6.07 (+/- 2.74) milliseconds of arc based on an intensity weighted mean of the determined R-rates of 0.3593 (+/- 0.0274) seconds of arc per second.

The fainter of the two stars underwent disappearance first. The brightness ratio of the components (brighter/fainter) was 2.82 (+/- 1.29). Hence, the difference in stellar magnitudes between the components is 1.13 (+0.66, -0.41). If the systemic apparent V magnitude is taken as 8.4 (as given by the DM catalog), the individual



components would have corresponding V magnitudes of 8.73 (+/- 0.13), and 9.86 (+/- 0.37).

The observations, extracted from the raw data used in the DC2 fit, the computed intensities and the resulting residuals are given in Table 5-22. The sensitivity of the determined two-star intensity curve to variations in the solution parameters is presented in Figure 5-43. A small amount of numerical "noise" is seen in the early part of the partial derivative curves for the velocity parameters. However, the noise amplitude is quite small, and as can be seen, is not in the region of the solution where such noise would be important in the parameter determinations. The numerical ranges of the partial derivatives, and supplemental statistical information are listed in Table 5-23.

The 6.5 percent shift of the mean of the residual amplitude distribution (Figure 5-44), over the 4096 milliseconds of the acquired data, was due to a small increase in background illumination over the event window. This can be seen if the RAWPLOT is examined carefully. The same gradual increase gives rise to the obvious Poisson tail in the distribution function. The slope of the increase is sufficiently small to be negligible on the time-scale of the data extracted for the DC2 solution (less than 1/20 of the entire data window).

The POWERPLOT for this occultation event is presented as Figure 5-45. The Coordinated Universal Times of geometrical occultation for the two component stars were found to be

TABLE 5-22X13534: OBSERVATIONS, COMPUTED VALUES, AND RESIDUALSFROM BIN 2250

NUM	OBS	COMP	RESID	NUM	OBS	COMP	RESID	NUM	OBS	COMP	RESID
		COMP 2218,7, 2218,4,1 2218,1 2217,	$ \begin{array}{c} \text{RESID} & -7.6 \\ \text{RESID} & -7.4 \\ RE$		$\begin{array}{c} 0 \text{ BS} & \\ -21999377\\ 2211605\\ 2431937\\ 2212022373\\ 220526933\\ 220526933\\ 220526933\\ 220526933\\ 220526933\\ 220526933\\ 220526933\\ 220526632\\ 220526632\\ 22056$	CMPF	$ \begin{array}{c} \text{RESTID} \\ \text{RESTID} \\ \text{38}, \text{7}, \text{7}, \text{234}, \text{9}, \text{7}, \text{234}, \text{1}, \text{7}, \text{234}, \text{1}, \text{7}, \text{234}, \text{1}, \text{1}, \text{2}, \text$	$\begin{array}{c} {\tt N} {\tt U} {\tt U} {\tt U} {\tt U} {\tt I} {\tt I}$	$\begin{array}{c} 0BS\\ -1135513774\\ 1135513774\\ 1135513774\\ 113232543\\ 1132323\\ 1132323\\ 1132323\\ 1132323\\ 1132323\\ 1132323\\ 1132323\\ 1132323\\ 1132323\\ 1132333\\ 1132333\\ 1132333\\ 1132333\\ 1132333\\ 1132333\\ 1132333\\ 11333333\\ 113333333\\ 11333333\\ 11333333\\ 11333333\\ 11333333\\ 11333333\\ 113333333\\ 113333333\\ 113333333\\ 113333333\\ 1133333333$	$\begin{array}{c} \text{COMP}\\ COM$	$\begin{array}{c} \textbf{R}=\textbf{S}=\textbf{ID}\\ \textbf{R}=\textbf{S}=\textbf{ID}\\ \textbf{R}=\textbf{S}=\textbf{ID}\\ \textbf{R}=\textbf{S}=\textbf{ID}\\ \textbf{R}=\textbf{ID}\\ \textbf{R}=\textbf{ID}\\\ \textbf{R}=\textbf{ID}\textbf$
40 41 42 43 44 45 46 47 48 49 50	2642 2008 2007 2116 2006 2391 2791 2168 1725 2290 2216	2221.2 2219.0 2216.0 2215.1 2217.1 2219.6 2221.4 2222.7 2222.7 2222.7 2222.7 2222.4 2213.4	420.8 -211.0 -209.0 -99.1 -211.1 171.4 569.6 -54.7 -497.7 70.6 2.6	107 108 109 110 111 112 113 114 115 116 117	1986 2292 1803 2927 1984 1558 1811 1783 2271 2367 2439	2181.7 2117.1 2067.3 2038.1 2031.7 2046.6 2078.5 2121.4 2168.7 2214.1 2252.2	-195.7 174.9 -264.3 888.9 -47.7 -488.6 -267.5 -338.4 102.9 186.8	174 175 176 177 178 179 180 181 182 183 183	1151 1156 1733 1433 1395 1870 1176 1071 987 1488 1178	1343.6 1342.9 1342.6 1342.3 1342.0 1341.8 1341.8 1341.6 1341.4 1341.2	-192.6 -187.3 390.1 90.4 52.7 528.0 -165.8 -270.6 146.8 -354.6 146.8
51 5235555555555555555555555555555555555	1951 1616 1811 2099 2232 2951 2403 2361 2887 2119 2338 2091 2502	2209.6 2212.4 2219.8 2225.9 2227.5 2226.2 2224.7 2221.4 2214.3 2205.1 2200.7 2200.7 2206.2 2219.4	-258.6 -596.4 -408.8 -126.9 724.8 178.3 139.6 672.7 -86.1 137.3 -115.2 282.6	118 119 120 121 122 123 124 125 126 127 128 129 130	2136 2274 2799 2093 1893 1804 2483 1612 2328 22328 22307 2165 2047	2279.0 2292.7 2290.7 2275.1 2246.6 2207.3 2159.5 2105.6 2047.9 1988.5 1929.3 1871.6 1816.5	-143.0 -18.13 -182.1 -353.6 -403.33 322.54 -435.9 349.55 277.42 293.55	1856 1867 1889 1901 1912 1934 1954 1956 197	1171 1431 1126 1234 1979 1711 1720 16963 1179 1077 1478 1023	1340.8 1340.7 1340.5 1340.2 1340.2 1340.1 1340.0 1339.8 1339.7 1339.6 1339.5 1339.4 1339.3	-169.8 90.3 -214.5 -106.8 370.9 380.0 359.2 123.3 -160.6 -262.6 -316.3
64 65 66	2487 2279 2867	2232.1 2238.0 2237.0	254.9 41.0 630.0	131 132 133	1587 1720 1795	1764.9 1717.0 1673.3	-177.9 3.0 121.7	1 98 1 99 200	1123 1255 1039	1339.2 1339.2 1339.1	-216.2 -84.2 -300.1







NOISEPLOT of the occultation of X13534.

IN FORMATION TABLE 5-23 SUPPLEMENTAL STATISTICAL •• K13534

IAN CE MATRIX VARIANCE/CO-VARIANCE

VEL2
TIM2
IN T2
DIA2
VEL1
TIM1
sky
IN T1
DIA1

CORRELATION MATRIX

VEL2 VEL2 0.988643 0.988643 0.970551 0.970557 0.971828 -0.991806 -0.991806 -0.991806 VEL2 TIM2 ------1.669E01 -1.781E01 TIM2 -0.989304 -0.989303 0.970768 -0.970768 -0.990968 0.417146 0.912630 1.092630 IN T2 -0.986257 -0.986257 0.986257 0.9866257 0.984625 0.984625 0.984625 0.984625 0.984625 0.984625 0.991806 -0.991806 DERIVATIVES THE PARTIAL VEL1 2.712E03 2.837E03 0F TIM1 5.625E01 5.948E01 TIMI TIMI -0.944109 -0.067580 -0.909000 -0.998766 0.2844568 0.984658 0.984653 0.970768 RANGES UMERICAL SKY 1.000E00 1.000E00 Ξi IN T1 0.973216 0.9747000000 0.944109 0.944109 0.948409 0.988620 0.988643 DIAI 1-000000 0.973216 0.973216 0.973216 0.936817 0.956013 -0.966973 0.986958 DIAI INTI SKY TIMI VELI DIA2 DIA2 TIM2 TIM2

1.224E03

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02:46:29.594 (+/- 0.008 seconds) and 02:46:29.632 (+/- 0.003 seconds), for the fainter and brighter components respectively.

X13607

The second occultation observed on the night of April 21, 1983, that of the F5 star X13607, occurred approximately one hour and forty eight minutes after the X13534 event. Because of the relatively early spectral type and moderate faintness (mV=8.2) of this star, a Johnson B filter had initially been chosen for the observation. However, a check on the sky brightness several minutes before the predicted time of disappearance showed that a significantly better StN/N ratio could be obtained with a Johnson V filter. The event, as indicated in the occultation summary, Table 5-24, occured with the lunar surface 59 percent illuminated. In addition, the selenocentric angle of the contact point was only 16 degrees from the southern cusp. This small cusp angle, along with the past first-quarter moon, made the sky rather bright and significantly brighter in B than V.

The RAWPLOT of the observation is shown in Figure 5-46, and the integration plot in Figure 5-47. There is no indication in either of these figures of any "wide" secondary components. A preliminary look at the subset of data selected for solution (200 milliseconds, beginning at millisecond 2395), which can be seen in Figure 5-48, rather obviously shows an unusual diffraction intensity variation. This was recognized as being attributable to a previously

TABLE 5-24 X13607: LUNAR OCCULTATION SUMMARY

STELLAR AND OBSERVING INFORMATION
Star: X13607, (SAO 080527, DM +20 2244) RA: 085559 DEC: +201558 mV: 8.2 Sp: F5 Filter: V Diaphragm: I Gain: C7+ Voltage: 1200
LUNAR INFORMATION
Surface Illumination: 59 percent Elongation from Sun: 100 degrees Altitude Above Horizon: 35 degrees Lunar Limb Distance: 366061 Kilometers Predicted Shadow Velocity: 347.5 meters/sec. Predicted Angular Rate: 0.1958 arcsec/sec.
EVENT INFORMATION
Date: April 21, 1983 UT of Event: 04:33:56 USNO V/O Code: 25 HA of Event: +603457 Position Angle: 180.2 Cusp Angle: 165 Contact Angle: -67.4 Watts Angle: 163.2
TWO-STAR MODEL PARAMETERS
Number of Data Points: 201 Number of Grid Points: 256, 256 Number of Spectral Regions: 53 Width of Spectral Regions: 50 Angstroms Limb Darkening Coefficients: 0.5, 0.05 Effective Stellar Temperatures: 5100 K, 5100 K
SOLUTIONS
Sky Level: 1429.3 (38.2) STAR 1 STAR 2
Stellar Diameter (ams): Point Source Point Source Time (relative to Bin 0): 2542.2 (3,4) 2488.7 (6.9) Stellar Intensity: 481.5 (53.1) 202.5 (49.1) Velocity (meters/sec.): 410.9 (29.7) 524.8 (100.7) Lunar Limb Slope (degrees) -16.1 (4.28) -24.3 (15.6) U.T. of Occultation: 04:33:58.827 (0.007) :58.773 (0.002)

Table 5-24. Continued.

TWO-STAR DERIVED QUANTITIES

53.4 (7.7) Temporal Separation (milliseconds): Intensity Weighted Mean L-Rate (meters/sec): 436.8 (37.6) Intensity Weighted Mean R-Rate (arcsec/sec): 0.246 (0.021) Projected Spatial Separation. Based on Predicted R-Rate: 10.46 arc-ms Based on Weighted Mean R-Rate: 13.15 (2.19) arc-ms. Brightness Ratio (Brighter/Fainter): 2.38 (0.634) 0.94 (+0.34, -0.26) Magnitude Difference: 8.58 (0.09) my of Star 1: mV of Star 2: 9.52 (0.20)

PHOTOMETRIC NOISE INFORMATION

Sum-of-Squares of Residuals:	13265444
Sigma (Standard Error):	257.541
Normalized Standard Error:	0.37652
Photometric (S+N)/N Ratio:	3.6558
(Change in Intensity)/Background:	0.4785
Change in Magnitude:	0.4246







unknown "close" stellar companion. Indeed, this was the case and was verified by the DC2 fit to the observed intensity curve.

Because of the early spectral type of the star (or properly the composite spectra of the two stars), an angular diameter determination was not expected, and indeed, the diameters of both components were found to be below the detection threshold.

The fainter of the two stars was occulted 53.4 (+/- 7.7) milliseconds before the brighter. The projected separation of the individual stellar components was found to be 13.2 (+/- 2.1) milliseconds of arc. This was determined using an intensity weighted mean R-rate of 0.2462 (+/- 0.02123) seconds of arc per second.

The determined ratio of component brightnesses (brighter/fainter) was 2.38 (+/- 0.63), or 0.94 (+0.34, -0.26) in terms of a stellar magnitude difference. Assuming an apparent V magnitude of 8.2 (as listed in the SAO catalog), the individual V magnitudes of the component stars are 8.58 (+/- 0.09) and 9.52 (+/- 0.20), respectively.

The observational data used in the DC2 fitting process, as well as the solution intensities and the residual values, are listed in Table 5-25. Table 5-26 gives the usual supplemental statistics, including the numerical ranges of the solution curve's partial derivatives which are depicted in Figure 5-49.

TABLE 5-25 X13607: OBSERVATIONS, COMPUTED VALUES, AND RESIDUALS FROM BIN 2395

1 1	NUM OBS COMP	RESID NUM	OBS CO	MP RESID	NUM OBS	COMP RE	==== SID
$ \begin{array}{c} 18 & 2019 & 2111.0 & -22.0 \\ 10 & 2101.2 & -20.0 & 83 & 2094 & 2095.8 & -34.2 & 123 & 399 & 1203.4 & -304.4 \\ 10 & 2103.4 & -205.4 & -105.4$	NUM OBS COMP 0 1845 2116.3 1 2330 2115.2 2 2404 2114.3 3 2207 2114.4 6 2132 2116.7 7 74097 2116.7 9 002 2114.9 9 202 2114.9 9 1212 2114.9 9 202 2114.5 9 122 2114.9 11 2090 2114.5 12 2231 2114.5 13 22231 2114.5 14 1748 2118.3 12 2020 2117.5 14 1748 2118.3	RESID NUL 	2647 21: 2499 22: 2529 22: 2271 22: 22203 22: 2203 22: 22499 20: 2299 20: 2299 20: 2299 20: 1883 20: 2059 20: 18807 20: 1807 20:	$\begin{array}{c} \\ \hline \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ $	NUM OBS 134 1832 135 1622 136 1312 137 1662 138 1633 139 1563 140 141 141 1005 144 1405 144 1405 145 1205 145 1305 151 130 151 155 151 155 151 155 155 155 1	CCMP RE: 1743.7 : 1743.7 : 1743.7 : 1743.1 - 33 1662.9 - 1662.9 - 1629.2 -11 1629.2 -11 1563.6 - 22 1574.3 -12 1563.6 - 1564.1 2 1543.4 - 12 1524.6 - 1524.6 - 1524.7 - 1524.1 2 1524.1 2 1	==== 30.0.91 30.0.91 30.0.91 30.0.91 30.0.91 30.0.91 30.0.92 30.02 30
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 18\\ 20(19) 2111 \\ 0 \\ 20(24) 2113 \\ 1 \\ 0 \\ 2113 \\ 1 \\ 201 \\ 105 \\ 2115 \\ 115 \\ 2115 \\$	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	2091 20 1993 20 1877 20 1752 20 1752 20 2321 20 2321 20 2321 20 2382 20 2382 20 2443 19 20 2382 20 1819 20 2382 10 1990 19 1443 19 1595 18 1575 18 1575 18 1575 18 1575 18 1575 18 1607 18	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 1505.4 \\ -505.4 \\ -505.4 \\ -805.$	463500823982667069668
	$\begin{array}{c} 19 & 2526 & 20300 \\ 40 & 2191 & 2086.1 \\ 41 & 2103 & 2088.1 \\ 42 & 1805 & 2099.6 \\ 43 & 2250 & 2113.9 \\ 44 & 2198 & 2129.1 \\ 45 & 1916 & 2124.7 \\ 46 & 2293 & 21249.2 \\ 47 & 2136 & 2150.3 \\ 48 & 2193 & 21449.2 \\ 49 & 1863 & 2136.2 \\ 50 & 2173 & 2125.0 \\ 51 & 2415 & 2114.2 \\ 52 & 2041 & 2105.6 \\ 53 & 1882 & 21007.4 \\ 55 & 2289 & 2097.4 \\ 55 & 2284 & 2097.4 \\ 56 & 2284 & 2097.4 \\ 57 & 2284 & 2097.4 \\ 58 & 2284 & 2097.4 \\ 58 & 2284 & 2097.4 \\ 58 & 2284 & 2097.4 \\ 58 & 2284 & 2097.4 \\ 58 & 2284 & 2097.4 \\ 58 & 2284 & 2097.4 \\ 58 & 2284 & 2097.4 \\ 58 & 2284 & 2097.4 \\ 58 & 2284 & 2097.4 \\ 58 & 2284 & 2097.4 \\ 58 & 2284 & 2097.4 \\ $	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	1863 19 2301 19 2496 19 12463 19 1759 20 1759 20 22231 20 12239 20 2215 20 2015 20 2015 20 2017 20 2047 20 2047 20 2047 20 2047 20 2047 20	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	$\begin{array}{c} 175 \\ 175 \\ 176 \\ 177 \\ 160 \\ 177 \\ 176 \\ 177 \\ 136 \\ 177 \\ 136 \\ 179 \\ 161 \\ 181 \\ 174 \\ 182 \\ 202 \\ 183 \\ 184 \\ 123 \\ 123 \\$	$\begin{array}{c} 1447.1\\ 1446.2\\ 11\\ 1445.2\\ -5\\ 1444.4\\ -4\\ 1443.6\\ -1\\ 1442.9\\ 1442.9\\ 11\\ 1442.9\\ 11\\ 1442.9\\ 11\\ 1442.9\\ 11\\ 1442.9\\ 11\\ 1442.9\\ 11\\ 1442.9\\ 11\\ 1442.9\\ 11\\ 1442.9\\ 11\\ 1442.9\\ 11\\ 1442.9\\ 11\\ 1442.9\\ 11\\ 1442.9\\ 11\\ 1442.9\\ 11\\ 1442.9\\ 11\\ 1442.9\\ 11\\ 1442.9\\ 11\\ 1442.9\\ 144$	52982461 57982431361 57982431361 57982431361 577766149595 5995359730 59953597730 59953597730 59953597730 59953597730 59953597730 59953597730 59953597730 59953597730 59953597730 59953597730 59953597730 59953597730 59953597730 59953597730 59953597730 59953597730 59953597730 59953597730 59953597730 59955597730 59955597730 59955597730 59955597730 59955577730 59955577730 5995577730 5995577730 5995577730 5995577730 5995577730 5995577730 5995577730 5995577730 5995577730 599557777730 59955777730 59955777777777777777777777777777777777
IN FORMATION TABLE 5-26 SUPPLEMENTAL STATISTICAL X13607:

MATRIX СE CE/ CO-VARIAN VARIAN

VEL2 5.212 5.29800 1.259800 1.259800 1.125801 1.125801 3.404804 5.3058101 -5.3058101 -5.340801
TIM2 -3.8365 -3.8365 -3.83620 -7.691300 -2.3396200 -2.339520 -3.394201 5.894801 5.894801 -5.940 E01
IN T2 -6.2855707 -6.285707 -11.750803 -14.600801 -4.600801 1.970806 2.637803 5.894801 -7.362801
DIA2 -3.092 -1.1612 -1.7592 -1.7592 -1.7592 -1.7592 -1.4612 -1.4612 -1.4612 -1.4612 -3.2058 -3.2058 -3.2058 -3.2058 -3.2058 -3.2058 -1.008 -1.00
VEL1 VEL1 1.855 1.855 1.25588 1.25588 1.25588 1.25588 1.25588 1.25588 1.25588 1.255
TIM1 -2.014E08 -1.590E01 -1.590E01 -4.583E01 -4.583E01 -4.583E02 -8.876E02 5.9910E02 5.9910E02 5.9910E02 -4.302E02
SKY -3.654E07 -3.654E07 -1.339E03 -4.5832E01 -4.5832E01 -1.438E02 -1.438E02 -1.438E02 -1.438E02 -1.438E02 -1.438E02 -1.125E01
INT1 1.020E06 1.020E06 1.020E06 1.020E06 1.1258E01 1.258E01 1.259E00 1.259E00
DIAI DIAI 3.525 3.525 3.525 15.225 16.287 10.2480 1.835 10.35810 -3.255810 -3.255810 -3.2558207 5.245810

CORRELATION MATRIX

0.873699 0.873699 0.873699 0.873699 0.55825902 0.590396 0.590396 0.9999222 1.000000EL2 5 1.376E01 -0.878964-0.8789640.558690.55840630.5840630.8831101.08999270.999927TIM2 TIM2 1.365 ± 00 1.628 ± 03 IN T2 -0.8187429 -0.8187429 0.7812055 -0.7812055 -0.7812055 -0.997529 0.997529 0.842121 -0.842421 5 З PARTIAL DERIVATIVES DIA2 ------3.602E08 -3.644E08 VEL1 3.330E03 -3.404E03 THE -----OF TIMI 2.565E01 2.732E01 -0.383576-0.313878-0.3138781.0508092-0.9999030.8104980.8104980.5781205-0.5782599UMERICAL RANGES IMIT SKY ------1.000E00 1.000E00 SKY -0.771404 -0.771404 -1.3558492 0.358492 0.358492 0.35665 0.452865 -0.442802 NUMER IN T1 1.366E00 1.038E02 IN TI IN TI 0.997229 0.91404 0.213474 0.213473 0.256433 0.256433 0.256433 0.856179 0.856179 1.000000 1.000000 0.395768 0.395768 0.395768 0.8570099 0.8736999DIAI MUMI NIM MINI NIM DIAL INTI SKY SKY VELI DIA2 INT2 INT2 VEL2

ī.







Figure 5-50. NOISEPLOT of the occultation of X13607.

The NOISEPLOT of the 4096 milliseconds of observational data (Figure 5-26), as well as the mean offset and one sigma values, are nearly identical to those obtained for the earlier occultation of X13567. This was expected as both stars were nearly the same apparent V magnitude, and the observing conditions were essentially identical.

Finally, the power spectra of the observation centered on the event, the pre-occultation and star-plus-sky signals, and the two-star model curve are shown in the POWERPLOT, Figure 5-51.

Since the WWUB time signal was of lower quality for this event, the Coordinated Universal Times of geometrical occultation were not quite as well determined as for the X13534 event. The fainter star underwent disappearance at 04:33:58.773 (+/- 0.009 seconds). The brighter component followed this at 04:34:58.827 (+/- 0.007 seconds).

ZC1462

The occultation of the K0 star ZC1462 on April 22, 1983, was observed under rather poor conditions, as detailed in the following excerpt from the observing log regarding that event: " ... seeing poor, often worse than 8 or 10 seconds of arc. Occasional image blooming causes the star to become invisible in the photometer viewing optics. ... " Despite the bad photometric conditions, a digital photoelectric record of the occultation was obtained and is presented in the RAWPLDT, Figure 5-52.







The integration plot, Figure 5-53, shows no evidence of stellar duplicity. Though there is an apparent change in slope at approximately 1800 milliseconds, it is a gradual change in level, rather than a sharp discontinuity typical of a second star. In addition, the original slope of the integration curve returns at approximately 2200 milliseconds. This behavior is due only to variations in atmospheric transparency.

The formal solution to the observed intensity curve is shown on the FITPLOT, Figure 5-54. As may be seen from the figure, or from the occultation summary, Table 5-27, the best fit to the raw data calls for a source diameter of roughly 9 milliseconds of arc. This anomalously large angular diameter is quite unexpected considering the relative faintness of this star, and is believed to be spurious.

The power spectrum of the star-plus-sky signal determined from 1000 millisecond data samples taken between 1.5 ond 0.5 seconds preceeding the determined time of geometrical occultation is shown on Figure 5-56. Contributions to the background noise due to the low frequency power components are quite important. Indeed, in the range of 10 to 70 Hertz these are dominant over the power spectral signature of the solution curve by one to two orders of magnitude, respectively. It is precisely, however, in this power spectral region that







Figure 5-55. FITPLOT of the occultation of ZC1462 to the Fourier transformed data.

TABLE 5-27 ZC1462: LUNAR OCCULTATION SUMMARY

STELLAR AND OBSERVING INFORMATION
Star: ZC1462 (SAO 098862, DM +17 2156) RA: 095639 DEC: +163225 mV: 7.38 Sp: KO Filter: V Diaphragm: I Gain: C7 Voltage: 1200
LUNAR INFORMATION
Surface Illumination: 70 percent Elongation from Sun: 114 degrees Altitude Above Horizon: 37 degrees Lunar Limb Distance: 366100 kilometers Predicted Shadow Velocity: 809.68 meters/sec. Predicted Angular Rate: 0.4969 arcsec/sec.
EVENT INFORMATION
Date: April 22, 1983 UT of Event: 05:13:38 USN0 V/O Code: 15 HA of Event: +562152 Position Angle: 124.6 Cusp Angle: 775 Contact Angle: -7.0 Watts Angle: 104.1
MODEL PARAMETERS
Number of Data Points: 201 Number of Grid Points: 256 Number of Spectral Regions: 53 Width of Spectral Regions: 50 Angstroms Limb Darkening Coefficient: 0.5 Effective Stellar Temperature: 5100K
SOLUTIONS Raw Fit Fourier Fit
Stellar Diameter (ams): 9.27 (3.69) 9.02 (2.44) Time: (relative to Bin 0): 2601.5 (1.1) 2601.1 (0.7) Pre-Event Signal: 2307.9 (25.1) 2299.8 (16.8) Background Sky Level: 1187.6 (25.3) 1184.9 (16.3) Velocity (meters/sec): 831.4 (78.6) 833.9 (50.4) Lunar Limb Slope (degrees): -6.57 (5.56) -6.29 (3.57) UT of Occultation: 05:13:37.1989 (0.0014) 37.1985 (0.0011)
PHOTOMETRIC NOISE DATA Raw Fit Fourier Fit
S-O-S of Residuals: 11526000 1253100 Sigma (Standard Error): 240.07 79.154 Normalized Standard Error: 0.2143 0.1420 Photometric (S+N)/N Ratio: 5.6663 8.0434 Intensity Change/Background: 0.74324 0.94092 Change in Magnitude: 0.72132 0.72002



the signature of the occultation curve is most important. Hence, the observed data are corrupted by intensity variations due to atmospheric noise which mimic the variations seen in an occultation curve.

This effect is seen rather dramatically in the second FITPLOT presented (Figure 5-55). Here, the observed data were Fourier transformed and all power components of frequency higher than 150 Hertz removed. After inverse transformation, the Fourier smoothed data were fit by the DC procedure. The solutions obtained in a fit to the smoothed data, as expected, are nearly identical to the fit to the raw data. What is obvious from visual inspection is that the undulating character of the data, in terms of both frequency and amplitude, is very similar to the fringing effects expected due to the occultation.

It is therefore apparent that the noise characteristics of the sky, at the time of this event, were such that the formal solution of the angular diameter is highly suspect, and in fact most probably erroneous. Under these conditions, with present reduction and analysis techniques, it is almost certainly impossible to ascertain the true angular diameter of the star. Though a formally determined diameter is presented here, a strong "caveat emptor" is placed on the result.

The observations extracted from the raw data set (as well as the computed intensity curve and the residuals) are listed in Table 5-28. The variance-covariance and correlation matrices of the formal solution are given in Table 5-29. Also found on this table are the ranges of the numerical values of the partial derivatives of the intensity curve, as depicted on the PDPLOT, Figure 5-57.

The photometrically poor conditions for occultation photometry are further noted by an examination of the NOISEPLOT of the 4096 milliseconds of observational data presented in Figure 5-58. The mean signal level does correspond with the peak of the distribution function of the residuals, but the distribution function itself is highly skewed. The mean signal level over the entire data window is more than 2 percent down from the mean level during the 201 milliseconds centered on the time of the event, and the one signa level of the distribution function is approximately 11.5 percent.

A final note on observations, such as this one, taken under photometrically noisy conditions is warranted. In this case, the noise associated with the raw data has a one sigma level of 21.4 percent of the mean intensity. This by itself is not sufficient to dismiss the possibility of obtaining a meaningful solution from the observed intensity curve. What is

TABLE 5-28 ZC1462: OBSERVATIONS, COMPUTED VALUES, AND RESIDUALS FROM BIN 2500 NUM OBS COMP F NUM OBS COMP RESID NUM OBS COMP RESID RESID -----2298.5 22.5 2293.2 -120.2 2297.6 -47.6 2313.8 -120.8 2338.1 214.9 2361.7 -117.7 2373.5 -256.5 2365.0 138.0 2334.8 272.2 2290.0 133.0 2243.6 -214.6 2210.8 15.2 ____ _____ _ _ _ 134 1458 135 1361 136 858 137 870 138 1379 139 1403 371.0 -200.1 337.0 1194.5 1194.1 1193.7 263.5 67 2321 68 2173 69 2250 70 2193 71 2553 72 2244 73 2117 74 2503 75 2607 76 2423 77 2029 78 2226 79 2371 80 2311 81 2095 82 1797
 667
 2311

 2111
 2311

 12648
 2311.0

 1923
 2311.0

 1923
 2311.0

 2155
 2311.0

 22055
 2311.0

 2555
 2311.0

 2555
 2311.0

 26107
 2311.0

 2767
 2311.0

 1779
 2311.0

 1682
 2311.0

 1682
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 2311.0
 0 -335.7 23 870 1193.4 1379 1193.1 1403 1192.8 -323 -388.0 185.9 210.2 131.4 -156.0 1403 1192.0 1324 1192.6 1104 1192.3 1074 1192.1 140 141 142 143 1324 244.0 -88 -118 .1 668.0 8 -8.9 1183 1191.9 308.0 ā -214.6 178.3 1370 778 144 1191.7 196.0 144 145 146 147 148 1191.5 1191.3 1191.0 1190.9 1190.6 1190.6 1190.6 1190.3 1190.3 1190.3 1190.0 1190.0 1190.0 1190.0 1190.0 1190.0 1190.0 1190.0 1190.0 1190.0 1190.0 1190.4 11 -413 2210.8 2204.0 2229.1 2284.2 2360.7 2445.5 2524.4 2584.9 2618.1 1014 1143 1257 1269 1109 1182 -429.0 167.0 -177 .3 -429.0 -200.0 -261.0 -620.0 -48.2 81.9 -189.2 81 82 83 84 66.0 14 78 .1 149 150 151 152 153 155 155 156 157 158 159 5 2639 2996 2478 2471 -532.0 -627.9 -196.0 272.0 -81.8 193.5 471.6 16 -8.6 1464 1336 1271 1247 273.5 85 86 -106.9 18 -147.1 ğ 80.7 2641 2543 2496 1643 2203 360.0 756.0 -117.0 2619.4 2588.7 2529.3 22529.3 22447.3 22238.3 22125.6 1908.3 1810.1 17212.2 1573.6 1573.6 1419.8 21.6 87 -45.7 $\begin{array}{c} 3067 \\ 2211 \\ 2194 \\ 2311 \\ 2215 \\ 2311 \\ 2755 \\ 2311 \\ 2755 \\ 2311 \\ 2755 \\ 2311 \\ 2781 \\ 2311 \\ 2216 \\ 2311 \\ 2216 \\ 2311 \\ 2213 \\ 2311 \\ 2213 \\ 2311 \\ 2213 \\ 2311 \\ 2213 \\ 2311 \\ 2213 \\ 2311 \\ 2213 \\ 2311 \\ 2213 \\ 2311 \\ 2213 \\ 2311 \\ 2213 \\ 2311 \\ 2213 \\ 2311 \\ 2213 \\ 2311 \\ 2213 \\ 2311 \\ 23$ 88 -125.1 13.0 -222.0 -33.3 1065 1203 890 912 992 995 995 997 999 999 999 -803.6 76.0 1643 2203 2313 2411 2218 1920 1553 1691 968 444.0 -144.324 25 26 27 28 444.0 329.9 472.0 -95.0 986 -203.9 74.7 285.4 1415 1278 $-95 \cdot 02$ $-268 \cdot 29$ $32 \cdot 01$ $91 \cdot 72$ $-228 \cdot 02$ $-864 \cdot 26$ $-1366 \cdot 42$ $-2096 \cdot 22$ $-2366 \cdot 35$ $-6196 \cdot 42$ $-74 \cdot 43$ $-74 \cdot 43$ $-136 \cdot 43$ 88.3 1189.7 1189.7 1189.6 1189.5 1189.5 1189.4 1189.4 1189.3 1278 1241 1357 1087 1032 1119 924 936 -257.0 51.3 162 163 .4 iq 29 -102.5 -30.1 -34.2 270.8 85.4 185.2 -154.6 103.8 187.1 -199.3 33.2 -79.6 154.5 164 165 166 167 Īĕ08 -157.5 -70.4 -253.3 -50.3 153.88 197.88 296.9 -50.1 -321.0 -422.0 -161.9 1844 1599 1475 1605 100 168 169 170 171 172 173 174 175 176 1189.2 1343 1189.2 1387 1189.2 1219 1189.1 1516 1189.1 1139 11 869 34 35 1383.6 36 1229 1457 1383.0 1353.2 1327.9 1306.9 1289.3 1274.8 1262.6 1252.5 104 105 1515 1318 38 39 106 1310 107 1090 108 1308 109 1183 110 1407 -1 40 41 767 1189.0 42 177 1027 1188.9 -136.3 27.3 269.8 279.2 293.1 -101.9 214.1 262.1 -223.8 278.2 -141.8 -67.8 1631 1015 1331 1234.0 1236.8 1225.7 1221.4 1217.7 1214.5 1207.3 1205.5 1205.5 1204.6 1201.4 387.0 -221.8 100.2 513.3 238.6 -144.7 86.5 -29.8 99.6 -40.3 182.0 -302.6 1403 1451 1188.9 1188.9 111 112 113 44 45 180 965 1188.8 46 1188.8 2604 2310.9 114 115 1739 1460 1073 181 182 1467 47 2310.9 2311.2 2313.0 2314.2 2312.7 2309.7 2309.1 215.8 360.0 188.8 -95.7 21.3 -76.1 1047 1188.8 1121 1188.8 48 2527 2527 2673 2503 2217 2331 2233 2140 2369 2080 2294 2107 2367 183 184 185 186 187 188 189 190 116 49 1139 1166 -49. 1188.7 30 1301 1188.7 118 1182 1242 1126 119 120 1309 1167 1188.7 1188.7 -62. 2312.4 2316.9 2317.9 2314.1 -172.4 52.1 -237.9 -20.1 -201.6 120 1107 121 1224 122 902 123 999 124 1045 125 1367 126 1119 1291 1075 1231 102.4 -113.6 42.4 -155.6 1188.6 1188.6 1188.6 1188.6 1188.6 56 1045 1367 1119 1115 1125 1345 91 -156.4 -1 1200.4 166.6 1 92 12 58 2308.6 1209 20. -80.4 193 1188.5 59 2367 2306.1 60.9 -106.7 2202 2010 2367 2237 1199.4 1198.6 1197.8 1197.1 1196.5 2308.7 127 128 129 -83.6 -72.8 147.9 194 195 196 1188.5 1188.5 1188.5 326.5 60 954 856 -234 61 2314.8 -304.8 62 2320.9 46.1 129 1345 130 1051 131 1273 132 913 133 1079 -145.5 197 198 1117 1188.5 5 -87.3 63 2324.3 178.6 212.6 136.6 2647 2539 2390 2323.3 2317.7 2308.4 323.7 1195.9 1195.4 1194.9 77.1 -282.4 -115.9 1367 1401 1188.4 64 άğ 1188.4 1325 66 81.6

TABLE 5-29 ZC1462: SUPPLEMENTAL STATISTICAL INFORMATION

VARIANCE/CO-VARIANCE MATRIX

DIAM	PREI	POST	TIME	VELO
3.207E ⁻ 16	6.839E 08	-2.060E 08	-4.716E ⁰ 9	3.890E ⁻ 10
6.839E 08	6.296E02	2.217E01	-6.631E00	3.594E-01
-2.060E 08	2.217E01	6.404E02	-7.124E00	3.994E-01
-4.716E 09	-6.631E00	-7.124E00	1.306E00	-7.414E ⁻ 02
3.890E ⁻ 10	3.594E 01	3.994E ⁻ 01	-7.414E ⁻ 02	6.182E-03

CORRELATION MATRIX

	DIAM	PREI	POST	TIME	VELO
DIAM	1.000000	0.959624	-0.441850	-0.367280	0.348224
PREI	0.959624	1.000000	-0.176967	-0.603263	0.586464
POST	-0.441850	-0.176967	1.000000	-0.669212	0.683702
TIME	-0.367280	-0.603263	-0.669212	1.000000	-0.999770
VELO	0.348224	0.586464	0.683702	-0.999770	1.000000

NUMERICAL RANGES OF THE PARTIAL DERIVATIVES

	MAXIMUM	MIN IMUM
DIAM	4.426E09	-3.867E09
PREI	1.278E00	6.941E ⁻ 04
POST	9.993E 01	-2.781E 01
TIME	1.127E02	-8.438E01
VELO	1.411E03	-1.846E03









important is whether the noise is dominant in the low or high frequency regime. To further verify that this is indeed true, a computational test was performed. A model intensity curve was needed for the test. For convenience, the solution curve for this occultation was used. To this model curve, noise (whose characteristics are described by the distribution function shown in Figure 5-58) was applied. The addition of this noise was subject to a temporal distribution such that the power components in the low frequency domain (less than 70 Hertz) were not highly dominant. In two trial attempts, the DC procedure had no trouble recovering the model parameters. In addition, the same noise distribution as used in the first trial was applied to a point source, and a point source solution was obtained.

Fortunately, the time-scale and zero reference of the occultation event were found through this, and additional numerical experimentation, to be moderately insensitive to this type of predominantly low frequency noise. Thus, the determination of the time of geometrical occultation of 05:13:37.199 (+/-0.001) Coordinated Universal Time is fairly reliable.

<u>X18067</u>

The KO star X18067 (mV=7.9) was occulted on June 18, 1983, under clear, photometrically steady skies. The fast photoelectric digital record of this event, shown in Figure 5-59, was obtained by John P. Oliver and

Martin England who (as noted in the occultation summary, Table 5-30) employed a Johnson V filter for the observation.

The integration plot of the observational data (Figure 5-60) hints at a possible "wide" stellar duplicity approximately 150 milliseconds before the obvious disappearance. To ascertain if this drop (also seen on the RAWPLOT) was real, a detailed integration plot was produced containing the data from 3100 to 3625 milliseconds. To aid in the visual interpretation of the graph, the integrated data were subjected to 5-point unweighted smoothing before being plotted. It is apparent in Figure 5-61 that a real, sharp change of slope does occur at approximately millisecond 3265.

The mean intensity (averaged over 125 milliseconds) prior to this was 2515 counts, and afterward. 2488 counts. This drop of 27 counts represents a lowering of the signal intensity by approximately 0.011 magnitudes. This seems like a rather small amount. However, as will be seen, the formal error of the pre-occultation intensity found by the DC fitting process at the time of the primary event was only roughly 0.007 magnitudes. Hence, the one sigma certainty of detection of a "wide" component is 64 percent better than the determination of the mean star-plus-sky level. Therefore, one can

TABLE 5-30

X18067: LUNAR OCCULTATION SUMMARY

STELLAR AND OBSERVING INFORMATION _____ Star: X18067 (SAD 119227, DM +05 2587) RA: 120641 DEC: +043627 mV: 7.9 Sp: K0 Filter: V Diaphragm: I Gain: B10+ Voltage: 1200 LUNAR INFORMATION _____ Surface Illumination: 53 percent Elongation from Sun: 93 degrees Altitude Above Horizon: 48 degrees Lunar Limb Distance: 370414 kilometers Predicted Shadow Velocity: 364.91 meters/sec. Predicted Angular Rate: 0.2032 arcsec/sec. EVENT INFORMATION _____ UT of Event: 02:16:54 Date: June 18, 1983 HA of Event: +354352 USNO V/O Code: 26 Position Angle: 63.8 Cusp Angle: 40N Contact Angle: +61.5 Watts Angle: 40.6 MODEL PARAMETERS ------Number of Data Points: 201 Number of Grid Points: 4096 Number of Spectral Regions: 53 Width of Spectral Regions: 50 Angstroms Limb Darkening Coefficient: 0.5 Number of Grid Points: Effective Stellar Temperature: 5100K SOLUTIONS _____ Stellar Diameter (ams): Point Source Stellar Diameter (amp): Time: (relative to Bin 0): 3422.1 (1.5) Pre-Event Signal: 2459.2 (16.8) Pre-Event Signal: 2459.2 (1.87) Background Sky Level: 1738.6 (20.8) Velocity (meters/sec.): 365.6 (10.7) Lunar Limb Slope (degrees): -1.83 (1.68) U.T. of Occultation: 02:16:52.406 (0.007) PHOTOMETRIC NOISE INFORMATION Sum-of-Squares of Residuals: 6146800 Sigma (Standard Error); 175.316 Normalized Standard Error: 0.24332 Photometric (S+N)/N Ratio: 5.1099 (Change in Intensity)/Background: 0.41441 Change in Magnitude: 0.37643







ascribe a confidence level to the detection of a "wide" secondary component of 90 percent.

The best solution for the observation of the primary occultation of X18067 (depicted graphically in Figure 5-62) found the star to be a point source. This was not unexpected as X18067 has been classified as a main sequence star, thus giving it a distance modulus of approximately 2.0, and an anticipated diameter of 0.85 solar radii. At this distance the corresponding angular diameter would be only 1/25 millisecond of arc, well below the detection threshold of roughly one millisecond of arc.

From the determined R-rate of 0.2036 (+/- 0.0059) seconds of arc per second, the "wide" secondary component was found to have a projected separation of 32 milliseconds of arc. The uncertainty in this is somewhat conjectural. In this case it is safe to assume that the time of secondary disappearance can be determined by inspection of the detailed integration plot to an accuracy of roughly 10 milliseconds. This gives an uncertainty of approximately 2 milliseconds of arc in the projected separation.

The subset of the observational data used in the fitting procedure, the computed intensities, and the resulting residuals are listed in Table 5-31. The supplementary statistical information, including



 TABLE 5-31

 X18067: OBSERVATIONS, COMPUTED VALUES, AND RESIDUALS FROM BIN 3300

NUM OBS	COMP	RESID	NUM	OBS	COMP	RESID	N UM	OBS	COMP	RESID
$ \begin{array}{c} \text{NUN} \\ \text{OSS} \\ OS$	$ \begin{array}{c} \begin{array}{c} \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\$	$\begin{array}{c} & = & \\ R = & 1 & 1 & 1 & 1 \\ S = & 1 & 1 & 1 & 1 \\ S = & 1 & 1 & 1 & 1 \\ S = & 1 & 2 & 1 & 5 & 1 & 3 \\ S = & 1 & 2 & 1 & 5 & 1 & 3 \\ S = & 1 & 2 & 1 & 5 & 1 & 3 \\ S = & 1 & 2 & 1 & 5 & 1 & 3 \\ S = & 1 & 2 & 1 & 5 & 1 & 2 & 3 \\ S = & 1 & 2 & 1 & 1 & 2 & 5 & 1 \\ S = & 1 & 2 & 1 & 1 & 2 & 5 & 1 \\ S = & 1 & 2 & 1 & 2 & 1 & 2 & 1 \\ S = & 1 &$	======================================	$= \begin{array}{c} 0 & -2 & -97 \\ 0 & -2 & -22 & 228 \\ 0 & -2 & -24 & -28 \\ 0 & -2 & -24 & -28 \\ 0 & -2 & -24 & -28 \\ 0 & -2 & -24 & -28 \\ 0 & -2 & -24 & -28 \\ 0 & -24 & -28 & -28 \\ 0 & $		$ \begin{array}{c} 1.16.099.97.55(28.07.845)(38.91.87715)(58.19.9976)(1.0.1443)(39.63750)(2.2818)(38.945)(39.14716)(3$	= 14364 + 14784 + 14884 + 14	$= \underbrace{0}_{-1} \underbrace{1}_{2773} \underbrace{1}_{5774} \underbrace{1}_{99} \underbrace{1}_{99} \underbrace{1}_{2773} \underbrace{1}_{773} \underbrace{1}_{97} \underbrace{1}_{99} \underbrace{1}_{97} \underbrace{1}_$	COUPL	$\begin{array}{c} \text{RESID} & \\ \text{RESID} & \\ -176, 47, 47, 47, 47, 47, 47, 47, 47, 47, 47$
48 2063 49 2217 50 2753 51 2729 53 2872 54 2818 55 2647 56 2586 57 2464 57 2464 57 2464 57 2464 57 2464	2403.2 2426.1 24253.6 2483.1 2512.4 25512.4 25581.2 2 2581.2 2 2 2581.2 2 2581.2 2 2 2581.2 2 2 2581.2 2 2 2581.2 2 2 2581.2 2 2 2 2581.2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2	-340 •2 -209 •1 299 •4 245 •9 269 •6 333 •2 257 •6 71 •7 -117 •2 -145 •7	115 116 117 118 119 120 121 122 123 124 125 126	2059 1966 1848 2056 2007 1779 1887 2062 2031 1872 1863 1995	2017.9 1998.8 1980.8 1964.1 1948.4 1933.8 1920.2 1907.5 1895.8 1884.9 1874.7 1865.3	41.1 -32.8 91.9 58.6 -154.8 -33.2 154.5 135.2 -12.9 -11.7 129.7	182 183 184 185 186 187 188 189 191 191 192	1783 1676 1841 1756 1772 1973 1850 1592 1585 1768 2008	1746.1 1745.9 1745.7 1745.2 1745.0 1745.0 1744.8 1744.8 1744.7 1744.5 1744.3 1744.1 1744.0	36.9 95.3 73.5 105.3 -152.5 -152.5 -159.3 264.0
60 2428 61 2271 62 2623 63 2459 64 2599 65 2315 66 2488	2531.2 2502.7 2471.1 2438.2 2406.0 2376.4 2351.0	-103.2 -231.7 151.9 20.8 193.0 -61.4 137.0	127 128 129 130 131 132 133	1874 1787 2055 2000 1888 1741 1703	1856.6 1848.6 1841.1 1834.2 1827.8 1821.9 1816.5	17.4 -61.6 213.9 165.8 60.2 -80.9 -113.5	194 195 196 197 198 199 200	1579 1680 1808 1596 1762 1711 1895	1743.8 1743.7 1743.6 1743.4 1743.3 1743.2 1743.1	-164.8 -63.7 64.4 -147.4 18.7 -32.2 151.9

range of the partial derivatives presented on the PDPLOT (Figure 5-63) are given in Table 5-32.

A cursory examination of the raw intensity plot reveals that while the expected high frequency scintillation and photon noise sources are present, there is virtually no variation in atmospheric transparency on longer time-scales. This led to a rapidly convergent solution with very low formal statistical errors. The distribution function of the observational noise is shown in Figure 5-64. The usual occultation power spectra are displayed in Figure 5-65.

The Coordinated Universal Time of geometrical occultation was determined to be 02:16:52.406 (+/- 0.007 seconds).

ZC2209 (32 Librae)

The bright (mV=5.92) star ZC2209 was occulted on June 22, 1983. ZC2209 is of spectral type K0 and has been classified as a giant (luminosity class III). The occultation event was observed by John P. Oliver, Martin England, and Howard L. Cohen with the instrumental configuration specified in Table 5-33. The observing log indicates that clouds appeared seven minutes after the event. Oliver, however, has indicated that prior to this the transparency of the sky was quite good and the seeing steady and calm. The post-event cloud cover was of the low patchy cumulus type, which is common in north Florida during the summer. Therefore, the onset of



Figure 5-63. PDPLOT of the occultation of X18067.



Figure 5-64. NOISEPLOT of the occultation of X18067.

TABLE 5-32

X18067: SUPPLEMENTAL STATISTICAL INFORMATION

VARIAN CE/CO-VARIAN CE MATRIX

DIAM	PREI	POST	TIME	VELO
1.008E ⁻ 14	2.581E ⁻ 07	-3.101E ⁰⁷	-1.077E-08	9.371E ⁻ 11
2.581E 07	2.874E02	1.917E01	-7.082E00	4.368E-02
-3.101E_07	1.917E01	4.509E02	-1.010E01	5.972E 02
-1.077E_08	-7.082E00	-1.010E01	2.379E00	-1.439E 02
9.371E ⁻ 11	4.368E ⁻ 02	5.972E ⁻ 02	-1.439E ⁻ 02	1.171E 04

CORRELATION MATRIX

	DIAM	PREI	POST	TIME	VELO
DIAM	1.000000	0.701669	-0.797370	0.254354	-0.231840
PREI	0.701669	1.000000	-0.130101	-0.498781	0.518495
POST	-0.797370	-0.130101	1.000000	-0.781331	0.766733
TIME	0.254354	-0.498781	-0.781331	1.000000	-0.999731
VELO	-0.231840	0.518495	0.766733	-0.999731	1.000000

NUMERICAL RANGES OF THE PARTIAL DERIVATIVES

	MAXIMUM	MIN IMUM
DIAM	3.274E08	-3.212E08
PREI	1.366E00	6.222E ⁰ 3
POST	9.938E 01	-3.659E ⁻ 01
TIME	3.393E01	-3.612E01
VELO	6.036E03	-6.035E03



TABLE 5-33 ZC2209: LUNAR OCCULTATION SUMMARY

STELLAR AND OBSERVING INFORMATION

Star: ZC 2209, (32 Librae, SAO 159280, DM -16 4089) RA: 152719 DEC: -163935 mV: 5.92 Sp: KO-III Filter: V Diaphragm: I Gain: B8 Voltage: 1000

LUNAR INFORMATION

Surface Illumination:	90	percent
Elongation from Sun:	143	degrees
Altitude Above Horizon:	44	degrees
Lunar Limb Distance:	386528	Kilometers
Predicted Shadow Velocity:	625.	71 meters/sec.
Predicted Angular Rate:	0 .	.3339 arcsec/sec.

EVENT INFORMATION

Date: June 22, 1983	UT of Event:	03:12:54
USNO V/O Code: 76	HA of Event:	+033323
Position Angle: 111.6	Cusp Angle:	85S
Contact Angle: +8.3	Watts Angle:	98.8

MODEL PARAMETERS

Number of Data Points:	201
Number of Grid Points:	256
Number of Spectral Regions:	53
Width of Spectral Regions:	50 Angstroms
Limb Darkening Coefficient:	0.5
Effective Stellar Temperature:	5100 K

SOLUTIONS

 Stellar Diameter (ams):
 12.18 (1.86)

 Time: (relative to Bin 0):
 631.2 (0.8)

 Pre-Event Signal:
 3020.2 (12.2)

 Background Sky Level:
 2025.8 (12.2)

 Velocity (meters/sec.):
 673.7 (45.5)

 Lunar Limb Slope (degrees):
 -10.88 (4.16)

 U.T. of Occultation:
 03112153.356 (0.007)

PHOTOMETRIC NOISE INFORMATION

Sum-of-Squares of Residuals:	2494300
Sigma (Standard Error):	111.68
Normalized Standard Error:	0.11231
Photometric (S+N)/N Ratio:	9.904
Intensity Change/Background:	0.49086
Change in Magnitude:	0.43359

cloud cover after the event did not adversley affect the observation.

The digital photoelectric record obtained is presented in Figure 5-66. As can be seen, the sky was quite well behaved throughout the 4096 milliseconds of data acquisition. The observation is seen early in the data window. One of the observers, who was using the LODAS system for the first time, nearly forgot to stop the data acquisition process after the event. The integration plot of the occultation record (Figure 5-67) shows no indication of stellar duplicity.

The solution to the observed intensity curve is depicted graphically in Figure 5-68. As indicated, a rather large angular diameter (12.18 +/- 1.86 milliseconds of arc) was determined by the differential corrections fitting procedure. Assuming the spectral type and luminosity classification for 2C2209 are correct, they would indicate a distance on the order of 115 parsecs; hence, the diameter found is an order of magnitude larger than one would expect. This disparity remains unresolved, as the fit is rather good and difficult to dismiss. Examination of the PDPLOT (Figure 5-69) clearly indicates that the region of sensitivity of the observed curve to variations in all the parameters was well considered and that numerical noise was minimal.







The distribution function of the residuals shown in the NOISEPLOT (Figure 5-70) is obviously Gaussian in nature. The mean observed intensity is down only 0.1 percent from the computed intensity, with a one sigma width in the distribution function of only 4 percent. In addition, the photometric (S+N)/N ratio of the observation was 9.9. Hence, the photometric noise characteristics of the sky were very well behaved.

The POWERPLOT (Figure 5-71) for this event shows the relative power contributions up to a limiting frequency of 250 Hertz. While the power components of the occultation signal and star-plus-sky background are normally examined to higher frequencies, this was not possible for this observation. Since only roughy 500 milliseconds of data prior to the event were retained in the data acquisition process, the power spectrum of the star-plus-sky background will have a Nyquist cutoff frequency of 250 Hertz. It is apparent, however, that the frequency components in the observed curve due to the occultation signal are dominant over the background noise in the region of importance in the occultation solution (i.e., the low frequency domain).

Hence, the angular diameter determined from the solution of the observed occultation intensity curve is unaffected by both the spatial and temporal characteristics of the background noise, and is not spurious due to any possible numerical problems. The





Figure 5-70. NOISEPLOT of the occultation of ZC2209.



reason for the disparity between the anticipated small angular diameter and the larger diameter inferred from the observation could lie in an erroneous spectral classification. However, this tentatve suggestion is put forth without much force, and the question at this time remains open.

The Coordinated Universal Time of geometrical occultation for this event was found to be 03:12:53.356 (+/- 0.007 seconds). The observed and computed intensity values (and the residuals) are given in Table 5-34. The variance-covariance and correlation matrices, and the numerical ranges of the partial derivatives are listed in Table 5-35.

ZC3214

As noted in the occultation summary, Table 5-36, the occultation of ZC3214 was observed on November 13, 1983. The photoelectric record of the event is shown in Figure 5-72. As can be seen in this RAWPLOT, the transparency of the sky was quite steady on time-scales of a few tenths of a second. This is reflected in the very flat nature of the integration plot (Figure 5-73). Examination of the integration plot shows no indication of any secondary events.

ZC3214 is an AO star, with an apparent V magnitude of 6.6. Thus, it was not suprising when the solution, depicted graphically in Figure 5-74, failed to reveal a

TABLE 5-34

ZC2209: OBSERVATIONS, COMPUTED VALUES, AND RESIDUALS FROM BIN 520 NUM OBS COMP RESTD NUM OBS COMP RESID NUM OBS COMP RESID ---_____ --------------2036.4 2035.8 2035.3 2034.8 2034.8 -74.0 67 2661 3007.6 -346.6 1894 1874 1959 2151 2068 2155 0 2949 3023.0 134 -142 .4 3023.0 3023.0 3023.0 3023.0 2661 3007.0 2848 3000.2 3145 2994.2 2884 2991.0 3019 2992.1 2870 2999.0 -152.2 150.8 -107.0 135 136 -161.8 2853 -170.0 68 69 70 66.0 3089 2884 3019 137 138 139 3176 153.0 116 . 3085 3082 3042 3093 4 3022.9 3022.9 62.1 71 26.9 33.7 121.1 16.5 2034.3 2033.9 2033.5 2033.1 2032.7 2032.4 2032.1 2031.6 5 72 2870 73 2873 74 2900 -129.0 3012.8 3033.7 3023.1 2050 1929 2028 18.9 -139.8 140 3023.1 3023.1 69.9 104. -133.7 141 -133.7 13.0 -43.1 -112.5 -14.2 3.9 -87.3 75 3074 2969 3023 .1 3023 .0 3023 .0 3023 .1 3023 .2 3023 .3 3023 .2 3023 .1 3023 .1 3023 .2 3023 .4 3023 .4 3061.0 142 -4.7 2020 2080 1937 1983 -69.0 9 2954 3142 3009 76 3050 3015 3147 3093.1 3127.5 3161.2 47.6 143 -14.1 10.8 298.7 95.8 -40.1 77 78 144 145 - 95 - 48 . .8 2031 2031.6 2175 2031.3 2033 2031.1 79 3195 80 3127 81 3109 23 3034 3191.1 146 -0.6 3322 3119 2983 3214.3 147 143.7 -119.2 4 148 1.9 81 3109 3228.2 82 3184 3231.1 83 3281 3222.0 84 3271 3200.7 85 3127 3167.7 86 3177 3124.2 1857 1948 2030.7 2030.7 2030.3 2030.3 2030.0 2029.8 2029.5 2029.5 2029.4 2029.3 2029.2 2029.0 2029.0 2029.0 2028.9 173.9 5 149 150 151 152 153 154 155 156 157 -4/.1 59.0 70.3 -40.7 -40.1 -20.1 158.8 175.6 167.6 -39.3 70.8 -173.9 -82.7 -51.5 -94.3 -44.1 3003 3182 3199 3191 2984 1979 1936 8 â 3023.4 3124.2 3071.9 52.8 60.1 1986 3023 .4 3023 .3 3023 .2 3023 .3 3023 .5 3023 .7 3023 .7 3023 .7 3023 .7 3023 .2 3023 .2 3023 .2 3023 .8 3024 .8 3024 .8 $\begin{array}{c} 86 \\ 87 \\ 3132 \\ 3071 \\ 98 \\ 9925 \\ 3071 \\ 99 \\ 3072 \\ 90 \\ 2057 \\ 3071 \\ 90 \\ 2057 \\ 3072 \\$ 87 2111 2161 2125 2020 20 22 22 22 23 22 24 25 27 27 3132 81.0 131.2 3094 2954 3160 -53.5 -23.2 -143.7 70.8 -69.3 136.5 85.3 -12.3 -99.1 -207.2 -52.8 -12.6 -319.8 -319.8 3109 3060 -99.4 46.9 158 159 160 1930 2101 2023 1907 -99.4 71.7 -6.2 22.0 26.1 3000 3011 2924 2816 $\begin{array}{c} -2.2\\ 98.69\\ 98.69\\ -209.7\\ -26.99\\ -209.7\\ -191.69\\ 91.4\\ 48.3\\ -1.1\\ -1.1\\ -34.5\\ 22.2\\ -114.5\\ -34.5\\ -34.5\\ -107.8\\ -93.2\\ -$ 161 -1 28 29 2055 2127 2027 2168 162 163 2971 2028.8 98.2 -1.7 3ó 3012 2705 2028.7 164 2028.6 139.4 165 -319.8 -330.2 -124.2 61.6 48.4 -41.5 -57.3 2694 3024 2 3023 2 30222 6 30222 6 30225 3 30226 7 30225 8 30226 7 30225 8 30226 7 3026 7 3 166 167 1966 1987 2028.6 2899 3084 3071 2028.5 34 35 36 168 169 170 171 172 173 174 175 176 177 3071 3285 3052 3000 3034 3080 3071 48.4 261.3 26.7 -26.5 7.3 54.2 46.9 -86.1 86.7 37 101.8 39 106 22,28 107 2228 108 2089 109 2006 110 2075 111 2334 112 2245 114 2147 115 1966 117 1956 118 1913 119 1976 121 2145 122 2066 123 1998 -79.1 40 -146.0 41 -145.9 -143.9 19.1 20.2 45.2 -154.7 42 2936 3131 2985 2733 2967 2955 3293 43 110.1 44 178 -30.0 -289.9 -59.3 -75.1 260.2 179 46 **1**80 61.3 4Ť 181 3030.1 3032.8 3033.1 3030.3 3025.1 3019.2 3014.7 3013.5 3016.3 48 182 -88 .6 49 3227 193.9 -41.3 11.9 -48.2 173.5 -114.8 212.4 276.3 129.9 225.4 30.1 -95.4 170.0 183 11.5 17.5 59.5 50 51 52 53 54 2989 -109.1 184 2065.1 2061.4 2058.2 2055.4 2052.9 2050.6 2048.7 3037 2971 -148.4 185 -148.4 -64.2 -79.4 92.1 -44.6 -50.7 205.1 87.7 119.6 -8.4 66.7 -33.3 -58.3 -33.2 186 187 3188 3019 2902 3235 3307 188 189 3022.6 190 191 1998 2252 2133 2190 2114 2024 1965 2015 2046.9 124 125 126 127 128 129 130 3038.4 58 3168 192 193 194 -111.2 59 3070 3044.1 2043.9 146.1 146.1 71.4 -17.4 -75.4 -24.4 -79.6 -54.8 96.9 199.8 3272 3046.6 2042.6 59.9 15.9 43.9 3046.0 3045.9 3042.4 3037.0 194 195 196 197 198 2043 2071 2031 61 3076 2041.4 2947 2040.4 2027.1 43.9 3.9 85.0 3207 2039.4 2038.6 2037.8 2037.1 64 3043 3030.4 12.6 -54.1 -278.4 1959 1983 2134 2112 2011 2053 2027.0 2027.0 2027.0 2027.0 2969 2737 3023.1 3015.4 99 -16.0 26.0

TABLE 5-35

ZC2209: SUPPLEMENTAL STATISTICAL INFORMATION

VARIANCE/CO-VARIANCE MATRIX

DIAM	PREI	POST	TIME	VELO
8.149E ⁻ 17	2.649E 08	2.525E ⁻ 10	-2.870E ⁻ 09	1.899E ⁻ 10
2.649E 08	1.476E02	9.753E00	-3.250E00	1.445E ⁰¹
2.525E ⁻ 10	9.753E00	1.478E02	-3.257E00	1.454E ⁻ 01
-2.870E 09	-3.250E00	-3.257E00	6.850E 01	-3.114E ⁻ 02
1.899E ⁻ 10	1.445E ⁰¹	1.454E ⁰¹	-3.114E ⁰²	2.066E 03

CORRELATION MATRIX

	DIAM	PREI	POST	TIME	VELO
DIAM	1.000000	0.996428	-0.137148	-0.647710	0.645788
PREI	0.996428	1.000000	-0.109177	-0.661815	0.659550
POST	-0.137148	-0.109177	1.000000	-0.663632	0.665138
TIME	-0.647710	-0.661815	-0.663632	1.000000	-0.999983
VELO	0.645788	0.659550	0.665138	-0.999983	1.000000

NUMERICAL RANGES OF THE PARTIAL DERIVATIVES

	MAXIMUM	MIN IMUM
DIAM	3.608E09	-4.206E09
PREI	1.212E00	1.209E 03
POST	9.988E 01	-2.122E01
TIME	6.867E01	-3.459E01
VELO	1.120E03	-1.217E03
TABLE 5-36

ZC3214: LUNAR OCCULTATION SUMMARY STELLAR AND OBSERVING INFORMATION _____ Star: ZC3214 (SAD 164756, DM -18 6037) RA: 215551 DEC: -175835 mV: 6.6 Sp: A0 Filter: V Diaphraom: I Gain: C5+ Voltage: 1200 LUNAR INFORMATION Surface Illumination: 54 percent Elongation from Sun: 95 degrees 20 degrees Altitude Above Horizon: 20 degrees Lunar Limb Distance: 402184 Kilometers Predicted Shadow Velocity: 694.7 meters/sec. Predicted Angular Rate: 0.3563 accsec/sec. EVENT INFORMATION -----Date: November 13, 1983 UT of Event: 03:35:44 USNO V/O Code: 28 HA of Event: 540443 Position Angle: 31.0 Cusp Angle: 51N Contact Angle: +28.4 Watts Angle: 51.8 MODEL PARAMETERS _____ Number of Data Points: 201 Number of Grid Points: 256 Number of Spectral Regions: 53 Width of Spectral Regions: 50 Angstroms Limb Darkening Coefficient: 0.5 Effective Stellar Temperature: 10800 K SOLUTIONS _____ Stellar Diameter (ams): Point Source Time: (relative to Bin 0): 2255.7 (0.8)
 Pre-Event Signal:
 2481.5
 (18.6)

 Background Sky Level:
 1484.8
 (18.3)

 Velocity (meters/sec.):
 732.3
 (20.4)
Lunar Limb Slope (degrees): -9.22 (1.68) U.T. of Occultation: 03:35:43.107 (0.018)

PHOTOMETRIC NOISE INFORMATION

Sum-of-Squares of Residuals:	6385000					
Sigma (Standard Error):	178.68					
Normalized Standard Error:	0.17927					
Photometric (S+N)/N Ratio:	6.5781					
(Change in Intensity)/Background:	0.67124					
Change in Magnitude:	0.5576					





measurable angular diameter. The fit of the model intensity curve to the observations appears to be quite good. This is reflected in the low formal errors of the solution parameters and is seen by visual inspection of the FITPLOT even as far out as the fourth order diffraction minimum.

As evidenced by the PDPLOT, Figure 5-75, the portions of the intensity curve most sensitive to parametric variation were well covered in the solution process. The model intensity curve was fit to the data presented in Table 5-37. The usual supplemental statistical information is compiled in Table 5-38.

The NOISEPLOT of the observation is presented as Figure 5-76. As noted, the RMS background noise through the 4096 milliseconds of the observation was approximately 8 percent. The power spectra of the occultation event, the star-plus-sky signal, and the best fit model curve are shown on the POWERPLOT, Figure 5-77.

As in the case of the occultation of ZC1221, a fine determination of the time of geometrical occultation, made possible by the small (0.8 millisecond) error of the formal solution, was thwarted by a noisy WWVB radio signal. The Coordinated Universal Time of geometrical occultation was 03:35:43.929 (+/- 0.018 seconds).





Figure 5-76. NOISEPLOT of the occultation of ZC3214.

TABLE 5-37 ZC3214 OBSERVATIONS, COMPUTED VALUES, AND RESIDUALS FROM BIN 2215

20	JZ1+ 4	OD D D D D D D D D D D D D D D D D D D	TTT TON 0 9	00112 0							
====	=====	=======					==========		=====		
AT 17M	ORC	COMB	DECTD	NTTM	OBC	COMP	RESTD	NIM	OBS	COMP	RESTD
NULL	003	CONF	VEDID	NOPI	005	COLL	TUDO ID	11 011	000	00112	ALD U LD
		A 1 A 1			~ ~			10/	1 2 0 7	1/02 1	106 1

1.011	OBS	CUMP	RESID	NOPI	UBS	COMP	RESID				
1 - 0 1 2 3 4 5 6 7 8 9 0 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	$D_{1-2}^{-1} Z_{2}^{-1} Z_{2}^{$	Luni- 2445.3, 444.4, 1, 4 2445.3, 4 2445.3, 4 2445.3, 4 2445.3, 4 2445.4, 4	$ \begin{array}{c} & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 2 \\ & 1 \\ & 2 \\ & 1 \\ & 2 \\ & 1 \\ & 2 \\ & 1 \\ & 2 \\ & 1 \\ & 2 \\ & 1 \\ & 2 \\ & 1 \\ & 2 \\ & 1 \\ & 2 \\ & 1 \\ & 2 \\ & 1 \\ & 2 \\ & 1 \\ & 2 \\ & 1 \\ & 2 \\ & 1 \\ & 2 \\ & 1 \\ & 2 \\ & 1 \\ & 2 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 2 \\ & 1 \\ & 1 \\ & 1 \\ & 2 \\ & 1 $	$ \begin{array}{c} & & \\ & & $	$ \sum_{i=1}^{n-1} \sum_{j=1}^{n-1} \sum_{i=1}^{n-1} \sum_{j=1}^{n-1} \sum_{i=1}^{n-1} \sum_{j=1}^{n-1} \sum_{j=1}^{n-1} \sum_{i=1}^{n-1} \sum_{j=1}^{n-1} \sum_{i=1}^{n-1} \sum_{j=1}^{n-1} \sum_{i=1}^{n-1} \sum_{j=1}^{n-1} \sum_{i=1}^{n-1} \sum_{j=1}^{n-1} \sum_{j=1}^{n-1} \sum_{i=1}^{n-1} \sum_{j=1}^{n-1} \sum_{j=1}^{n-1} \sum_{i=1}^{n-1} \sum_{j=1}^{n-1} \sum_{i=1}^{n-1} \sum_{j=1}^{n-1} \sum_{j=1}^{n-1} \sum_{i=1}^{n-1} \sum_{j=1}^{n-1} \sum_{j=1}^{n-1} \sum_{j=1}^{n-1} \sum_{i=1}^{n-1} \sum_{i=1}^{n-1}$	$\begin{array}{c} 1&1&1&2&2&2&2&2&2&2&2&2&2&2&2&2&2&2&2&2$	$\begin{array}{c} 1 \\ 1 \\ 1 \\ 2 \\ 2 \\ 1 \\ 2 \\ 2 \\ 1 \\ 2 \\ 2$	$\sum_{i=1}^{n-1} \sum_{i=1}^{n-1} $	$\begin{array}{c} 1 \\ 1 \\ 1 \\ 1 \\ 1 \\ 1 \\ 1 \\ 1 \\ 1 \\ 1 $	$\begin{array}{c} 1_{442}\\ 1_{443}\\$	$\begin{array}{c} 1 \\ 1 \\ -1 \\ -1 \\ -1 \\ -1 \\ -1 \\ -1 \\ $
389 40 42 43 445 445 447 48 49	2185 2436 2507 2521 2174 2322 2348 2588 2705 2437 2003	2492.6 2509.44 2506.42 2457.25 2457.25 2471.44 2534.1 2534.1 2534.1 2534.3 2450.32 2450.32	-307.6 -75.4 -70.4 23.8 -276.5 -149.4 -159.6 53.9 1749.4 -59.3 -59.3 -59.3 -447.2	105 106 107 108 109 110 111 112 113 114 115 116	1330 1331 1240 1627 1860 1575 1546 1474 1645 1654 1677 1771	1590 •9 1578 •1 1567 •1 1557 •5 1542 •3 1536 •1 1536 •8 1526 •1 1528 •6 1515 •4	-120.09 -247.1 -327.1 69.5 310.7 32.7 9.9 -56.8 118.9 138.4 255.6	172 173 174 175 176 177 178 178 181 182 183	1378 1511 1384 1335 1439 1795 1670 1463 1523 1644 1563 1498	$1486 \cdot 877$ $1486 \cdot 766$ $1486 \cdot 766$ $1486 \cdot 514$ $1486 \cdot 514$ 148	-108.8 24.3 -102.7 -151.6 308.5 183.5 -23.4 36.6 157.6 76.7
50 51 52 53 55 55 55 55 50 55 60 61	2239 2413 2315 2563 2728 3179 2628 2281 2265 2427 2668	2419.92 2425.2 25669.44 25239.4 25579.7 2547.7 2485.6 2418.5 2368.5 2368.5 2368.5 2368.5 2407.3	-180.9 -12.2 -151.3 -20.9 -6.4 148.3 631.3 142.4 -137.5 -108.3 58.5 260.7	11/ 118 119 120 121 122 123 124 125 126 127	1607 1503 1560 1416 1363 1455 1468 1089 1195 1612 1503 1552	1512.7 1510.3 1508.2 1504.5 1502.9 1501.6 1500.4 1499.3 1497.4 1496.7	94.3 -7.3 -90.2 -141.5 -47.9 -33.6 -411.4 -304.3 113.7 5.6 55.3	184 185 187 1887 1889 1991 192 193 194	1351 1523 1387 1560 1287 1303 1379 1495 1892 1718 1698 1527	1486.3 1486.2 1486.2 1486.2 1486.1 1486.1 1486.0 1486.0 1486.0 1486.0 1486.0 1486.0	-135.3 36.8 -99.2 73.8 -199.1 -183.1 -107.1 9.0 406.0 232.0 212.0 41.1
62 63 64 65 66	2722 2482 2765 2487	24/7.7 2557.5 2622.4 2653.5 2641.8	-0.7 164.5 -140.4 111.5 -154.8	129 130 131 132 133	1451 1564 1513 1527 1614	1496.0 1495.3 1494.7 1494.1 1493.6	-45.0 68.7 18.3 32.9 120.4	196 197 198 199 200	1311 1526 1623 1721	1485.9 1485.9 1485.9 1485.9 1485.8	-36.9 -174.9 40.1 137.1 235.2

30.4

TABLE 5-38 ZC3214: SUPPLEMENTAL STATISTICAL INFORMATION

VARIANCE/CO-VARIANCE MATRIX

DIAM	PREI	POST	TIME	VELO
5.694E ⁻ 16	5.662E 08	-4.205E 08	-1.344E ⁰ 9	4.047E ⁻ 11
5.662E 08	3.371E02	7.832E00	-2.875E00	6.537E_02
-4.205E 08	7.832E00	3.262E02	-2.672E00	5.986E_02
-1.344E 09	-2.875E00	-2.672E00	5.657E_01	-1.283E_02
4.047E ⁻ 11	6.537E 02	5.986E 02	-1.283E ⁻ 02	4.051E ⁰ 4

CORRELATION MATRIX

	DIAM	PREI	POST	TIME	VELO
DIAM	1.000000	0.845265	-0.692455	-0.164543	0.173361
PREI	0.845265	1.000000	-0.200047	-0.660013	0.666554
POST	-0.692455	-0.200047	1.000000	-0.592852	0.585544
TIME	-0.164543	-0.660013	-0.592852	1.000000	-0.999958
VELO	0.173361	0.666554	0.585544	-0.999958	1.000000

NUMERICAL RANGES OF THE PARTIAL DERIVATIVES

	MAXIMUM	MINIMUM
DIAM	1.811E09	-1.802E09
PREI	1.365E00	1.019E ⁰³
POST	9.990E 01	-3.655E 01
TIME	9.068E01	-9.686E01
VELO	3.942E03	-3.974E03



X31590

The rather slow occultation disappearance of the K0 star X31590 is apparent on the plot of the 4096 milliseconds of raw observational data shown in Figure 5-78. The relative faintness of the star (mV=8.7) and the high sKy brightness, due to a 72 percent illuminated moon, resulted in a final photometric S+N/N of only 2.55. Fortunately, the seeing was unusually good, and the photometer J diaphragm (as noted in the occultation summary, Table 5-39) was selected.

The integration plot of the event, Figure 5-79, shows no indication of any disappearances other than that of X31590 itself.

The graphical depiction of the solution is shown in Figure 5-80. Three hundred and fifty milliseconds of data (given in Table 5-40) were included to be fit by the DC process. This rather lengthy data set was necessary due to the somewhat lengthened time-scale of the event. Not unexpectedly, given the poor S+N/N ratio and the faintness of the star, no detectable stellar disc was found.

The partial derivatives of the intensity curve, with respect to each of the solution parameters, are shown graphically in Figure 5-81. The numerical ranges of these derivatives are listed, along with the other usual supplemental statistics, in Table 5-41.

TABLE 5-39 X31590: LUNAR OCCULTATION SUMMARY

STELLAR AND OBSERVING INFORMATION Star: X31590, (SAO 165704, DM -10 6166) RA: 222634 DEC: -095159 mV: 8.7 Sp:K0 Filter: V Diaphragm: J Gain: C7 Voltage: 1200 LUNAR INFORMATION _____ Surface Illumination: 72 percent Elongation from Sun: 116 degrees Altitude Above Horizon: 50 degrees Lunar Limb Distance: 397056 kilometers Predicted Shadow Velocity: 574.2 meters/sec. Predicted Angular Rate: 0.2983 arcsec/sec. EVENT INFORMATION _____ Date: November 15, 1983 UT of Event: 01:51:29 USND V/0 Code: 33 HA of Event: +071416 Position Angle: 75.6 Cusp Angle: 785 Contact Angle: -24.5 Watts Angle: 99.4 MODEL PARAMETERS _____ 351 Number of Data Points: Number of Grid Points: 256 Number of Spectral Regions: 50 Width of Spectral Regions: 53 Angstroms Limb Darkening Coefficient: 0.5 Effective Stellar Temperature: 5100 K SOLUTIONS _____ Stellar Diameter (ams): Point Source Time: (relative to Bin 0): 2401.3 (5.3) Pre-Event Signal: 2396.0 (16.5) Background Sky Level: 1999.5 (29.4) Velocity (meters/second): 229.6 (14.1) Lunar Limb Slope (deorees): +33.2 (1.4) U.T. of Occultation: 01:02:40.701 (0.014) PHOTOMETRIC NOISE INFORMATION ____ Sum-of-Squares of Residuals: 22911940 Sigma (Standard Error): 225.8576 0.654175 Normalized Standard Error: Photometric (S+N)/N Ratio: 2.54997 (Change in Intensity)/Background: 0.19834 Change in Magnitude: 0.19645



Figure 5-79. INTPLOT of the occultation of X31590.



TABLE 5-40 X31590: OBSERVATIONS, COMPUTED VALUES, AND RESIDUALS FROM BIN 2140

===:	=====			=====	=====			=====		=======	
NUM	OBS	COMP	RESID	NUM	OBS	COMP	RESID	NUM	OBS	COMP	RESID
NUM 0 123345 6789 10 11	OBS 2529 2570 2398 2643 2495 2291 2681 2517 24681 2517 24631 2211	COMP 23 97 .5 23 98 .1 23 98 .5 23 98 .8 23 98 .8 23 98 .6 23 98 .0 23 97 .3 23 96 .0 23 95 .5 23 95 .5 23 95 .5	RESID 	NUM 67 68 69 70 71 72 73 74 75 76 77 77 8 79	OBS 2719 2303 2128 2267 1800 1981 2600 2751 2403 3080 2830 2830 2830	COMP 2400.2 2402.3 2404.0 2405.3 2405.6 2405.6 2405.6 2404.6 2402.8 2400.8 2400.8 2400.8 2400.8 2407.7 2397.7 2394.7 2399.1	RESID 	NUM 134 135 136 137 138 139 140 141 142 143 144 145 146	OBS 21999 2160 1941 25170 23910 2391 2316 2354 2218 2118 2118 2195 2375	COMP 2353.0 2351.1 2351.0 2352.6 2355.9 2367.3 2374.9 2383.2 2392.9 2402.7 2412.7 2412.4	RESID -154.00 -191.1 -410.0 164.4 34.1 -350.9 -29.3 -7 -58.9 -292.7 -174.9 -292.7 -352.6
13 14 15 16 17 18 20 21 223 24 5 25	2156 2408 2175 2271 2321 2591 2744 2388 2553 2632 2400 2281 2455	23 95 9 23 96 5 23 97 3 23 98 1 23 98 1 23 98 3 23 99 3 23 99 3 23 99 3 23 99 3 23 98 1 23 98 1 23 97 3 23 98 1 23 97 3 23 98 1 23 96 5 23 96 5 23 96 5 23 96 5 23 96 5 23 96 5 23 97 3 23 98 1 23 96 5 23 96 5 23 97 5 23 96 5 23 96 5 23 97 5 23 98 5 23 99 5 23 90 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5 5	-239.9 -11.5 -222.3 -127.1 -77.8 191.7 344.5 -11.3 154.2 233.9 2.8 -115.3 59.4 2.8	80 81 82 83 85 86 87 88 90 91 92	2317 2132 2303 2619 2573 1699 2681 2400 2495 2243 2151 2436 2200	2387.0 2385.7 2385.4 2386.6 2390.1 2393.4 2397.3 2401.4 2405.5 2409.3 2412.5 2414.8	-70.0 -253.7 -82.4 233.0 185.4 -691.1 287.6 -287.6 -287.6 -287.6 -287.6 -287.6 -287.6 -287.6 -287.6 -287.6 -28.5 -214.8	147 148 150 151 152 153 155 156 157 158 159	2436 1913 2183 2151 2076 2255 3051 2719 2637 2157 2262 2539 2444	2431.72 2440.2 2447.8 2454.9 2452.9 2462.9 2463.9 22463.9 22463.9 22463.9 224553.9 224553.9 24553.9 24553.9 224553.9 24555.9 24555.9	4.3 -527.2 -264.88 -303.1 -382.9 -207.3 587.1 174.8 -301.9 -191.9 91.4 -4.1
207 228 200 312 3345 367 899	2843 2717 2256 2019 2655 21301 22370 2236 2375 2131 2777 2476 2639	2395 • 1 2395 • 9 2395 • 5 2395 • 5 2396 • 3 2398 • 1 2398 • 1 2399 • 1 2397 • 2 2398 • 7 2397 • 7	447.91 -139.00 -376.5258.7 -237.3 -67.22 170.9 -163.7 -25.0 -268.9 377.3 -25.0 -268.9 377.3 241.3	93 94 95 97 98 100 101 102 103 104 105	2725 2883 2831 2586 2707 2805 2515 2496 2507 2423 2112 2265 2410	2416.0 2414.7 2412.3 2408.8 2404.3 23993.8 23993.8 2388.4 2383.3 2378.9 2375.4 2373.1 2372.2	309.0 467.0 173.7 298.2 397.7 115.8 102.2 -428.4 123.7 -428.4 123.7 -428.4 123.7 -44.1 -108.1	160 1612 1623 1654 1656 1667 1689 1670 1711 1723	2248 2872 2177 27065 1954 2156 2303 2477 2105 2479 2070 2019	2431.1 2421.5 2411.1 2400.3 2389.4 2378.4 2367.6 2348.1 2339.6 2332.1 2325.7 2320.7 2317.1	-463-1 4565-51 -223-36-66 -413-66 -413-66 -413-66 -45-1 137-4 137-4 -227-13 -250-7 -298-1
40 42 44 44 44 44 44 44 44 55 12 55 55 55	2851 2346 2435 2380 2372 2443 2644 2169 2280 2516 2599 2433 2649 2599 2433 2599 24337 2900	2396.6 2395.7 2394.5 2394.5 2394.4 2394.8 2395.6 2395.6 2397.9 2397.9 2399.1 2400.1 2400.9 2401.1	404.4 -49.7 40.1 -14.5 -22.4 48.2 248.4 -227.6 -117.9 116.9 198.9 30.1 735.7 498.9	107 108 109 110 111 112 113 114 115 116 117 118 119 120	1872 2203 2239 2423 2881 2486 2098 2655 2512 2231 2184 2219 2219 2084	2372 *8 2374 *8 2378 *2 2382 *9 2388 *6 2395 *1 2402 *0 2408 *9 2415 *6 2421 *7 2426 *8 2430 *8 2433 *3 2434 *2	-500.8 -171.8 -139.2 40.1 492.4 90.9 -304.0 246.1 96.4 -190.7 -242.8 -211.8 -211.8 -350.2	174 175 176 177 178 182 182 182 183 185 185 187	2405 2304 2394 2150 2231 2178 1995 1797 2127 2691 3139 2457 2257 2307	2314.9 2314.2 2315.0 2317.3 2321.0 2326.1 2332.5 2340.0 2348.6 2358.1 2368.4 2379.4 2379.4 2390.7	90.1 -10.2 -167.3 -90.0 -148.1 -337.5 -543.0 -221.6 332.9 770.6 -133.9 -70.6 -133.9
5557890 5557890 66123456 665 665 665	2933 2235 2317 2509 2195 2495 2471 2743 2654 2089 2250 2380 2111	2400.55 2399.4 2398.1 2396.5 2395.6 2392.6 2392.6 2392.1 2392.9 2392.9 2394.2 2394.2 2396.0 2398.0	532.5 -164.4 -81.1 112.5 -200.0 101.4 78.4 350.9 261.8 -303.9 -144.2 -16.0 -287.0	121 122 123 124 125 126 127 128 129 130 131 132 132	2328 2771 2802 2572 2600 2148 2029 2356 2335 2679 2407 2157 2157	2433.52 2431.1 2427.2 2421.8 2415.8 2415.8 2407.8 2399.7 2399.3 2382.9 2382.9 2382.9 2387.7 2361.4 2356.5	-105.45 339.9 374.8 150.2 184.7 -259.8 -370.7 -35.3 304.1 39.3 -204.4 242.5	188 189 190 191 192 193 194 195 196 197 198 199 198 200	2396 2396 2313 2787 2578 2242 2223 1887 2579 2746 2482 2589 2690	2414 2426 2438 2450 2462 2462 2462 2462 2462 2483 250 2483 2502 3 2502 3 2517 8 2524 2524 2529 6	51.38 -30.88 -125.84 336.49 -231.00 -260.53 -606.3 765.55 -35.8 64.84 140.4

TABLE 5-40. CONTINUED.

=== NUM	0BS	COMP	RESID	NUM	OBS	COMP	RESID	NUM	OBS	COMP	RESID
$\begin{array}{c} N & -2222222222222222222222222222222222$	$\begin{array}{c} 0 \\ 0 \\ -8 \\ -8 \\ -8 \\ -2 \\ -2 \\ -2 \\ -2 \\ -2$	$\begin{array}{c} \text{COMP}\\ \hline \text{COMP}\\ 2353, \text{C}, \text{C}\\ 2537, \text{C}, \text{C}\\ 2537, \text{C}\\ 2539, \text{C}\\ 2235, \text{C}\\ 223$	$ \begin{array}{c} {\rm EES} {\rm ID} \\ {\rm asympt} \\ {\rm asy$	$ \begin{array}{l} {\tt N} {\tt UM} \\ -2552 \\ 22553 \\ 22554 \\ 22554 \\ 22554 \\ 22554 \\ 22554 \\ 22554 \\ 22554 \\ 22557 \\ 22557 \\ 22557 \\ 22557 \\ 22557 \\ 22557 \\ 22557 \\ 22577 \\ 22577 \\ 22577 \\ 22774 \\$	$\begin{array}{c} 0BS\\ -2075313\\ 227506\\ 2277506\\ 1980\\ 3222429\\ 11849\\ 8498\\ 224299\\ 9591\\ 11920\\ 22321\\ 11920\\ 22321\\ 11920\\ 22321\\ 11920\\ 22321\\ 11920\\ 22321\\ 11920\\ 22321\\ 11920\\ 22321\\ 11920\\ 2232222\\ 22321\\ 11920\\ 2232222\\ 22323\\ 11920\\ 22322222\\ 22323\\ 11920\\ 2232222\\ 22323\\ 11920\\ 2232222\\ 22323\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 223222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 223222\\ 22332\\ 11920\\ 223222\\ 22332\\ 11920\\ 223222\\ 22332\\ 11920\\ 223222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 223222\\ 22332\\ 11920\\ 223222\\ 22332\\ 11920\\ 2232222\\ 22332\\ 11920\\ 22322222\\ 22332\\ 11920\\ 22322222\\ 22332\\ 11920\\ 22322222\\ 22332\\ 11920\\ 223222222\\ 22332\\ 11920\\ 22322222\\ 22332\\ 11920\\ 22322222\\ 22332\\ 11920\\ 223222222\\ 22332\\ 11920\\ 2232222222\\ 22332\\ 11920\\ 2232222222\\ 22332\\ 11920\\ 223222222222222\\ 22332\\ 11920\\ 223222222222222222222222222222222222$	COMP 2147.9 2147.9 2141.7 2135.7 2113.0 21124.3 21124.3 2113.9 2109.0 2109.0 2109.0 2109.0 2005.5 2009.7 2005.5 2009.7 2005.7 20	$ \begin{array}{c} {\rm RESID} \\ -772, 93\\ -772, 93\\ -772, 93\\ -772, 93\\ -772, 93\\ -772, 93\\ -772, 93\\ -772, 93\\ -732, 93\\ -732, 90, 93\\ -732, 90, 93\\ -732, 9$	$\label{eq:nonlinear} \begin{split} & n - 300350788901121545677899012234556789012333333333333333333333333333333333333$	$\begin{array}{l} 0BS\\ -065757217601\\ 217601\\ 1172954221202931\\ 22029312929221230322212312922212323221222122212222$	COMP 2018.7. 2018.1. 2017.5. 2015.2. 2017.2. 2017.2. 2017.2. 2017.2. 2016.2. 2007.2. 2005.2. 2	RESID 486.3,9 486.3,9 486.3,9 486.3,9 486.3,9 486.3,9 486.3,9 486.2,9 486.2,9 486.2,9 486.2,9 497.7,1 496.2,9 497.7,2 497.7
250	2099	2154.5	-55.5	300	2123	2019.4	103.6	350	2580	2004.5	575.5





Figure 5-82. NOISEPLOT of the occultation of X31590.

TABLE 5-41 X31590: SUPPLEMENTAL STATISTICAL INFORMATION

VARIAN CE/CO-VARIAN CE MATRIX

DIAM	PREI	POST	TIME	VELO
3.067E ⁻ 14	3.795E ⁻ 07	-9.342E ⁻ 07	-2.879E 08	1.217E ⁻ 10
3.795E ⁰⁷	2.716E02	2.809E01	-2.066E01	4.788E 02
-9.342E07	2.809E01	8.649E02	-5.633E01	1.270E-01
-2.879E 08	-2.066E01	-5.633E01	2.765E01	-6.389E 02
1.217E ⁻ 10	4.788E 02	1.270E ⁰¹	-6.389E ⁻ 02	1.983E 04

CORRELATION MATRIX

	DIAM	PREI	POST	TIME	VELO
DIAM	1.000000	0.463093	-0.915775	0.629985	-0.621074
PREI	0.463093	1.000000	-0.071425	-0.330036	0.338810
POST	-0.915775	-0.071425	1.000000	-0.874404	0.868639
TIME	0.629985	-0.330036	-0.874404	1.000000	-0.999931
VELO	-0.621074	0.338810	0.868639	-0.999931	1.000000

NUMERICAL RANGES OF THE PARTIAL DERIVATIVES

	MAXIMUM	MIN IMUM	
DIAM	2.073E08	-2.086E08	
PREI	1.367E00	1.276E ⁰ 2	
POST	9.872E ⁰¹	-3.666E ⁰¹	
TIME	1.133E01	-1.208E01	
VELO	4.875E03	-5.024E03	



The distribution function of the residual intensities is shown in the NOISEPLOT, Figure 5-82. Despite the weakness of the stellar signal against the background sky, the noise characteristics of the sky itself were well behaved.

The POWERPLOT showing the usual Fourier power spectra is presented as Figure 5-83. As typical with longer duration events, the power components important to the occultation curve are shifted toward lower frequencies.

The Coordinated Universal Time of geometrical occultation was determined to be 001:02:40.703, with a relatively large one sigma uncertainty of 0.014 seconds. 2C3458 (336 B. Aquarii)

The second occultation observed on the night of November 15, 1983 was that of the moderately bright K0 star ZC3458 (336 B. Aquarii). Even with the moon 73 percent illuminated, this event was quite promising, given the good seeing conditions which continued to prevail since the early part of the evening.

Disaster (literally, "bad star", appropriately enough), however, usually strikes at the best of times. Blow attests that observers at the University of Texas have " . . . found at least 57 ways to foul up an occultation [observation] . . ." (1983, p. 9-14). A fifty-eighth may now be contributed by the University of Florida.

The loss of a high speed photoelectric record of the ZC3458 event was due to a mechanical failure, resulting in two self-cancelling errors as far as the LODAS video strip chart display was concerned. The "O" Key on the LODAS command Keyboard had become mechanically sticky and electrically "bouncy". As a result, the data acquisition rate, which should have been entered as 001 samples per millisecond was entered as 010 samples per millisecond. Similarly, the video strip chart recorder display rate which was typed in as 010 display points per two acquisition times was taken as 001. Hence, the video display was updating at a rate indicative of millisecond data acquisition, while data were actually being sampled and stored at a slow rate of 100 samples per second.

The observer freely admits his error for not checking the printed observing log at the time of setting up the instrumental system. The RAWPLOT of the 40.96 seconds of data obtained (with a much-too-fast amplifier time constant of 2 KiloHertz (see Chapter 1)), is shown in Figure 5-84. The integration plot of the event, Figure 5-85, has no indication of any very widely separated components.

All, however, was not lost as the time of occultation was determined to an accuracy as well as could be obtained at this sampling rate. The





Figure 5-85. INTPLOT of the occultation of ZC3458.

Coordinated Universal Time of "geometrical" occultation was seen at 03:37:57.996 (+/- 0.006 seconds). X01217

Four occultation disappearances were observed on the night of November 17, 1983. The first of these was the occultation of the K0 star X01217. The raw photoelectric data which was obtained is shown in Figure 5-86. The observation of this 7.7 magnitude star was made with a Johnson V filter (see Table 5-42). The integration plot of the 4096 milliseconds of acquired data, Figure 5-87, shows no indication of stellar duplicity.

Based on the predicted R-rate of 516.3 meters/second, the usual number of data points (two hundred milliseconds) were extracted for a solution determination. The preliminary run of the DC fitting procedure indicated that the actual shadow velocity was considerably slower and that the data set used was of insufficient length. As a result, the observation was re-reduced using 350 milliseconds of data. The final solution which was obtained gave an R-rate of 284.8 (+/- 16.0) meters/second. This corresponds to a moderately large local lunar limb slope of +28.3 (+/- 1.8) degrees.

The resulting best fit is shown graphically in Figure 5-88. As can be seen the observational data were rather noisy and led to a photometric S+N/N ratio of

	TA	ABLE 5-42	
X01217:	LUNAR	OCCULTATION	SUMMARY

STELLAR AND OBSERVING INFORMATION

Star: X01217, (SAO 129029, DM -00 0139) RA: 005444 DEC: -000404 mV: 7.7 Sp: K0 Filter: V Diaphragm: I Gain: B10+ Voltage: 1200

LUNAR INFORMATION

Surface Illumination:	87	percent
Elongation from Sun:	138	degrees
Altitude Above Horizon:	42	degrees
Lunar Limb Distance:	391396	Kilometers
Predicted Shadow Velocity:	516	.3 meters/sec.
Predicted Angular Rate:	0	.2721 arcsec/sec.

EVENT INFORMATION

Date: November	17, 1983	UT of Event:	00:02:09
USNO V/O Code:	74	HA of Event:	-401416
Position Angle:	99.0	Cusp Angle:	52S
Contact Angle:	-44.6	Watts Angle:	121.4

MODEL PARAMETERS

Number of Data Points:	351
Number of Grid Points:	256
Number of Spectral Regions:	53
Width of Spectral Regions:	50 Angstroms
Limb Darkening Coefficient:	0.5
Effective Stellar Temperature:	5100 K

SOLUTIONS

Stellar Diameter (ams):	2.88	(4.98)
Time: (relative to Bin 0):	2091.6	(3.5)
Pre-Event Signal:	2927.6	(13.5)
Background Sky Level:	2495.6	(20.7)
Velocity (meters/sec.)	284.8	(16.0)
Lunar Limb Slope (degrees):	+28.3	(1.8)
U.T. of Occultation: 00:0	2:07.790	(0.004)

PHOTOMETRIC NOISE INFORMATION

Sum-of-Squares of Residuals:	14186840
Sigma (Standard Error):	201.330
Normalized Standard Error:	0.46613
Photometric (S+N)/N Ratio:	3.1453
(Change in Intensity)/Background:	. 0.1730
Change in Magnitude:	0.1733







TABLE 5-43 X01217: OBSERVATIONS, COMPUTED VALUES, AND RESIDUALS FROM BIN 1850

=== NUM	OBS	COMP	RESID	n UM	OBS	COMP	RESID	N UM	OBS	COMP	RESID
=N1-0123456789011234567890123456789012345678901234567890112345678901123222222233333333333333344444444	$\begin{array}{c} \textbf{z} = \textbf{z} = \textbf{z} \\ \textbf{z} = \textbf{z} = \textbf{z} \\ \textbf{z} \\ \textbf{z} = \textbf{z} \\ $	$\begin{array}{c} \\ -& -& -& -& -& -& -& -& -& -& -& -& -& $	$\begin{array}{c} \text{RESID} \\ \text{RESID} \\$	 NUMM 6788 700 7723 77456 800 802 802 802 802 970 7723 77477 7723 77477777777777777777777	$\begin{array}{c} \textbf{z} = \textbf{z} \\ \textbf{z} = $	$\begin{array}{c} \textbf{C} \\ $	$ \begin{array}{c} \text{RESID} \\ \text{RESID} \\ \text{C}_{-1} \\ $	$\begin{array}{c} & & & & & & \\ & & & & & & & \\ & & & & $	BS 32996 33071 227777 30751 227777 31072 227752 22900 227255 22725 22900 227255 22725 229000 22725 22725 22900 22725 22725 207777 27725 207777777777	COMP 	$\begin{array}{c} {\rm RESID} \\ {\rm RESID} \\ {\rm result} \\ $
38 39 40 41 42 43 45	2702 3124 3317 2890 2598 2743 2894	2928.8 2928.3 2928.0 2927.9 2928.1 2928.4 2928.4 2928.9	-226.8 195.7 389.0 -37.9 -330.1 -185.4 -34.9	105 106 107 108 109 110 111	3069 2696 2906 3071 2901 2986 3204	2941.6 2943.7 2944.3 2943.5 2941.0 2937.2 2932.3	127.4 -247.7 -38.3 127.5 -40.0 48.8 271.7	172 173 174 175 176 177 178	2721 3007 2799 2971 3296 3083 3264	2845.2 2845.7 2848.5 2853.6 2860.9 2870.1 2881.0	-124.2 161.3 -49.5 117.4 435.1 212.9 383.0
4478 4478 4490 55123	2748 3055 3171 2967 2959 2824 2925 2855	2929.8 2930.0 2930.0 2929.7 2929.6 2928.6 2928.1 2928.1	-181.8 124.9 241.0 37.3 29.8 -104.6 -3.1 -72.7	113 114 115 116 117 118 119	3073 2624 2671 2953 2703 3051 2867 2620	2 921 .1 2 915 .9 2 911 .7 2 908 .9 2 907 .8 2 908 .6 2 911 .4 2 916 .0	151.9 -291.9 -240.7 44.1 -204.8 142.4 -44.4	180 181 182 183 184 185 186 187	2895 2806 2475 2753 3327 3245 3039 2663	2907.2 2921.8 2937.0 2952.6 2968.3 2983.7 2983.7 2983.6 3012.8	-12.2 -115.8 -462.0 -199.6 358.7 261.3 -40.4
54567 555555 560 61	2604 3350 3112 2903 3171 2800 3013 2741	2927.6 2927.9 2928.3 2929.0 2929.6 2930.2 2930.4 2930.4	-323.6 422.1 183.7 -26.0 241.4 -130.2 82.6 -189.4	121 122 123 124 125 126 127	2999 2844 2957 3227 2969 2960 3239 3166	2922.1 2929.1 2936.7 2944.2 2951.0 2956.5 2950.3 2960.3	- 30 • 0 - 85 • 1 20 • 3 282 • 8 18 • 0 3 • 5 278 • 7 204 • 0	188 189 190 191 192 193 194	3071 3084 3212 3015 3367 2983 3172 3135	3026.1 3038.2 3049.1 3058.5 3066.3 3072.6 3077.2 3080.1	44.9 45.8 162.9 -43.5 300.7 -89.6 94.8
62 63 64 65 65	2991 3023 2629 2823 2863	2930.0 2929.4 2928.7 2928.0 2927.5	61.0 93.6 -299.7 -105.0 -64.5	129 130 131 132 133 133	3209 2886 2722 2524 2799	2961.5 2958.7 2953.7 2946.8 2938.5	247.5 -72.7 -231.7 -422.8 -139.5	196 197 198 199 200	3012 3108 2927 2974 2895	3081.4 3080.9 3078.9 3075.3 3070.2	-69.4 27.1 -151.9 -101.3 -175.2

TABLE 5-43. CONTINUED.

NUM OB	G COMP	RESID	NUM O	==== BS	COMP	RESID	N UM	OBS	COMP	RESID
	$\begin{array}{c} -1 \\ -1 \\ -1 \\ -1 \\ -1 \\ -1 \\ -1 \\ -1 $	$\begin{array}{c} -& -& -& -& -& -& -& -& -& -& -& -& -& $	$\begin{array}{c} 1 \\ 2 \\ 2 \\ 2 \\ 2 \\ 2 \\ 2 \\ 2 \\ 2 \\ 2 \\$	-54443333452544540945573328925388202324354364565581335=-68769511384452954455362216321897432897512845655813355=-		$\begin{array}{c}$	$\begin{array}{c} - & - & - & - & - & - & - & - & - & - $	$\begin{array}{c} -2222222222222222232232232232222222222$	$\begin{array}{c} 533,3,0\\ 2533,2,02\\ 2533,2,02\\ 2533,2,02\\ 2532,2$	$\begin{array}{c} -1 \\ -1 \\ -1 \\ -1 \\ -1 \\ -1 \\ -1 \\ -1 $

\$

TABLE 5-44 X01217: SUPPLEMENTAL STATISTICAL INFORMATION

VARIAN CE/CO-VARIAN CE MATRIX

DIAM	PREI	POST	TIME	VELO
9.193E ⁻ 16	5.253E 08	-8.800E ⁻ 08	-6.868E 09	3.992E ⁻ 11
5.253E 08	1.835E02	1.329E01	-1.093E01	4.225E ⁰ 2
-8.800E 08	1.329E01	4.297E02	-2.321E01	8.805E ⁻ 02
-6.868E 09	-1.093E01	-2.321E01	1.239E01	-4.798E ⁰ 2
3.992E ⁻ 11	4.225E ⁰ 2	8.805E ⁻ 02	-4.798E 02	2.552E04

CORRELATION MATRIX

	DIAM	PREI	POST	TIME	VELO
DIAM	1.000000	0.609252	-0.854026	0.443445	-0.435567
PREI	0.609252	1.000000	-0.113136	-0.372319	0.378780
POST	-0.854026	-0.113136	1.000000	-0.826759	0.821489
TIME	0.443445	-0.372319	-0.826759	1.000000	-0.999956
VELO	-0.435567	0.378780	0.821489	-0.999956	1.000000

NUMERICAL RANGES OF THE PARTIAL DERIVATIVES

	MAXIMUM	MIN IMUM
DIAM	1.063E09	-1.063E09
PREI	1.356E00	5.680E 03
POST	9.943E-01	-3.561E 01
TIME	1.522E01	-1.567E01
VELO	3.706E03	-3.855E03

......

only 3.15. The solution yielded an angular diameter of of 2.88 milliseconds of arc. However, the one sigma uncertainty in this value of 4.98 milliseconds of arc is quite high, and hence, the diameter must be viewed with caution. This large formal error was partially due to the low S+N/N ratio and partially to the small dynamic range of the change in signal level in comparison to the background level (only 10.5 percent of full scale).

The 300 millisecond extract from the raw intensity data used in the DC solution, the computed intensity values, and their residuals are given in Table 5-56. The accompanying supplemental statistical information is presented in Table 5-47. Figure 5-89 shows the usual partial derivatives of the solution intensity curve. As can be seen the regions of high sensitivity to variations of each of the parameters were well covered in the 300 millisecond solution.

The noise figure of the residual amplitudes is shown on the NOISEPLOT, Figure 5-90. The usual comparative power spectra are shown in Figure 5-91. Here, one may note that even at low frequencies, the background noise is dominant over the occultationcentered data. This was another contributing factor to the large statistical uncertainty in the angular diameter determination.

The Coordinated Universal Time of geometrical occultation was 00:02:07.790 (+/- 0.004 seconds).







Figure 5-90. NOISEPLOT of the occultation of X01217.



X01246

This relatively early (spectral class F0) star was occulted 78 minutes after X01217. In the intervening time, the seeing conditions improved considerably, and the "J" diaphragm was selected for this observation (as noted in the occultation summary, Table 5-45). Neither the trace of the raw intensity data (Figure 5-92) nor the integration plot (Figure 5-93) show any signs of stellar duplicity.

Three hundred milliseconds of data were used in the DC fitting process, and the best solution curve is shown in Figure 5-95. As it turned out, this was more data than was needed to be considered for a proper solution. To enable a better visualization of the fit, a detailed FITPLOT is presented as Figure 5-95. In this figure, covering only 100 milliseconds, 5-point smoothing was applied to the raw data before being plotted along with the fit to unsmoothed data.

The star was, not unexpectedly, indistinguishable from a point source. The 300 milliseconds of data subjected to DC fitting are listed (with the computed intensities and the residuals) on Table 5-46. Table 5-47 contains the variance/covariance and correlation matrices for the solution parameters, as well as the ranges of the numerical values of the partial derivatives of the intensity curve (which are shown in Figure 5-92).

TABLE 5-45 X01246: LUNAR OCCULTATION SUMMARY

STELLAR AND OBSERVING INFORMATION

Star: X01246, (SAO 109548, DM -00 0145) RA: 005610 DEC: +002154 mV: 8.2 Sp: F0 Filter: V Diaphragm: J Gain: B10+ Voltage: 1200

LUNAR INFORMATION

Surface Illumination:	88	percent
Elongation from Sun:	139	degrees
Altitude Above Horizon:	55	degrees
Lunar Limb Distance:	390238	Kilometers
Predicted Shadow Velocity:	648	.2 meters/sec.
Predicted Angular Rate:	0	.3426 arcsec/sec.

EVENT INFORMATION

Date: November 1	7, 1983	UT of Event:	01:20:02
USNO V/O Code:	53	HA of Event:	-210352
Position Angle:	66.1	Cusp Angle:	855
Contact Angle:	-15.2	Watts Angle:	88.5

MODEL PARAMETERS

Number of Data Points:	301
Number of Grid Points:	256
Number of Spectral Regions:	53
Width of Spectral Regions:	50 Angstroms
Limb Darkening Coefficient:	0.5
Effective Stellar Temperature:	7300 K

SOLUTIONS

Stellar Diameter (ams):	Point	Source
Time: (relative to Bin 0):	1920.4	(3.3)
Pre-Event Signal:	2716.5 (15.9)
Background Sky Level:	2488.9 (18.2)
Velocity (meters/sec.):	688.2 (80.1)
Lunar Limb Slope (degrees):	-9.82	(7.08)
U.T. of Occultation: 01:	20:01.938	(0.011)

PHOTOMETRIC NOISE INFORMATION

Sum-of-Squares of Residuals:	12365170
Sigma (Standard Error):	203.0202
Normalized Standard Error:	0.75862
Photometric (S+N)/N Ratio:	2.3182
(Change in Intensity)/Background:	0.10928
Change in Magnitude:	0.11261











Figure 5-95. FITPLOT of the occultation of X01246 showing the raw fit to a smoothed data set.

TABLE 5-46 X01246: OBSERVATIONS, COMPUTED VALUES, AND RESIDUALS FROM BIN 1750

TABLE 5-46. CONTINUED.

	====
A COMP REST NUM ORS COMP REST NUM ORS COMP REST	SID
NUM OBS COMP RESID NON OBS COMP RESID	
$ \begin{array}{c} 1 \\ 1 \\ 1 \\ 2 \\ 1 \\ 1 \\ 2 \\ 1 \\ 1 \\ 2 \\ 1 \\ 1$	
TABLE 5-47 X01246: SUPPLEMENTAL STATISTICAL INFORMATION

VARIAN CE/CO-VARIAN CE MATRIX

DIAM	PREI	POST	TIME	VELO
3.372E ⁻ 14	2.876E 07	-2.754E 07	-4.922E ⁻ 08	1.347E 09
2.876E 07	2.534E02	5.360 E00	-8.449E00	1.733E 01
-2.754E ⁻ 07	5.360 E00	3.316E02	-1.072E01	2.183E 01
-4.922E 08	-8.449E00	-1.072E01	1.079E01	-2.223E 01
1.347E ⁻ 09	1.733E ⁻ 01	2.183E 01	-2.223E 01	6.409E ⁻ 03

CORRELATION MATRIX

 DIAM
 PREI
 POST
 TIME
 VELO

 DIAM
 1.000000
 0.806934
 -0.727659
 -0.010259
 0.013761

 PREI
 0.806934
 1.000000
 -0.193646
 -0.481362
 0.482215

 POST
 -0.727659
 -0.193646
 1.000000
 -0.623821
 0.619762

 TIME
 -0.010259
 -0.481362
 -0.623821
 1.000000
 -0.999972

 VELO
 0.013761
 0.482215
 0.619762
 -0.999972
 1.000000

NUMERICAL RANGES OF THE PARTIAL DERIVATIVES

	MAXIMUM	MIN IMUM
DIAM	2.617E08	-2.596E08
PREI	1.367E00	6.770E 04
POST	9.993E-01	-3.667E ⁰¹
TIME	2.315E01	-2.468E01
VELO	1.092E03	-1.140E03









The steadiness of the sky, in terms of atmospheric transparency, is attested to by the NOISEPLOT, Figure 5-97. The mean of the residual distribution is only 0.8 percent below that of the computed intensities, with a one sigma width of 6.5 percent. Despite the prevailing observing conditions, the faintness of the stellar signal in comparison to the bright sky background (10.9 percent, or 5.5 percent of full scale) led to a very low S+N/N ratio of only 2.32. This in turn was the primary cause for the large statistical uncertainties in the time of geometrical occultation (3.3 milliseconds) and the L-rate (11.6 percent of the determined value of 688.2 meters/second).

The characteristics of the comparative power spectra (Figure 5-98) are very similar to those seen for the previous occultation. The change in scale, with contributing power components more important at somewhat higher frequencies, is due to the different time-scales of the two events.

The occultation disappearance was determined to have occurred at 01:20:01.938 (+/- 0.011 seconds), Coordinated Universal Time.

ZC0126

The occultation of ZC0126 followed that of X01246 by only eight minutes. This star was of later spectral type than X01246 (G5 in comparison to F0), as well as half a magnitude brighter in V. These two differences



in characteristics led to a somewhat better occultation signal, with a S+N/N ratio of 3.45 (as noted in Table 5-48).

The raw intensity data are presented in Figure 5-99. The integration plot, Figure 5-100, suggests a possible secondary disappearance at approximately 1200 milliseconds. To see this sudden change in the integrated signal level more clearly, Figure 5-101 details the supposed secondary disappearance only. As is apparent, there is indeed a sharp decline in the integrated signal level at approximately 1205 milliseconds. No other such abrupt changes in the signal level (except, of course at the time of the primary disappearance) are seen in the data. Hence, a fainter "wide" secondary component was detected at 1205 milliseconds.

The ascending and decending branches of the integrated data, presented in Figure 5-101, were fit via least squares to linear equations. The statistical uncertainty in the intersection of the two regression lines was 0.008 seconds, which is, therefore, the uncertainty in the time of secondary disappearance.

Four hundred milliseconds of intensity data centered on the estimated time of primary occultation were fit via the DC process. This, in retrospect, was roughly 100 milliseconds longer than needed in the post-occultation portion of the data. The best fit to

TABLE 5-48 ZC0126: LUNAR OCCULTATION SUMMARY

STELLAR AND OBSERVING INFORMATION

Star: ZC0126, (SAO 109552, DM -00 0146) RA: 005623 DEC: -001517 mV: 7.71 Sp: G5 Filter: V Diaphragm: J Gain: B10 Voltage: 1200

LUNAR INFORMATION

Surface Illumination:	88	percent
Elongation from Sun:	139	degrees
Altitude Above Horizon:	56	degrees
Lunar Limb Distance:	390149	kilometers
Predicted Shadow Velocity:	449	.8 meters/sec.
Predicted Angular Rate:	0	.2378 arcsec/sec.

EVENT INFORMATION

Date: November 1	7, 1983	UT of Event:	01:28:05
USNO V/O Code:	64	HA of Event:	-190552
Position Angle:	98.3	Cusp Angle:	535
Contact Angle:	-47.6	Watts Angle:	120.7

MODEL PARAMETERS

Number of Data Points: 401 Number of Grid Points: 266 Number of Spectral Regions: 53 Width of Spectral Regions: 50 Angstroms Limb Darkening Coefficient: 0.5 Effective Stellar Temperature: 5700 K

SOLUTIONS

 Stellar Diameter (ams):
 Point Source

 Time: (relative to Bin 0):
 982.4 (2.7)

 Pre-Event Signal:
 2721.6 (14.0)

 Background Sky Level:
 2279.2 (12.6)

 Velocity (meters/sec.):
 322.4 (14.3)

 Lunar Limb Slope (degrees):
 +22.11 (1.83)

 U.T. of Occultation:
 01:28:01.981 (0.010)

PHOTOMETRIC NOISE INFORMATION

Sum-of-Squares of Residuals:	13087870	
Sigma (Standard Error):	180.886	
Normalized Standard Error:	0.40882	
Photometric (S+N)/N Ratio:	3.4461	
(Change in Intensity)/Background:	0.19413	
Change in Magnitude:	0.19263	









Figure 5-101. INTPLOT of the ZC0126 secondary event.

the data is shown in Figure 5-102. A detail of the disappearance is plotted, along with the data subjected to 5-point smoothing, in Figure 5-103.

The solution found ZCO126 to have no detectable disc, and hence has been classed as a point source solution. The intensity data, computed intensities and residuals are compiled in Table 5-49. The usual supplemental statistics are presented on Table 5-50.

The PDPLOT, Figure 5-104 exhibits no anomalous behavior. It does indicate, however, that a better data sample might have been selected from data points beginning roughly 50 milliseconds earlier. But, given the point source nature of the solution and the fact that the regions of sensitivity to parametric variation are well covered, this selection would not have made a significant difference in the solution.

The NOISEPLOT and the POWERPLOT for this observation are presented as Figures 5-105 and 5-106, respectively. Here, too, there is no unusual behavior.

The Coordinated Universal Time of geometrical occultation for the primary event was 01:28:01.981 (+/- 0.010 seconds). The secondary event was seen at 01:28:02.204 (+/- 0.013 seconds) C. U. T.

<u>X01309</u>

The last occultation observed on the night of November 17, 1983, was that of the star X01309. The observational conditions at the time of disappearance of





Figure 5-103. FITPLOT of the occultation of ZC0126 showing the raw fit to a smoothed data set.

TABLE 5-49

2C0126: OBSERVATIONS, COMPUTED VALUES, AND RESIDUALS FROM BIN 800

N UM	OBS	COMP	RESID	NUM	OBS	COMP	RESID	N UM	OBS	COMP	RESID
0	2826	2724.8	101.2	67	2831	2730.9	100.1	134	2770	2808.0	-38.0
1	2883	2723.8	-150.8	68 69	2565	2722.6	-15/.6	135	2905	2824.3	594.7
3	2864	2722.7	141.3	70	3183	2705.7	477.3	137	2863	2851.9	11.1
5	2595	2720.9	-125.9	72	2738	2695.5	42.5	139	2721	2871.4	-150.4
67	2628	2720.8	-92.8	73	2657 3015	2694.8	-37.8	140	2886	2877.8	-166.0
8	2767	2722.7	44.3	75	3035	2703.2	331.8	142	2779	2883.8	-104.8
10	2303	2725.1	-422.1	27	2820	2722.0	98.0	144	2959	2880.7	78.3
11	2652	2725.6	-150.6	78	2705	2743.9	-38.9	145	2949	2869.2	411.1
13	2257	2724.3	-467.3	80 81	2567	2753.3	-186.3	147 148	3132 2851	2860.6	271.4
15	3103	2721.4	381.6	82	2971	2764.0	207.0	149	2936	2838.6	97.4
17	2573	2720.2	-147.2	84	2783	2760.6	22.4	151	2527	2811.2	-284.2
18 19	2434 2666	2720.8	-286.8	85 86	2565	2753.5	208.5	152	2619	2795.9	-160.8
20	2778	2723.7	54.3	87	2762	2731.4	30.6	154	2819 2597	2763.1	55.9 -148.8
22	3134	2726.1	407.9	89	2819	2704.8	114.2	156	2882	2728.2	153.8
24	2788	2725.4	62.6	91	2790	2682.4	107.6	158	2928	2692.5	235.5
25 26	3260 2923	2723.9	200.9	92 93	2579	2675.1	- 92.4	160	2666	2656.8	9.2
27 28	2339 2654	2720.5	-381.5	94 95	2564 2395	2671.4	-107.4	161 162	2535	2639.3	-104.3
29	2379	2719.6	-340.6	96	2343	2683.1	-340.1	163	2487	2605.1	-118.1
31	2831	2722.4	108.6	98	3108	2707.4	400.6	165	2505	2572.7	-67.7
32	2864	2726.4	137.6	100	2545	2738.2	-193.2	167	2385	2542.2	-157.2
34 35	2639 2456	2727.6	-88.6	101	2634 2879	2753.7	-119.7	168 169	2602 2720	252/.8	206.0
36	2657	2726.7	-69.7	103	2859	2780.1	78.9	170	2363 2503	2500.8	-137.8
38	2497	2722.2	-225.2	105	2591	2795.5	-204.5	172	2395	2476.2	-81.2
40	2587	2717.7	-130.7	107	2847	2796.3	50.7	174	2471	2453.8	17.2
41 42	2823	2716.8	186.6	108	2855	2791.1	24.6	175	2350	2443.5	-83.8
43 44	2543 2731	2719.4	-176.4	110	2724 2511	2770.6	-46.6	177	2263 2531	2424.5	-161.5
45	2774	2726.0	48.0	112	2775	2740.4	34.6	179	2833	2407.6	425.4
47	2636	2731.9	- 95.9	114	2581	2705.9	-124.9	181	2706	2392.6	313.4
40	2585	2732.3	-147.3	116	3066	2673.1	392.9	183	2316	2379.4	-63.4
50 51	2776	2729.7	46.3 -334.7	117	2881	2659.2	221.8	184	2335	23/3.3 2367.7	-32.7
52	2651	2720.8	-69.8	119	2389	2638.6	-249.6	186 187	2258 2307	2362.4	-104.4
54	2503	2712.1	-209.1	121	2403	2630.3	-227.3	188	2433	2352.8	80.2
56	2738	2709.4	28.6	123	2889	2635.3	253.7	190	2387	2344.3	42.7
58	2739	2715.7	23.3	125	2775	2652.8	122.2	191	2184	2337.0	-153.0
59 60	3043 2637	2/21.5 2728.1	321.5 -91.1	126 127	2936 2551	2665.6	-129.5	193 194	2456 2528	2333.6	197.5
61 62	2836 2600	2734.7	101.3	128	2777	2697.2	79.8	195 196	2135 2287	2327.6	-192.6
63	2815	2743.6	71.4	130	3059	2733.9	325.1	197	2496	2322.3	173.7
65	2844	2742.5	101.5	132	2629	2772.0	-143.0	199	2499	2317.7	181.3
	2/28	2/3/.8	-9.8	133	23/2	2/90.5	-210.0	200	2014	0.6162	198.4

TABLE 5-49. CONTINUED.

====			DECID	=====:			PECTD	NIIM	085	COMP	RESID
N UM	085	COMP	KESID				RESID				
201	2241	2313.7	-72.7	268	2366	2282.2	180.9	335	2180	2280.1	206.9
203	2183	2310.1	-127.1	270	2352	2282.1	69.9	337	2451	2280.1	170.9
204	2255	2308.5	-53.5	271	22231	2282.0	-62.0	339	2493	2280.1	212.9
206	2375	2305.6	69.4	273	2393	2281.9	-241.9	340	2164	2280.1	-116.1
208	2407	2303.1	103.9	275	2040	2281.8	-234.8	342	2079	2280.1	-201.1
209	2308	2301.9	-188.8	276	2251	2281.8	-30.8	343	2103	2280.1	-153.0
211	2265	2299.8	-34.8	278	2343	2281.7	61.3	345 346	2319	2280.0	39.0
213	2304	22 97 .9	6.1	280	2175	2281.6	-106.6	347	2173	2280.0	-107.0
214 215	2136 2235	22 97 .1 22 96 .3	-161.1	281	21308	2281.5	-144.5	349	2250	2280.0	-30.0
216	2287	2295.5	-8.5	283 284	2376	2281.4	94.6	350 351	2131 2164	2280.0	-149.0
218	2460	2294.1	165.9	285	2463	2281.3	181.7	352	2059	2280.0	-221.0
220	2243	2292.9	-49.9	287	2257	2281.2	-24.2	354	2211	2279.9	-68.9
221	2603	22 92 .3	310.7	288 289	2279	2281.2	-124.1	355	2381	2279.9	101.1
223	2417	22 91 .3	125.7	290	2259	2281.1	-22.1	357 358	2475	2279.9	195.1
225	2331	22 90 .3	40.7	292	2127	2281.1	-154.1	359	1921	2279.9	-358.9
226 227	2199	2289.9	-83.5	293	2238	2281.0	-43.0	361	22.96	2279.9	16.1
228	2207	2289.1	-82.1	295	2455	2281.0	-161.9	362 363	2512	2279.9	-216.9
230	2160	2288.4	-128.4	297	2299	2280.9	18.1	364	2003	2279.9	-276.9
232	2130	2287.7	-157.7	299	2047	2280.9	-233.9	366	2500	2279.8	220.2
233	2131	2287.4	-156.4	300	2295	2280.8	15.2	368	2135	2279.8	-144.8
235	2237	2286.8	-49.8	302	2295	2280.8	-37.7	369 370	2238	22/9.8	-41.8
237	2234	2286.3	-52.3	304	2616	2280.7	335.3	371	2581	2279.8	301.2
238	2412	2285.9	126.1	306	2431	2280.6	150.4	373	2393	2279.8	113.2
240	2459	2285.7	173.3	307 308	2360	2280.6	79.4 24.4	374	2528	2279.8	183.2
242	1859	2285.3	-426.3	309	2424	2280.6	143.4	376	2204	2279.8	-75.8
245	2582	2284.9	297.1	311	2236	2280.6	-44.6	378	2223	2279.8	-56.8
245 246	2307	2284.8	170.4	312	2399	2280.5	118.5	380	2480	2279.8	200.2
247	2151	2284.4	-133.4	314	2672	2280.5	391.5	381 382	2767 2166	2279.7	487.3
249	2058	2284.2	-226.2	316	2068	2280.4	-212.4	383	2208	2279.7	-71.7
251	2260	2283.9	-23.9	318	2361	2280.4	80.6	385	2451	2279.7	171.3
252 253	2040 2138	2283.8	-243.8	319 320	2496 2336	2280.4	215.6	386 387	2071	2279.7	-208.7
254	1965	2283.5	-318.5	321	2210	2280.3	-70.3	388	2047	2279.7	-232.7
256	2374	2283.3	90.7	323	2540	2280.3	259.7	3 90	2445	2279.7	165.3
257	22495	2283.2	-35.1	325	2353	2280.3	72.7	392	2195	2279.7	-84.7
259 260	2327	2283.0	44.0	326 327	2243	2280.3	-37.3	393 394	2246	2279.7	-132.7
261	2103	2282.8	-179.8	328	2123	2280.2	-157.2	3 95 3 96	2137 2291	2279.7	-142.7
263	2539	2282.6	256.4	330	1800	2280.2	-480.2	3 97	2185	2279.7	-94.7
265	2360	2282.4	77.6	332	2219	2280.2	-61.2	399	2099	2279.7	-180.7
266 267	2511 2116	2282.4	228.6	333 334	2295 2228	2280.2	-52.2	400	2227	22/9.6	-52.6
===	=====			=====	=====	======	========	=====		========	=======

TABLE 5-50 ZC0126: SUPPLEMENTAL STATISTICAL INFORMATION

VARIANCE/CO-VARIANCE MATRIX

DIAM	PREI	POST	TIME	VELO
3.719E ⁻ 15	1.224E ⁰⁷	-7.310E 08	-1.376E 08	8.217E ⁻ 11
1.224E_07	1.970E02	5.010E00	-8.764E00	4.015E_02
-7.310E 08	5.010E00	1.600E02	-6.918E00	3.135E_02
-1.376E 08	-8.764E00	-6.918E00	7.146E00	-3.278E02
8.217E ¹¹	4.015E ⁰ 2	3.135E ⁻ 02	-3.278E ⁻ 02	2.070E 04

CORRELATION MATRIX

	DIAM	PREI	POST	TIME	VELO
DIAM	1.000000	0.891094	-0.591950	-0.334484	0.338583
PREI	0.891094	1.000000	-0.167240	-0.663298	0.665761
POST	-0.591950	-0.167240	1.000000	-0.509287	0.504986
TIME	-0.334484	-0.663298	-0.509287	1.000000	-0.999987
VELO	0.338583	0.665761	0.504986	-0.999987	1.000000

NUMERICAL RANGES OF THE PARTIAL DERIVATIVES

	MAXIMUM	MIN IMUM
DIAM	4.982E08	-4.917E08
PREI	1.367E00	1.104E ⁰³
POST	9.989E 01	-3.665E01
TIME	1.793E01	-1.914E01
VELO	3.912E03	-3.984E03







Figure 5-105. NOISEPLOT of the occultation of ZC0126.



this KO star (mV=7.8) were virtually unchanged from those which prevailed at the times of the two previous occultations. As noted on Table 5-51, the instrumental configuration was the same as that used for the ZCO126 event.

The RAWPLOT and INTPLOT (Figures 5-107 and 5-108, respectively) typify those of a single star disappearance, with no indication of stellar duplicity. A DC fit to 200 milliseconds of data was performed, and the resulting best fit is depicted in Figure 5-109. No discernible stellar disc could be detected for this star.

The observations, computed intensities and the residuals for the solution are listed in Table 5-52. The PDPLOT, Figure 5-110, shows nothing unusual in the numerical partial derivatives. Table 5-32 contains the numerical ranges of these partial derivatives along with the variance/covariance and correlation matrices of the solution parameters.

As might have been expected both the NOISEPLOT and POWERPLOT for this event (Figures 5-111, and 5-112, respectively) are very similar to those seen for the observations of ZC0126 and X01246.

The Coordinated Universal Time of geometrical occultation was 03:32:06.631 (+/- 0.005 seconds).

TABLE 5-51 X01309: LUNAR OCCULTATION SUMMARY

______ STELLAR AND OBSERVING INFORMATION

_____ Star: X01309, (SAO 109577, DM +00 0159, GC1185K) RA: 005834 DEC: +004131 mV: 7.8 Sp: K0 Filter: V Diaphragm: J Gain: B10 Voltage: 1200

LUNAR INFORMATION

Surface Illumination:	88	percent
Elongation from Sun:	139	degrees
Altitude Above Horizon:	59	degrees
Lunar Limb Distance:	389597	kilometers
Predicted Shadow Velocity:	384	9 meters/sec.
Predicted Angular Rate:	0.	2038 arcsec/sec.

EVENT INFORMATION _____

Date: November	17, 1983	UT of Event:	03:32:06
USNO V/O Code:	64	HA of Event:	+112736
Position Angle:	104.1	Cusp Angle:	47S
Contact Angle:	-54.1	Watts Angle:	126.4

MODEL PARAMETERS

Number of Data Points:	201
Number of Grid Points:	256
Number of Spectral Regions:	53
Width of Spectral Regions:	50 Angstroms
Limb Darkening Coefficient:	0.5
Effective Stellar Temperature:	5100 K

SOLUTIONS _____

Stellar Diameter (ams):	Point	Source
Time: (relative to Bin 0):	1954.5	(2.0)
Pre-Event Signal:	3134.6	(18.6)
Background Sky Level:	2621.9	(23.8)
Velocity (meters/sec.):	514.6	(27.7)
Lunar Limb Slope (degrees):	-20.8	(4.13)
U.T. of Occultation: 03:3	32:06.631	(0.005)

PHOTOMETRIC NOISE INFORMATION

Sum-of-Squares of Residuals:	7797315
Sigma (Standard Error):	197.4502
Normalized Standard Error:	0.38512
Photometric (S+N)/N Ratio:	3.59662
(Change in Intensity)/Background:	0.19555
Change in Magnitude:	0.19392







TABLE 5-52 X01309: OBSERVATIONS, COMPUTED VALUES, AND RESIDUALS FROM BIN 1830

===:	=====							=====			
N UM	OBS	COMP	RESID	N UM	OBS	COMP	RESID	N UM	085	COMP	RESID
=M - 01234567890123456789012234567890123555555555555	$\begin{smallmatrix} & & & & & & & & & & & & & & & & & & &$	COMP 	$\begin{array}{c} \\ \textbf{RESID} \\ \textbf{RESID} \\ \textbf{RESID} \\ \textbf{10} \\ \textbf{2244}, \textbf{719} \\ \textbf{1767}, \textbf{224}, \textbf{719} \\ \textbf{1970}, \textbf{220} \\ \textbf{218}, \textbf{4370}, \textbf{200} \\ \textbf{218}, \textbf{4320}, \textbf{2310}, \textbf{218}, \textbf{4320}, \textbf{2310}, $	$ \begin{array}{l} = & = & = & = & = & = & = & = & = & = $	$\begin{array}{c} = & = & = & = & = & = & = & = & = & = $		$ \begin{array}{c} \text{Resc} \\ \text{Resc} $	$\begin{array}{c} & & & & & & \\ & & & & & & & \\ & & & & $	$ \begin{array}{c} 0 \\ = \\ 0 \\ 0$	COMP 	$\begin{array}{c} {\rm RESID} \\ {\rm result} \\$
45512345555555566666666666666666666666666666	29517 3203981 333281 3355597 3293746 3555947 33555947 33425594 3342590 32293 33555947 33293 32293 33293 3200 3200	31550.5 31553.3 31553.3 31427.5 31123.4 31105.5 31125.9 31105.5 311105.5 31155.9 31155	-2066.57 485.8 1853.07 467.74 -167.74 325.19 325.9 41.7.4 -1004.8 -1004.8	116 117 118 119 1221 1223 1244 1256 127 128 1290 1301 132	3068 32832 226505 22655 22635 2275 2275 2275 2275 2275 2275 2275 22	22884223 8892227774 8892227774223 8900 890 890 890 890 890 890 890 890 89	179.7 418.4 -103.4 -203.8 -211.8 -122.2 -33.4 216.6 91.0 91.0 91.0 91.0 215.1 158.0 130.4 -161.3	183 184 185 186 188 189 190 191 193 194 195 196 197 198 199	24030 227647 22599641 22599642 22599642 22599642 22592 224018 225754 225754 225754 225754 225754 225754 225754 225754 225754 225754 225754 225754 225754 225966 225967 225966 225967 225967 225967 225967 225967 225967 225967 22597 22597 22597 22597 22597 22597 22597 22597 22597 22597 22597 22597 22597 22597 22597 22597 22577 22597 225772 22577 225772 2257772 2257772 22577272 22577272 225772 225772 25	0,87,65,4,33,2,1,100,99,88 0,87,65,4,33,2,1,100,99,88 0,87,65,4,33,2,1,100,99,88 0,87,65,4,33,2,1,100,99,88 0,87,65,4,33,2,1,100,99,88 0,87,65,4,33,2,1,100,99,88 0,87,65,4,33,2,1,100,99,88 0,87,65,4,33,2,1,100,99,88 0,87,65,4,33,2,1,100,99,88 0,87,65,4,33,2,1,100,99,88 0,87,65,4,33,2,1,100,99,88 0,87,65,4,33,2,1,100,99,88 0,87,65,4,33,2,1,100,99,88 0,87,65,4,33,2,1,100,99,88 0,87,65,4,33,2,1,100,99,88 0,87,65,4,33,2,1,100,99,88 0,87,65,4,33,2,1,100,99,88 0,87,65,4,33,2,1,100,99,88 0,87,65,4,33,2,1,100,99,88 0,87,65,4,4,4,4,4,4,4,4,4,4,4,4,4,4,4,4,4,4,	-221 135 -277 -63 -300 -28 -390 -172 -223 -76 -49 0 -469 -216 -246 -246 -246 -247 -295
_66	3312	3106.8	205.2	133	2615	2673.4	-58.4	200	2511	2623.7	-112.7



Figure 5-110. PDPLOT of the occultation of X01309.



Figure 5-111. NOISEPLOT of the occultation of X01309.

TABLE 5-53 X01309: SUPPLEMENTAL STATISTICAL INFORMATION

VARIAN CE/CO-VARIAN CE MATRIX

DTAM	PREI	POST	TIME	VELO
3.004E ⁻ 14	4.253E 07	-5.242E ⁻ 07	-2.350E 08	3.963E_10
4.253E 07	3.435E02	1.772E01	-8.716E00	1.041E_01
-5.242E07	1.772E01	5.679E02	-1.309E01	1.504E_01
-2.350E 08	-8.716E00	-1.309E01	4.109E00	-4.807E_02
3.963E-10	1.041E ⁰¹	1.504E ⁰¹	-4.807E ⁰ 2	7.689E 04

CORRELATION MATRIX

	DIAM	PREI	POST	TIME	VELO
DIAM	1.000000	0.707272	-0.807680	0.275065	-0.255759
PREI	0.707272	1.000000	-0.155474	-0.463281	0.480530
POST	-0.807680	-0.155474	1.000000	-0.779561	0.767024
TIME	0.275065	-0.463281	-0.779561	1.000000	-0.999799
VELO	-0.255759	0.480530	0.767024	-0.999/99	1.000000

NUMERICAL RANGES OF THE PARTIAL DERIVATIVES

	MAXIMUM	MINIMUM
DIAM	2.403E08	-2.322E08
PREI POST	9.965E 01	-3.667 E 01
TIME VELO	3.313E01 2.824E03	-3.536E01 -2.852E03



ZC3158 (37 Capricorni)

The bright (mV=5.79) F5 star 2C3158 was occulted under very favorable circumstances on the night of December 10, 1983. Though the event occurred early in the evening (approximately 00:07 U. T.), astronomical twilight had been over for approximately ten minutes. Thus, despite the relatively small solar elongation of 62 degrees, the event was seen in dark skies. The penalty for this was a moderately low lunar altitude, though still above two air masses. The darkened portion of the 27 percent illuminated lunar disc was quite dark, and virtually no Earthshine could be seen, even in a 6-inch finding telescope.

The seeing was unusually good for RHO and estimated at 2 seconds of arc or better. These almost extraordinary combinations of fortuitous circumstances allowed the observation to be made with the narrow bandwidth "b" filter. Though the integrated bandpass of this filter is less than 1/10 that of a Johnson B filter, the S+N/N ratio obtained for this observation (given in Table 5-54) was still a respectable 5.52.

As easily seen on the RAWPLOT, Figure 5-113, the stellar signal was highly dominant (by a factor of 6.5) over the background sky brightness. The integration plot of the 4096 milliseconds of raw data (Figure 5-114) shows no evidence of any secondary disappearances.

TABLE 5-54 ZC3158: LUNAR OCCULTATION SUMMARY

STELLAR AND OBSERVING INFORMATION

Star: 2C3158, (37 Cap, SAO 190461, DM -20 6237) RA: 213355 DEC: -200935 mV: 5.79 Sp: F5 Filter: "b" Diaphragm: I Gain: C9+ Voltage: 1200

LUNAR INFORMATION

Surface Illumination:	27	percent		
Elongation from Sun:	62	degrees		
Altitude Above Horizon:	31	degrees		
Lunar Limb Distance:	400741	kilometers		
Predicted Shadow Velocity:	355	.7 meters/sec.		
Predicted Angular Rate:	0	.1831 arcsec/sec.		

EVENT INFORMATION

Date: December	10, 1983	UT of Event:	00:06:42
USNO V/O Code:	59	HA of Event:	
Position Angle:	118.3	Cusp Angle:	46S
Contact Angle:		Watts Angle:	137.9

MODEL PARAMETERS

Number of Data Points: 301 Number of Grid Points: 256

Number of Spectral Regions:	5
Width of Spectral Regions:	50 Angstroms
Limb Darkening Coefficient:	0.5
Effective Stellar Temperature:	6600K

SOLUTIONS

Stellar Diameter (ams):	1.74	(1.68)
Time: (relative to Bin 0):	1521.9	(1.6)
Pre-Event Signal:	1509.9	(23.7)
Background Sky Level:	231.7	(49.3)
Velocity (meters/sec.):	269.9	(3.34)
Lunar Limb Slope (degrees):	+20.3	(0.5)
U.T. of Occultation: 00:	:06:39.731	(0.004)

PHOTOMETRIC NOISE INFORMATION

Sum-of-Squares of Residuals:	38293450
Sigma (Standard Error):	357.274
Normalized Standard Error:	0.2795
Photometric (S+N)/N Ratio:	4.5775
(Change in Intensity)/Background:	5.5159
Change in Magnitude:	2.0349
Normalized Standard Error: Photometric (S+N)/N Ratio: (Change in Intensity)/Background: Change in Magnitude:	0.2795 4.5775 5.5159 2.0349





Because of the somewhat slower R-rate for this event (predicted to be 0.1831 arc seconds per second), 300 milliseconds of data were used in the DC fitting process. The resulting best fit is shown graphically in Figure 5-115 and as can be seen is rather good. The slowly decaying diffraction fringes arose from the near monochromaticity of the optical bandpass. This is reflected in the PDPLOT (Figure 5-116) as well. Examination of the PDPLOT (see Table 5-56 for the ranges of numerical values in the partial derivatives) shows that the sensitivity of the solution curve to variations in the stellar diameter and velocity parameter (and to a lesser extent time of geometrical occultation) continues to be high before the start of the choosen data set. While this is true, because of the decaying amplitude of the diffraction fringes in the occultation curve itself, to include data from times earlier than considered would not be helpful. Data extracted from times before about 1280 milliseconds would be dominated by variations due to scintillation noise rather than intrinsic variations in the intensity curve.

A stellar angular diameter of 1.74 (+/- 1.68) milliseconds of arc was determined by the DC fitting procedure, for ZC3158. This star, according to Schlesinger and Jenkins (1940), has a parallax of 0.032 arc seconds, or a corresponding distance of 31 parsecs.





Figure 5-116. PDPLOT of the occultation of ZC3158.



TABLE 5-55 ZC 3158: OBSERVATIONS, COMPUTED VALUES, AND RESIDUALS FROM BIN 1280 _____ NUM OBS COMP RESID NUM OBS COMP RESID NUM OBS COMP RESID -----------1332.2 1323.6 1324.7 1335.6 1355.4 1355.4 1383.6 1418.4 1418.4 1458.2 1501.4 1546.0 1590.0 1631.6 1669.3 1447 1367 1955 2049 2083 1462 210.4 -190.2 24.5 555.0 534.9 179.6 1677 1300 1544 67 1671 1491.4 134 114.8 0 1466.6 -207.9 -282.2 324.5 345.2 279.3 43.4 630.3 713.4 727.6 78.4 1455.9 1427.2 1410.5 135 136 1490.2 1519.5 1548.0 68 1248 1145 23 69 1735 1752 1696 žó 2103 15756.9 15766.9 155766.9 155766.9 15576.9 1485.9 1485.9 1485.9 14451.6 1459.5 14496.7 14596.7 14596.8 15576.8 15576.8 1406.8 138 4 71 72 73 74 75 77 77 -86.3 1416.7 139 5 -314.4 -448.7 -125.7 494.0 140 141 142 991 1439.7 1104 1439.7 1471.7 1510.0 1548.9 1058 1121 1121 963 1346 2004 -400.2 1232 1087 1258 - 380 ġ. •4 196.1 185.5 -233.7 -365.5 15.8 -425 143 144 145 .0 1745 1770 ā 1584.5 .0 -288 1210 78 1378 •6 1611.7 1628.5 1631.2 1621.2 1219 79 1263 1647 8Ó 210.8 1443 1791 1775 81 82 1832 1980 14 1980 1598.6 1280 1566.3 1479 1528.2 -286.3 -49.2 781.9 1420.6 83 1280 1560.5 1479 1528.2 2270 1488.1 2871 1450.4 1759 1419.3 1170 1398.1 84 576.8 582.6 573.5 551.5 520.7 488.1 460.6 443.9 2185 1923 856 422 85 86 18 1450.4 19 87 339.7 -228.1 265.4 -81.5 -494.5 ***** 88 1388.6 1392.5 1408.5 1799 1122 1747 310.9 -338.6 303.1 89 1654 90 1311 89 9ĭ 914 1595 1717 1023 153.6 263.1 -455.7 -430.0 1441 4 1453 9 1478 7 1511 0 1544 2 1571 9 1588 4 1590 7 1578 0 1552 9 1519 9 1485 8 1435 8 1435 8 92 93 9Ô6 1435.8 1471.2 -529.8 92 906 93 2183 94 1641 95 2079 96 2276 97 1671 98 2307 99 1740 100 1384 1592 1519 1511.2 1552.1 1590.1 1621.4 129.8 161 526.9 1604 1081 162 163 1463 -81.2 -580.9 -696.4 163 1585 164 1471 165 898 49.6 663.7 86.1 991 892 1643.3 1081 -509.7 1653.9 1652.0 166 167 1591 1560 -509.7 41.0 65.1 -192.9 777.5 272.2 341.4 -322.2 434.2 434.2 459.5 -268.0 100 1619 1618 1327 2263 1728 1778 1109 1875 1923 1455 1843 1071 1638.1 168 1485 169 1795 170 938 171 1186 172 1003 173 1037 174 1172 175 1675 176 1423 178 981 178 01278 180 1278 180 1278 181 1328 182 1733 183 1521 184 2471 185 1651 101 1744 168 1484 105.9 389.8 211.4 263.6 509.4 614.2 2003 1613.2 36 37 1791 1579.6 104 1804 1540.4 2008 105 1498.6 38 1436.6 1431.2 1440.8 1463.7 1495.5 1531.1 1564.1 1564.1 1589.1 žă 106 2072 1457.8 -592.4 1421.4 829 1392.2 1372.4 1363.7 1366.5 -603.2 -183.4 -103.7 41 108 789 109 1189 110 1260 42 -40.5 43 110 111 -103.7 -416.5 -241.9 -182.6 21.3 -29.8 140.2 -423.7 213.7 213.7 742.0 302.7 -186.7 -32.5 -5 158 4 -297.4 0 -295.0 65.0 91.6 -493.1 950 1139 44 1366.5 1380.9 1405.6 1438.7 1477.8 1520.2 1562.8 1602.7 1637.3 1664.0 586.9 -58.2 32.0 45 2176 1543 1601.2 1599.0 46 1223 47 1631 114 1460 1599.0 1582.5 1554.6 1519.5 1483.3 1451.4 1429.4 1420.0 1425.6 1444.7 1460 1448 1567 1703 1179 1851 1633 2424 208.5 -159.6 -56.5 -50.3 48 1791 115 116 117 118 119 120 121 122 4ğ 1395 1463 1433 1536 1355 1435 50 51 52 53 53 54 55 -50.3 84.6 -74.4 15.0 2.4 186 187 1511 1368 1832-1 -321 1865-4 -497-4 1894-9 -535-9 1920-2 -479-2 188 1359 188 1359 189 1441 190 1763 191 1779 192 2431 193 1928 194 1935 195 2095 196 2338 197 1985 198 2628 199 1791 200 3150 1428 1992 2.4 336.3 52.5 33.2 -774.0 -417.6 746.2 81.2 123 124 125 1941.1 1957.4 -178 56 1781 1499 1474.5 1510.8 1548.0 1580.6 -100.7 -32.5 -415.7 -174.9 -109.9 -184.9 -317.4 1527 1544 57 1639 •4 1232 1441 969.1 976.1 461 .9 774 59 126 127 128 .1 978.6 60 1163 2350 1468 -43.6 1978.6 1976.5 1969.9 1959.1 1944.4 1925.8 1903.7 1603.8 1351 1175 118.5 1696 1614.8 129 1920 1570 1780 1535 1611.0 1594.4 1565.5 1529.9 130 131 132 133 879 1469 1327 1431 1492.4 1449.9 1410.7 1376.9 1350.2 63 309.0 -570.9 58.3 -49.9 25 g -24.4 214.5 5.1 •6 64 -134 .8 65 80.8 1246.3

TABLE 5-55. CONTINUED.

TABLE 5-56

ZC 3158: SUPPLEMENTAL STATISTICAL INFORMATION

VARIANCE/CO-VARIANCE MATRIX

DIAM	PREI	POST	TIME	VELO
4.209E ⁻ 16	1.051E-07	-1.973E 07	-4.960E ⁻ 10	2.342E[12
1.051E ⁰⁷	8.947E02	8.525E01	-1.498E01	3.970E_02
-1.973E 07	8.525E01	2.254E03	-3.618E01	1.069E_01
-4.960E ⁻ 10	-1.498E01	-3.618E01	3.897E00	-1.164E_02
2.342E ⁻ 12	3.970 = 02	1.069E ⁻ 01	-1.164E ⁻ 02	4.462E 05

CORRELATION MATRIX

	DIAM	PREI	POST	TIME	VELO
DIAM	1.000000	0.548987	-0.891538	0.642859	-0.673466
PREI	0.548987	1.000000	-0.110924	-0.283453	0.243923
POST	-0.891538	-0.110924	1.000000	-0.918646	0.933638
TIME	0.642859	-0.283453	-0.918646	1.000000	-0.999159
VELO	-0.673466	0.243923	0.933638	-0.999159	1.000000

NUMERICAL RANGES OF THE PARTIAL DERIVATIVES

	MAXIMUM	MIN IMUM
DIAM	4.690E09	-4.917E09
PREI	1.367E00	2.017E 02
POST	9.798E ⁰¹	-3.667E ⁻ 01
TIME	4.488E01	-4.674E01
VELO	2.702E04	-2.786E04

This then gives a linear diameter for the star of approximately 6 solar radii. This is entirely consistent with the spectral and luminosity characteristics of 2C3158. Again, assuming a distance of 31 parsecs, an F5 star would have an absolute V magnitude of -0.4. This value, according to Allen (1963), would place it between luminosity classes II and III. Also, interpolating from Allen's tables, an F5 star of this luminosity class would have a diameter of roughly 8 solar diameters. Hence, a "ballpark" estimate of the anticipated diameter of the star shows the observationally determined angular diameter to be in good agreement with what might have been inferred from the previously known physical properties of ZC3158.

The NOISEPLOT of the raw data is presented as Figure 5-117. As can be seen in the RAWPLOT itself, the dispersion of the residual amplitudes about the mean intensity level is quite high. This is to be expected for photometric data whose dominant noise source is, as in this case, determined by photon arrival statistics.

The POWERPLOT of the comparative power spectra (Figure 5-118) does indeed show that the background sky level had no significant effect on the observations. The flat nature of the star-plus-sky power spectrum is typical of photon shot noise. The only low frequency components in the occultation power spectrum of any significance are due to diffraction fringing effects,



with no contributions from atmospheric scintillation or transparency variations.

The Coordinated Universal Time of geometrical occultation was 00:06:42.333 (0.004).

ZC0835

The night of March 11, 1984 U. T., gave rise to four occultations which were observed from RHO. The first of these, that of 2C0835, occurred at approximately 00:30 U. T. and hence was observed 25 minutes before the end of astronomical twilight. Though the sky was not yet completely dark, the brightness of this star (mV=6.92) led to a S+N/N ratio of 11.81 with a Johnson B filter. The seeing was quite good and enabled the use of the small "J" diaphragm, as noted in the occultation summary (Table 5-57).

The photoelectric record of the event is presented in Figure 5-119. The rather clean signal and precipitous drop are attributed to the good seeing, small background noise, and the brightness of the star. The integration plot (Figure 5-120) is remarkably clean and attests to the steadiness of the sky over time-scales of tenths of seconds to seconds. There is no indication of stellar duplicity in this figure.

The parametric solution obtained for the 200 milliseconds of observational data fit by the DC procedure, given in Table 5-58, did give rise to a value for the stellar diameter. However, the diameter determined is very close to the lower limit of detectability due to both the finite

TABLE 5-57 ZC0835: LUNAR OCCULTATION SUMMARY

STELLAR AND OBSERVING INFORMATION

Star: ZC0835, (SAO 077252, DM +24 0854) RA: 053232 DEC: +243711 MV: 6.92 Sp: B8 Filter: B Diaphragm: J Gain: B9 Voltage: 1200

LUNAR INFORMATION

Surface Illumination:	53 percent
Elongation from Sun:	93 degrees
Altitude Above Horizon:	77 degrees
Lunar Limb Distance:	374852 kilometers
Predicted Shadow Velocity:	385.64 meters/sec.
Predicted Angular Rate:	0.2345 arcsec/sec.

EVENT INFORMATION

Date: March 11,	1984	UT of Event:	00:39:30
USNO V/O Code:	17	HA of Event:	+130138
Position Angle:	33.7	Cusp Angle:	36N
Contact Angle:	+47.3	Watts Angle:	35.2

MODEL PARAMETERS

Number of Data Points:	201
Number of Grid Points:	256
Number of Spectral Regions:	42
Width of Spectral Regions:	50 Angstroms
Limb Darkening Coefficient:	0.5
Effective Stellar Temperature:	12000 K

SOLUTIONS

Stellar Diameter (ams):	1.08 (1.46)
Time: (relative to Bin 0):	2238.1 (0.6)
Pre-Event Signal:	2441.7 (15.0)
Background Sky Level:	791.1 (18.0)
Velocity (meters/sec.):	388.0 (5.5)
Lunar Limb Slope (degrees):	-3.15 (0.82)
U.T. of Occultation: 00:3	39:28.8 (0.1)

PHOTOMETRIC NOISE INFORMATION

	Sum-of-Squares of Residuals:	4660483			
	Sigma (Standard Error):	153.6513			
	Normalized Standard Error:	0.092524			
	Photometric (S+N)/N Ratio:	11.80804			
	Intensity Change/Background:	2.083527			
	Change in Magnitude:	1.222619			
	Sigma (Standard Error): Normalized Standard Error: Photometric (S+N)/N Ratio: Intensity Change/Background: Change in Magnitude:	0.092524 11.80804 2.083527 1.222619			




TABLE 5-58 ZCO835: OBSERVATIONS, COMPUTED VALUES, AND RESIDUALS FROM BIN 2120

====							DEGID	======	ORC	COMP	DECTD
NUM	OBS	COMP	RESID	N UM	OBS	COMP	KESID	N 0M		COPIF	KE91D
0 1	2488 2481	2443.7 2451.3	44.3 29.7	67 68	2487 2561	2476.3	10.7	134 135	857 957	898.2	-41.2
23	2403	245/./	-156.4	70	2398	2236.0	236.0	137	839	878.0	-39.0
4	2330	2455.3	-125.3	71	2540	2176.4	363.6	138	843	872.3	-29.3
26	2145	2438.4	-293.4	73	2191	2115.5	75.5	140	885	862.4	22.6
7	1923	2433.4	-510.4	74	2103	2118.4	-15.4	141 142	1010	858.0	-86.0
9	2021	2441.3	-420.3	76	2143	2189.8	-46.8	143	879	850.3	28.7
10	2163	2451.0	-288.0	78	2280	2253.9	-52.4	144	807	843.7	-36.7
12	2372	2462.8	-90.8	79	2463	2421.2	41.8	146	784	840.8	-56.8
14	2485	2451.7	33.3	81	2861	2611.8	249.2	148	871	835.5	35.5
15 16	2341 2224	2442.2	-101.2	82	2769	2791.6	2.4	150	874	831.0	43.0
17	2219	2434.1	-215.1	84	2782	2868.5	-86.5	151	767	829.0	-62.0
19	2472	2448.4	23.6	86	2832	2983.6	-151.6	153	712	825.3	-113.3
20	2718	2457.9	260.1	87	2905	3018.6	-113.6	154	816	823.7	-6.2
22	2547	2462.6	84.4	89	3040	3040.6	-0.6	156	889	820.7	68.3
23	26 90	2454.0	247.0	90	2997	3001.1	-4.1	158	952	818.2	133.8
25	2651	2431.4	219.6	92 93	2978	2960.6	17.4	159 160	1036	817.0	219.0
27	2672	2426.6	245.4	94	3071	2845.3	225.7	161	665	814.9	-149.9
28	2280	2437.4	-15/.4 44.3	95 96	2846	2695.1	93.9	162	786	813.1	-27.1
30	2520	2473.8	46.2	97 98	2899	2611.1	287.9	164	789	812.3	-23.3
32	2638	2493.6	144.4	99	2531	2433.0	98.0	166	840	810.8	29.2
33 34	2419	2486.3	-67.3	100	2160	2341.6	-181.6	167	768	810.1	-41.4
35	2443	2439.9	3.1	102	2100	2159.8	-59.8	169	776	808.8	-32.8
37	2549	2387.3	161.7	104	1895	1985.1	-90.1	171	861	807.6	53.4
38	2215	2376.0	-161.0	105	1768	1902.2	-134.2	173	888	807.0	81.5
40	2467	2402.2	64.8	107	1865	1747.0	118.0	174	1001	805.9	195.1
41	2630	2479.5	150.5	109	1746	1607.7	138.3	176	820	805.0	15.0
43	2823	2520.9	302.1	110	1529	1544.3	-15.3	177	839 912	804.5	34.5
45	2647	2568.8	78.2	112	1206	1429.8	-223.8	179	802	803.7	-1.7
47	2627	2539.0	88.0	114	1419	1331.0	88.0	181	680	803.0	-123.0
48	2419	2496.1	-77.1	115	1546 1368	1287.1	258.9	182 183	799	802./	-153.4
50	2135	2384.8	-249.8	117	1238	1209.5	28.5	184	703	802.1	-99.1
31 52	2140	2296:9	-156.9	119	1186	1144:0	42.0	182	726	801.6	-75:6
53	2240	2280.5	-40.5	120	1021	1115.3	-94.3	187	806 718	801.3	-83.0
55	2060	2318.7	-258.7	122	974	1065.0	- 91.0	189	858	800.8	57.2
57	22490	2435.6	-188.6	124	1001	1023.1	-22.1	1 91	723	800.3	-77.3
58	2193	2508.0	-315.0	125	966 954	1004.9	-38.9	1 92	665 823	800.0	-135.0
60	2928	2638.8	289.2	127	867	973.1	-106.1	194	953	799.6	153.4
62	2780	2703.6	76.4	128	940	946.6	-0.2	195	714	799.2	-85.2
63 64	2599	2700.5	-101.5	130	1113	935.1 924.6	177.9	197	838 746	799.0	39.0
65	2411	2623.0	-212.0	132	1056	915.0	141.0	199	813	7 98 .7	14.2
===	2310	C+CCC2	-37.5	222	500	900.2	0+ 1 C	200		/ 90.0	-0/.0

aperture of the telescope and the optical bandpass. This is reflected in the formal error of the determined angular diameter which was 1.08 (+/- 1.46) milliseconds of arc.

The low noise level, in comparison to the stellar signal intensity, resulted in a very good fit to the observational data. This can be seen in the FITPLOT, Figure 5-121. The PDPLOT of the solution curve, Figure 5-122, indicates that the regions of high sensitivity to parametric variation were well considered in the solution. The numerical ranges of the partial derivatives are listed, along with the other usual supplementary statistics, in Table 5-69.

The distribution function of the residual intensities is shown in the NOISEPLOT, Figure 5-21. The purely Gaussian nature of the distribution bespeaks the dominance of the stellar signal over the unimportant sky background.

The comparative power spectra, shown in Figure 5-124. are additional indicators of the quality of this observation. The frequency components of the occultation power spectrum dominates those of the star-plus-sky signal to frequencies in excess of 100 Hertz.

The determination of the Coordinated Universal Time of geometrical occultation to high precision was hampered by the inability to successfully decode the WWVB time signal, and as previously noted, this can often be a problem until well after sunset. As a result, a post-occultation calibration of the SPICA-IV/LODAS clock, nearly an hour after the time of the event, degraded the quality of the event timing. The





Figure 5-122. PDPLOT of the occultation of ZC0835.



VARIANCE/CO-VARIANCE MATRIX ------

DIAM	PREI	POST	TIME	VELO
5.021E ⁻ 17	1.821E 08	-1.885E 08	-3.245E ⁻ 10	3.783E ⁻ 12
1.821E 08	2.259E02	1.218E01	-2.286E00	1.933E_02
-1.885E 08	1.218E01	3.239E02	-2.984E00	2.444E_02
-3.245E ¹⁰	-2.286E00	-2.984E00	3.202E ⁰¹	-2.679E 03
3.783E ⁻ 12	1.933E 02	2.444E ⁰ 2	-2.679E 03	3.033E 05

CORRELATION MATRIX -----

	DIAM	PREI	POST	TIME	VELO
DIAM	1.000000	0.749272	-0.771171	0.169356	-0.150545
PREI	0.749272	1.000000	-0.156314	-0.523365	0.539440
POST	-0.771171	-0.156314	1.000000	-0.756773	0.744181
TIME	0.169356	-0.523365	-0.756773	1.000000	-0.999818
VELU	-0.150545	0.539440	0.744181	-0.999818	1.000000

NUMERICAL RANGES OF THE PARTIAL DERIVATIVES

	MAXIMUM	MINIMUM	
DIAM	4.426E09	-4.589E09	
PREI	1.363E00	4.073E 03	
POST	9.959E ⁰¹	-3.630E 01	
TIME	9.160E01	-9.600E01	
VELO	1.081E04	-1.121E04	
			= =



time of geometrical occultation was determined to be 00:39:28.8 (+/- 0.1 seconds) C. U. T.

X07145

The moderately bright (mV=7.4) K0 star, X07145, was occulted 49 minutes after the ZC0835 event. The good photometric conditions which were present at the time of the earlier event continued to prevail, as they did for the two events which were to follow. As a result, the plot of the raw photometric data (Figure 5-119) was extremely similar to that of ZC0835. The S+N/N ratio was not as good, since this star was roughly one half of a magnitude fainter, but was still a very substantial 7.74.

This observation was made with a Johnson V filter (as noted in Table 5-60) in accord with the spectral type of the star. The integration plot of the event (Figure 5-126) was practically free of complicating effects, which sometimes arise under noisy conditions. The linearity of the ascending and descending branches of the plotted integrated intensity function precludes the possibility of "wide" stellar duplicity.

The 200 milliseconds of data used in the DC fitting process are listed in Table 5-61, along with the computed intensities and their residuals. The FITPLOT of the observed and computed intensity curves is presented as Figure 5-127. The determined L-rate for this event of 528.9 (+/- 11.3) meters per second caused the time-scale of this event to be 56 percent faster than that of the 200835 event. Thus, there





TABLE 5-60 X07145: LUNAR OCCULTATION SUMMARY

_____ -----STELLAR AND OBSERVING INFORMATION _____ Star: X07145 (SA0 077276. DM +24 0868) RA: 053404 DEC: +241655 mV: 7.4 Sp: K0 Filter: V Diaphram: J Gain: B11+ Voltage: 1100 LUNAR INFORMATION _____ 53 percent Surface Illumination: Elongation from Sun: 93 degrees 67 degrees Altitude Above Horizon: Lunar Limb Distance: 374996 kilometers Predicted Shadow Velocity: 494.14 meters/sec. Predicted Angular Rate: 0.2718 arcsec/se 0.2718 arcsec/sec. EVENT INFORMATION

Date: March 11.	1984	UT of Event:	01:28:06
USNO V/O Code:	17	HA of Event:	+245103
Position Angle:	125.2	Cusp Angle:	52S
Contact Angle:	-40.7	Watts Angle:	126.6

MODEL PARAMETERS

Number of Data Points:	201
Number of Grid Points:	256
Number of Spectral Regions:	53
Width of Spectral Regions:	50 Angstroms
Limb Darkening Coefficient:	0.5
Effective Stellar Temperature:	5100K

SOLUTIONS

 Stellar Diameter (ams):
 Point Source

 Time: (relative to Bin 0):
 2138.8
 (0.8)

 Pre-Event Signal:
 1547.2
 (13.0)

 Background Sky Level:
 658.7
 (14.8)

 Velocity (meters/sec.):
 528.9
 (11.3)

 Lunar Limb Slope (degrees):
 -10.4
 (1.3)

 U.T. of Occultation:
 01:28:05.604 (0.006)
 0.006)

PHOTOMETRIC NOISE INFORMATION

Sum-of-Squares of Residuals:	3497400
Sigma (Standard Error):	131.9
Normalized Standard Error:	0.14844
Photometric (S+N)/N Ratio:	7.7369
Intensity Change/Background:	1.3491
Change in Magnitude:	0.92725

TABLE 5-61 X07145: OBSERVATIONS, COMPUTED VALUES, AND RESIDUALS FROM BIN 2025

====				=======		COMP	DESTD	NIIM	OBS	COMP	RESID
NUM	OBS	COMP	RESID	NOPI		COMP					
0	1885 1539	1549.2 1547.3	335.8	67 68	1438 1504	1693.5 1699.3	-255.5	134	735 587	688.0	-101.0
23	1609	1547.7	61.3	69 70	16/9	1684.0	278.3	130	677	684.1	-7.1
4	1791	1553.1	237.9	71	1881	1601.1	279.9	138 139	628 655	682.5	-25.9
é	1520	1550.0	-30.0	73	1673	1486.7	186.3	140	672	679.5 678.2	-7.5 140.8
8	1718	1545.8	172.2	75	1446	1394.7	51.3	142	597	677.0	-80.0
9 10	1574 1671	1548.6	25.4 118.2	77	1232	1364.2	-132.2	145	739	674.9	64.1
$\frac{11}{12}$	1347 1256	1555.0	-208.0 -297.1	78 79	1299 1401	13/6./	-77.7	145	519	673.1	-154.1
13	1626	1548.6	77.4	80 81	1676 1560	1450.8	225.2 54.0	14/ 148	6/5 712	671.6	40.4
15	1447	1545.4	-98.4	82	1484	1567.9	-83.9	149 150	715 787	671.0 670.4	44.0 116.6
17	1701	1554.3	146.7	84	1433	1694.0	-261.0	151	759	669.8	89.2 70.7
19	1359	1553.4	-194.4	86	2201	1798.3	402.7	153	574	668.8	-94.8
20 21	1387	1548.0	-207.0	88	1956	1860.0	96.0	155	574	667.9	-93.9
22 23	1458 1331	1544.5	-86.5	89 90	1918	1870.9	47.1	157	828	667.1	160.9
24	1555	1555.8	-0.8	91 92	2015 1937	1857.8	157.2	158	705	666.4	38.6
26	1477	1556.2	-79.2	93	1936	1799.4	136.6	160 161	507 558	666.1 665.8	-159.1
28	1520	1541.3	-21.3	95	1783	1708.2	74.8	162	591 749	665.5	-74.5 83.8
30	1760	1543.0	217.0	97	1506	1597.0	- 91 .0	164	8 97	664.9	232.1
32	1503	1564.0	-61.0	99	1484	1477.6	6.4	166	709	664.5	44.5
33	1589	1564.9	82.1	101	1356	1359.2	-3.2	168	697	664.1	32.9
35 36	1652 1408	1552.2	99.8 -128.5	102 103	1455	1302.5	-63.2	170	728	663.8	64.2
37	1388 1291	1525.2	-137.2	104 105	995 1009	1196.8 1148.6	-201.8	172	675	663.4	11.6
39	1425	1535.9	-110.9	106	871	1103.5	-232.5	173 174	814 909	663.3 663.1	150./ 245.9
41	1735	1577.2	157.8	108	752	1023.4	-271.4	175	668 613	662.9 662.8	5.1 -49.8
43	1565	1590.0	-25.0	110	909	956.1	-47.1	177	800	662.7	137.3
44	1367	1548.3	-181.3	112	857	900.5	-43.5	179	644	662.5	-18.5
46	1363	1499.5	-136.5	114	785	855.2	-70.2	181	577	662.3	-85.3
48 49	1312 1424	1494.3	-182.3	115	749 864	835.8	45.5	182	622	662.1	-40.1
50 51	1615	1535.6	79.4	117 118	870 758	802.9 789.0	-31.0	184 185	671	662.0	27.1
52	1692	1606.0	86.0	119 120	772	776.6	-4.6	186 187	631 638	661.8 661.7	-30.8 -23.7
54	1544	1631.6	-87.6	121	695	755.5	-60.5	188 189	613 643	661.6	-48.6
56	1540	1580.8	-40.8	123	841 701	738.6	102.4	1 90	673	661.5	11.5
58	1463	1494.4	-31.4	125	719	725.1	-6.1	1 92	569	661.3	-92.3
59 60	1620	1446.5	173.5	127	595	714.3	-119.3	194	644	661.2	-17.2
61 62	1484	1452.7	-74.7	128	712	705.6	6.4	1 95	731	661.1	69.9
63 64	1444	1522.8	-78.8 146.6	130 131	/55 840	/01.8 698.5	53.2 141.5	197	679	661.0	18.0
65 66	1611 1338	1625.5	-14.5	132 133	603 829	695.5 692.8	-92.5 136.2	199 200	827 661	660.9 660.9	166.1
									=====		

were fewer observed points per diffraction fringe than for the previous event. Nevertheless, the fit was quite good, as an examination of the FITPLOT will reveal. The variance/covariance and correlation matrices of the solution parameters are given in Table 5-62.

The PDPLOT for this event is shown as Figure 5-128. A small amount of numerical noise was present in the computation of the partial derivative of the intensity curve with respect to the angular diameter. This numerical noise, however, was present only in the partial derivatives for times early in the solution space. At those times, as is apparent in the PDPLOT, the intensity curve was very insensitive to changes in the diameter. Hence, in this case the numerical noise which arose did not adversely affect the final solution determination.

The best fit to the observations was for the intensity curve of a point source. Thus, the angular diameter of X07145 was below the capability of detection by the occultation method.

The distribution function of the residual amplitudes, seen in Figure 5-192, is similar to that previously shown for the previous event. In this case, due to a smaller stellar signal in comparison to a slightly more variable background, a Poisson tail can be seen in the distribution function.

The POWERPLOT (Figure 5-130) indicates that the frequency components that are of most importance in the



TABLE 5-62

X07145: SUPPLEMENTAL STATISTICAL INFORMATION

VARIANCE/CO-VARIANCE MATRIX

DIAM	PREI	POST	TIME	VELO
1.430E ⁻ 16	2.593E-08	-2.853E 08	-4.964E ⁻ 10	8.640E_12
2.593E 08	1.703E02	4.570E00	-2.251E00	2.834E_02
-2.853E08	4.570E00	2.231E02	-2.817E00	3.503E_02
-4,964E 10	-2.251 E00	-2.817E00	5.799E ⁰¹	-7.327E_03
8.640E ¹²	2.834E 02	3.503E 02	-7.327E 03	1.271E ⁰ 4

CORRELATION MATRIX

	DIAM	PREI	POST	TIME	VELO
DIAM	1.000000	0.750746	-0.795461	0.161136	-0.153926
PREI	0.750746	1.000000	-0.197010	-0.523228	0.529208
POST	-0.795461	-0.197010	1.000000	-0.721394	0.716211
TIME	0.161136	-0.523228	-0.721394	1.000000	-0.999972
VELO	-0.153926	0.529208	0.716211	-0.999972	1.000000

NUMERICAL RANGES OF THE PARTIAL DERIVATIVES







Figure 5-129. NOISEPLOT of the occultation of X07145.



occultation spectral signature dominate over those in the star-plus-sky signal.

Geometrical occultation was found to have occurred at 01:28:05.604 (+/- 0.006 seconds) Coordinated Universal Time. X07202

The third event observed on the night of March 11, 1983 U. T., was the disappearance of X07202. This K2 star was substantially fainter (mV=8.21) than either of the stars previously observed on that night. The lesser brightness of this star in comparison to either X07145 or ZC0835 and a consequently poorer S+N/N ratio of 3.07 caused the derived solution to be of lower quality.

The raw photometric data are presented in Figure 5-131. The smaller change from pre-occultation to post-occultation signal intensity is rather obvious. Although the photometric conditions throughout the night remained virtually unchanged, the ascending branch of the integration plot (see Figure 5-132) does indicate that the sky became less transparent for approximately one second just prior to the event. This behavior, however, did not seem to persist at the time of the event itself. The integration plot also shows no evidence of "wide" stellar duplicity.

The observed and computed intensity curves are shown in Figure 5-133. Two hundred milliseconds of data were used in the DC fit. This seemed appropriate given the predicted L-rate of 633.40 meters/second. The observed L-rate, determined by the fitting process to be 1132 (+/-110) TABLE 5-63 X07202: LUNAR OCCULTATION SUMMARY

_____ STELLAR AND OBSERVING INFORMATION _____ Star: X07202, (SAO 077312, DM +24 0882) RA: 053529 DEC: +242136 mV: 8.21 Sp: K2 Filter: V Diaphragm: J Gain: C8 Voltage: 1100 LUNAR INFORMATION _____ 53 percent Surface Illumination: Elongation from Sun: 94 deorees 57 degrees Altitude Above Horizon: 57 degrees Lunar Limb Distance: 375321 kilometers Predicted Shadow Velocity: 633.40 meters/sec.

Predicted Angular Rate: 0.3481 arcsec/sec.

EVENT INFORMATION

USNO V/O Code:	15 H	A of Event:	+362114
Position Angle: 1	11.1 C	usp Angle:	66S

MODEL PARAMETERS

Number of Data Points: 201 Number of Grid Points: 256 Number of Spectral Regions: 50 Angstroms Limb Darkening Coefficient: 0.5 Effective Stellar Temperature: 4700 K

SOLUTIONS

 Stellar Diameter (ams):
 Point Source

 Time: (relative to Bin 0):
 2224.3
 (1.6)

 Pre-Event Signal:
 1679.6
 (18.5)

 Background Sky Level:
 1299.7
 (18.2)

 Velocity (meters/sec.):
 1132.1
 (110.3)

 Lunar Limb Slope (degrees):
 -28.1
 (10.0)

 U.T. of Occultation:
 0.21530.3666 (0.012)
 (0.012)

PHOTOMETRIC NOISE INFORMATION

Sum-of-Squares of Residuals:	6706927
Sigma (Standard Error):	183.1246
Normalized Standard Error:	0.4821239
Photometric (S+N)/N Ratio:	3.074155
Intensity Change/Background:	0.292238
Change in Magnitude:	0.2783562









Figure 5-134. Detailed FITPLOT of the occultation of X0702 showing the raw fit to a smoothed data set.

meters/second, was considerably faster. Thus, the observed diffraction fringes were comprised of correspondingly fewer observed data points than expected. This, in addition to the lower signal intensity, was responsible for the somewhat larger statistical uncertainties in the solution parameters.

In order to allow a better visualization of the solution, the computed intensity curve is shown in detail (over 100 milliseconds only) in Figure 5-134. The observed data were subjected to 5-point smoothing before being plotted on the detailed computed curve.

As expected, no diameter was determined. The PDPLOT illustrating the partial derivatives of the computed intensity curve is shown in Figure 5-135. Unquestionably, the regions of sensitivity to parametric variation were well covered in the solution. The observed and computed intensities and the residuals are given in Table 5-64. The variance/covariance and correlation matrices of the solution parameters and the numerical ranges of the partial derivatives are contained in Table 5-65.

The NOISEPLOT (Figure 5-136), as expected, is a bit more skewed toward a Poisson distribution than that of the previous event. Also, not unexpectedly, the one sigma width of the distribution was increased to 11.98 percent.

The POWERPLOT for this event indicates that the power components of the star-plus-sky signal at low frequencies were comparable to that of the occultation signal itself. Thus, any type of Fourier smoothing was out of the question.





Figure 5-136. NOISEPLOT of the occultation of X07202.

TABLE 5-64 K07202: OBSERVATIONS, COMPUTED VALUES, AND RESIDUALS FROM BIN 2125

==== N U M	OBS	COMP	RESID	N UM	OBS	COMP	RESID	N UM	OBS	COMP	RESID
=M - 0123456789012345678901234567890123456789012345678901234456789012234556789 =U- =N-	$\begin{smallmatrix} & & & & & & \\ & & & & & & \\ & & & & & $	 COMP	$\begin{array}{c} \text{RESID} \\ \text{RESID} \\359, \text{Residual} \\ -79, \text{r}, \text{r},$	LUM 	$\begin{array}{l} = = = & = & = & = & = & = & = & = & = $	COMP 1674.3 1764.3 1764.3 1764.3 1764.3 1764.3 1764.3 1764.3 1764.3 1764.3 1764.3 1742.9 1734.3 1742.9 1734.3 1742.9 1734.3 1742.9 1734.3 1742.9 1734.3 1742.9 1734.3 1742.9 1734.3 1742.9 1734.3 1742.9 1734.3 1742.9 1734.3 1742.9 1734.3 1734.3 1734.3 1734.3 1734.3 1735.0 1864.4 1735.0 1864.4 1735.0 1864.4 1735.0 1864.4 1735.0 1864.4 1735.0 1864.4 1735.0 1864.4 1735.0 1364.4 1337.4 1336.4 1337.4	$ \begin{array}{c} & = & 1 \\ = & $		$\begin{array}{c} \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\ \\$	$\begin{array}{c} \hline \\ \hline $	$ \begin{array}{c} \text{RESID} \\ \text{RESID} \\ \text{result} \\ 1000, 200, 200, 200, 200, 200, 200, 200$
59 61 62 63 65	1775 2063 1770 1443 1817 1439	1676.3 1680.4 1688.5 1682.0 1670.2 1676.7 1694.9	-115.3 94.6 374.5 88.0 -227.2 140.3 -255.9	126 127 128 129 130 131 132	1296 1405 1228 1365 1321 1447 1427	1301.6 1301.6 1301.5 1301.4 1301.3 1301.2 1301.1	-5./ 103.4 -73.5 63.6 19.7 145.8 125.9	193 194 195 196 197 198 199	1214 1207 1326 1141 1380 1423 1405	1299.9 1299.9 1299.9 1299.9 1299.9 1299.9 1299.9	-85.9 -92.9 26.1 -158.9 80.1 123.1 105.1
	1662	1694.8	-32.8	133	1092	1301.0	-209.0	200	1349	1299.9	49.1

TABLE 5-65 X07202: SUPPLEMENTAL STATISTICAL INFORMATION

VARIANCE/CO-VARIANCE MATRIX

DIAM	PREI	POST	TIME	VELO
4.401E-13	1.328E-06	-8.203E ⁻ 07	-7.665E ⁻ 08	5.838E ⁻ 09
1.328E 06	3.429E02	5.490E00	-4.822E00	2.791E_01
-8.203E 07	5.490E00	3.303E02	-4.658E00	2.704E_01
-7.665E 08	-4.822E00	-4.658E00	2.549E00	-1.489E_01
5.838E 09	2.791E ⁰¹	2.704E 01	-1.489E ⁰¹	1.218E 02

CORRELATION MATRIX

	DIAM	PREI	POST	TIME	VELO
DIAM	1.000000	0.886113	-0.638250	-0.216493	0.214986
PREI	0.886113	1.000000	-0.210694	-0.610273	0.607797
POST	-0.638250	-0.210694	1.000000	-0.582763	0.582383
TIME	-0.216493	-0.610273	-0.582763	1.000000	-0.999953
VELO	0.214986	0.607797	0.582383	-0.999953	1.000000

NUMERICAL RANGES OF THE PARTIAL DERIVATIVES

	MAXIMUM	MIN IMUM	
DIAM	6.264E08	-6.116E08	
PREI	1.362E00	4.004E ⁰ 4	
POST	9.996E 01	-3.620E ⁰¹	
TIME	5.504E01	-5.869E01	
VELO	9.125E02	-8.903E02	



The determined Coordinated Universal Time of geometrical occultation was 02:15:30.366 (+/- 0.012) seconds.

The relatively faint (mV=8.9) B9 star, X07202, occulted by the moon at approximately 03:17 U. T. on March 11, 1984, was the last of the four events observed on that night. As noted in the occultation summary table (Table 5-66) a Johnson B filter was used for this observation.

The photoelectric record of the observation is shown in the RAWPLOT, Figure 5-138. On cursory examination, the integration plot (Figure 5-139) suggests a possible "wide" stellar companion roughly 150 milliseconds after the anticipated disappearance. Further detailed examination of the occultation record, however, revealed that this visual artifact in the INTPLOT was due only to the noise in the observational data. The probability of this being due to the presence of a secondary star was found to be less than 0.1.

As anticipated, given the very early spectral type of the star and its faintness, the best solution found by the DC procedure was that of a point source. This point source solution is depicted graphically in the FITPLOT, Figure 5-140. The S+N/N ratio of this observation was only 2.69. This is reflected in the obviously noisy observational data shown along with the fit. The observational data used in the DC fitting, the computed intensities and the residuals are given in Table 5-67.

TABLE 5-66 X07247: LUNAR OCCULTATION SUMMARY

STELLAR AND OBSERVING INFORMATION Star: X07247, (SAO 077341, DM +24 0895) RA: 053708 DEC: +241825 mV: 8.9 Sp:B9 Filter: B Diaphragm: J Gain: C7+ Voltage: 1100 LUNAR INFORMATION ------Surface Illumination: 54 percent Elongation from Sun: 94 degrees 44 degrees Altitude Above Horizon: 44 degrees Lunar Limb Distance: 375983 kilometers Predicted Shadow Velocity: 618.7 meters/sec. Predicted Angular Rate: 0.3394 arcsec/s 0.3394 accsec/sec. EVENT INFORMATION ------ Date: March 11, 1984 UT of Event: 03:17:08 USN0 V/O Code: 14 HA of Event: +51:2416 Position Angle: 126.3 Cusp Angle: 51S Contact Angle: -36.1 Watts Angle: 127.5 MODEL PARAMETERS -----Number of Data Points: 201 Number of Grid Points: Number of Grid Points: 256 Number of Spectral Regions: 50 Width of Spectral Regions: 42 Angstroms Limb Darkening Coefficient: 0.5 Effective Stellar Temperature: 5100 K SOLUTIONS _____ Stellar Diameter (ams): Point Source Time: (relative to Bin 0): 1218.9 (3.4) Pre-Event Signal: 1967.4 (15.4) Background Sky Level: 1647.7 (31.3) Velocity (meters/second): 444.3 (44.1) Lunar Limb Slope (degrees): -20.1 (7.4) U.T. of Occultation: 03:17:07.618 (0.004) PHOTOMETRIC NOISE INFORMATION Sum-of-Squares of Residuals: 7116248 Sigma (Standard Error): 188.6304 Normalized Standard Error: 0.59007 Photometric (S+N)/N Ratio: 2.6937 (Change in Intensity)/Background: 0.19401 Change in Magnitude: 0.19252







X07247: OBSERVATIONS,	TABL COMPUTED	E 5-67 VALUES,	AND RESI	DUALS 1	FROM BIN 1	060
NUM OBS COMP RESID	NUM OBS	COMP	RESID	NUM OB	S COMP	RESID
$ \begin{array}{c} 0.00 \\ 0$	111 111 67 2013 68 1886 68 1887 68 1886 68 1887 67 2103 70 2505 71 1877 72 1877 74 1877 75 2137 76 2137 77 1911 83 1941 84 1941 84 1941 84 1943 89 1815 91 22758 89 1815 93 1262 84 1944 94 1599 95 1264 96 2348 99 1216 1001 124 102 1217 103 128 104 128 105 1936 106 1936 107 </td <td></td> <td>$\begin{array}{c} -4.44, 0.0 \leq 3.2$</td> <td>$\begin{array}{c} -223; 134; 22000, 2316; 2345,$</td> <td></td> <td>$\begin{array}{c} -1.18 & 5, 7, 5, 7, 5, 8, 10, 10, 10, 10, 10, 10, 10, 10, 10, 10$</td>		$\begin{array}{c} -4.44, 0.0 \leq 3.2 $	$\begin{array}{c} -223; 134; 22000, 2316; 2345,$		$\begin{array}{c} -1.18 & 5, 7, 5, 7, 5, 8, 10, 10, 10, 10, 10, 10, 10, 10, 10, 10$
66 2095 1967.7 127.3	133 1774	2083.6	-309.6	200 16	40 1651.4	-11.4

The PDPLOT of the solution intensity curve (Figure 5-141) shows nothing unusual or of concern in the final numerical computations of the partial derivatives. The numerical ranges of the partial derivatives, as well as the usual supplementary statistics, are listed in Table 5-68.

The stellar signal itself was only 319.7 counts, or only 19.4 percent of the background level. Hence, the slow rise (over 4096 milliseconds) in the now strongly dominant background contribution resulted in a stronger Poisson tail in the distribution function of the residuals (see Figure 5-142).

The Coordinated Universal Time of geometrical occultation for X07247 was 03:17:07.618 (+/0.004 seconds).

The four events which occurred on this night were of progressively fainter stars, observed under similar conditions. It is of interest to compare the degradation of the solutions (in terms of the formal errors of the determined parameters), the broadening and increasing Poisson nature of the residual amplitude distributions, and the increasing dominance of the low frequency power contributions in the pre-occultation (star-plus-sky) signal. The latter can be seen for X07202 in Figure 5-143.

X09514

The first of two occultations observed on the night of March 12, 1984 U. T., was that of the K5 star X09514. Due to an intermittent problem in the A-to-D converter input selected for this event, approximately half of the data which





VARIANCE/CO-VARIANCE MATRIX

DIAM	PREI	POST	TIME	VELO
1.371E ⁻ 14	2.264E ⁰ 7	-6.947E ⁻ 07	-1.822E 08	3.630E ⁻ 10
2.264E ⁰⁷	2.457E02	2.824E01	-1.200E01	1.398E 01
-6.947E 07	2.824E01	1.031E03	-4.002E01	4.384E_01
-1.822E 08	-1.200E01	-4.002E01	1.149E01	-1.297E_01
3.630E ⁻ 10	1.398E 01	4.384E ⁰¹	-1.297E ⁰¹	1.972E 03

CORRELATION MATRIX

	DIAM	PREI	POST	TIME	VELO
DIAM	1.000000	0.405449	-0.944025	0.762444	-0.749555
PREI	0.405449	1.000000	-0.083764	-0.242888	0.260790
POST	-0.944025	-0.083764	1.000000	-0.927240	0.919840
TIME	0.762444	-0.242888	-0.927240	1.000000	-0.999806
VELO	-0.749555	0.260790	0.919840	-0.999806	1.000000

NUMERICAL RANGES OF THE PARTIAL DERIVATIVES

	MAXIMUM	MIN IMUM
DIAM	3.548E08	-3.771E08
PREI	1.363E00	1.138E 02
POST	9.886E 01	-3.634E 01
TIME	2.070E01	-2.166E01
VELO	1.838E03	-1.876E03



should have been acquired were not. This problem was not known to exist at the time of the event but came to light after the fact. Fortunately, data dropouts did not occur during the occultation event itself, and hence, did not adversely affect the solution process. The RAWPLOT (Figure 5-144) shows only that portion of the data record for which photometric data were successfully acquired.

The integration plot, Figure 5-145, (for data taken over an interval of 1920 milliseconds) shows no indication of stellar duplicity.

The best fit to the observations is shown graphically on Figure 5-146. As expected, given the poor (S+N)/N ratio (only 2.27) and an apparent V magnitude of 8.7 for the star, no angular diameter could be determined. The solution parameters are given in the occultation summary, Table 5-69. The observed and computed intensities, and the residuals are listed in Table 5-70.

Figure 5-147 shows the sensitivity of the intensity curve to variations in the solution parameters. The numerical range of the values plotted on this PDPLOT can be found on Table 5-71, along with the variance/covariance and correlation matrices of the solution parameters.

The noise figure of the residuals, taken from the 1920 milliseconds of data successfully acquired, is presented as Figure 5-148. Figure 5-149 is the POWERPLOT of the solution curve, the occultation event, and the star-plus-sky signal.


TABLE 5-69 X09514: LUNAR OCCULTATION SUMMARY

STELLAR AND OBSERVING INFORMATION _____ Star: X09514, (SAO 078598, DM +25 1364) RA: 063835 DEC: +253153 mV: 8.7 Sp: K5 Filter: B Diaphragm: J Gain: C7+ Voltage: 1200 LUNAR INFORMATION _____ Surface Illumination: 65 percent Elongation from Sun: 107 degrees Altitude Above Horizon: 51 degrees Lunar Limb Distance: 369551 kilometers Predicted Shadow Velocity: 606.3 meters/sec. Predicted Angular Rate: 0.3384 arcsec/sec. EVENT INFORMATION ----- Date: March 12, 1984 UT of Event: 03:45:29 USN0 V/O Code: 14 HA of Event: +440653 Position Angle: 61.2 Cusp Angle: 57N Contact Angle: +37.1 Watts Angle: 56.4 MODEL PARAMETERS _____ Number of Data Points: 201 Number of Grid Points: 256 Number of Spectral Regions: 53 Width of Spectral Regions: 50 Angstroms Limb Darkening Coefficient: 0.5 Effective Stellar Temperature: 4200 K SOLUTIONS _____ Stellar Diameter (ams): Point Source Time: (relative to Bin 0): 3918.4 (3.7) Pre-Event Signal: 2596.5 (21.8) Background Sky Level: 2305.5 (24.3) Velocity (meters/sec.): 621.6 (77.2) Lunar Limb Slope (degrees): -6.36 (7.29) U.T. of Occultation: 03:35:32.8 (0.1) PHOTOMETRIC NOISE INFORMATION -------Sum-of-Squares of Residuals: 10568390 229,8738 Sigma (Standard Error): 0.78984 Normalized Standard Error: Photometric (S+N)/N Ratio: 2.26608 (Change in Intensity)/Background: 0.12624 Change in Magnitude: 0.12908



TABLE 5-70 X09514: OBSERVATIONS, COMPUTED VALUES, AND RESIDUALS FROM BIN 3800 NUM OBS COMP RESID NUM OBS COMP RESID NUM OBS COMP RESID NUM OBS COMP RESID 134 1951 2317-1 -366.1 135 2599 2316.1 282.9 136 2651 2315.2 331.4 137 2358 2314.4 55.6 138 2286 2313.7 -227.7 140 2361 2313.5 -315.5 142 2286 2313.5 -315.5 142 2086 2311.5 -315.5 143 2086 2311.5 -315.5 145 2086 2311.5 -315.5 145 2086 2311.5 -315.5 145 2086 2311.7 -225.7 145 2086 2311.7 -225.7 145 2086 2310.7 -215.5 145 2086 2310.7 -215.5 145 2086 2310.7 -215.5 145 2086 2310.7 -215.5 145 2086 2310.7 -215.5 145 2086 2310.7 -215.5 145 2086 2310.7 -215.5 145 2086 2309.5 -257.5 145 2039 2308.5 -140.2 155 2015 2308.5 -140.2 155 2015 2308.5 -140.2 155 2015 2308.5 -140.2 155 2015 2308.5 -140.2 155 2015 2308.5 -140.2 155 2015 2308.5 -140.2 155 2015 2308.5 -160.2 156 2191 2308.0 -117.0 157 2024 2307.7 -162.7 161 2409 2307.7 -162.7 161 2409 2307.3 -425.2 161 2407 3207.7 -425.7 161 2407 3207.7 -425.7 162 3407 3207.7 -425.7 163 2307.7 -425.7 164 213 3207.7 -475.7 164 213 3207.7 -475.5 154 2175 3207.7 -475.5 155 2077 3207.7 -475.5 155 2077 3207.7 -475.5 155 2077 3207.7 -475.5 155 2077 3207.7 -475.5 155 2077 3207.7 -475.5 155 2077 3207.7 -475.5 155 2077 3207.7 -475.5 155 2077 3207.7 -475.5 155 2077 3207.7 -475.5 155 2077 3207.7 -475.5 155 2077 3207.7 -475.5 155 2077 3207.7 -475.5 155 2077 3207.7 -475.5 155 2077 3207.7 -475.5 155 2077 3207.7 -475.5 155 2077 3207.7 -475.5 155 2077 3207.7 -475.5 155 2077 32077 3207.5 155 2077 3207.7 -475.5 --------217.7 416.8 118.7 -82.3 2815 2597.3 3014 2597.2 2716 2597.3 ٥ $\begin{array}{c} 79 & 2578 & 2643 \cdot 6 & -9 & -6 \cdot , 9 \\ 80 & 2659 & 2643 \cdot 6 & -9 & -8 \cdot , 9 \\ 81 & 2876 & 2643 \cdot 6 & -394 \cdot 1 \\ 82 & 2806 & 2651 \cdot 2 & -294 \cdot 8 \\ 81 & 2806 & 2651 \cdot 2 & -394 \cdot 1 \\ 83 & 2806 & 2551 \cdot 9 & -91 \cdot 1 \\ 85 & 2643 \cdot 2 & 551 \cdot 5 & -91 \cdot 1 \\ 85 & 2643 & 2551 \cdot 5 & -124 \cdot 1 \\ 88 & 2250 & 2554 \cdot 6 & -124 \cdot 1 \\ 89 & 2426 & 2554 \cdot 6 & -124 \cdot 1 \\ 89 & 2426 & 2554 \cdot 6 & -124 \cdot 1 \\ 89 & 2426 & 2554 \cdot 6 & -124 \cdot 1 \\ 99 & 2425 & 2556 \cdot 5 & -124 \cdot 1 \\ 99 & 2425 & 2650 \cdot 4 & -888 \cdot 4 \\ 99 & 2426 & 2644 \cdot 5 & -188 \cdot 5 \\ 99 & 2427 & 2666 \cdot 1 & 105 \cdot 9 \\ 99 & 2637 & 2767 \cdot 1 & -124 \cdot 1 \\ 99 & 2425 & 2650 \cdot 4 & -888 \cdot 4 \\ 99 & 2477 & 2767 \cdot 6 & -186 \cdot 5 \\ 99 & 2677 & 2707 \cdot 7 & -74 \cdot 3 \\ 99 & 2677 & 2707 \cdot 7 & -74 \cdot 3 \\ 99 & 2677 & 2707 \cdot 7 & -74 \cdot 3 \\ 99 & 2677 & 2767 \cdot 6 & -166 \cdot 5 \\ 100 & 2543 & 2685 \cdot 5 & -144 \cdot 5 \\ 100 & 2543 & 2685 \cdot 5 & -144 \cdot 5 \\ 100 & 2543 & 2685 \cdot 6 & -77 \cdot 24 \cdot 3 \\ 100 & 2543 & 2685 \cdot 7 & -248 \cdot 7 \\ 100 & 2543 & 2685 \cdot 7 & -248 \cdot 7 \\ 100 & 2543 & 2686 \cdot 9 & -36 \cdot 9 \\ 110 & 2520 & 2566 \cdot 7 & -28 \cdot 7 \\ 100 & 2543 & 2449 \cdot 6 & -78 \cdot 69 \cdot 5 \\ 1110 & 2520 & 2440 \cdot 6 & -78 \cdot 69 \cdot 5 \\ 112 & 2430 & 2443 \cdot 5 & -78 \cdot 5 \\ 114 & 2437 & 2443 \cdot 5 & -78 \cdot 5 \\ 116 & 2352 & 2356 \cdot 5 & -128 \cdot 7 \\ 116 & 2350 & 2358 \cdot 4 & -78 \cdot 69 \cdot 5 \\ 117 & 2351 & 2338 \cdot 4 & -29 \cdot 8 \\ 118 & 2449 & 2337 & 2443 \cdot 5 & -78 \cdot 5 \\ 122 & 2244 & 2347 \cdot 0 & -78 \cdot 69 \cdot 5 \\ 114 & 2237 & 2338 \cdot 4 & -29 \cdot 8 \\ 118 & 2449 & 2326 \cdot 5 & -128 \cdot 7 \\ 122 & 2244 & 2328 \cdot 4 & -118 \cdot 4 \\ 121 & 2237 & 2338 \cdot 5 & -114 \cdot 3 \\ 122 & 2244 & 2328 \cdot 4 & -115 \cdot 4 \\ 122 & 2244 & 2328 \cdot 4 & -115 \cdot 4 \\ 122 & 2244 & 2328 \cdot 4 & -115 \cdot 4 \\ 122 & 2244 & 2328 \cdot 4 & -115 \cdot 4 \\ 122 & 2244 & 2328 \cdot 4 & -118 \cdot 4 \\ 121 & 2237 & 2338 \cdot 5 & -78 \cdot 5 \\ 122 & 2244 & 2328 \cdot 4 & -115 \cdot 4 \\ 121 & 2237 & 2338 \cdot 5 & -114 \cdot 3 \\ 221 & 2238 \cdot 2328 \cdot 4 & -115 \cdot 4 \\ 221 & 2238 & 2328 \cdot 4 & -115 \cdot 4 \\ 221 & 2238 & 2328 \cdot 4 & -115 \cdot 4 \\ 221 & 2242 & 2328 \cdot 4 & -115 \cdot 4 \\ 221 & 2242 & 2328 \cdot 4 & -115 \cdot 4 \\ 221 & 2242 & 2328 \cdot 4 & -115 \cdot 4 \\ 221 & 2242 & 2328 \cdot 4 & -115 \cdot 4 \\ 221 &$ 161 2879 162 2476 163 1875 164 2133 165 2265 166 1894 -413.1 -413.1 -34.0 -223.0 387.1 -74.8 44.2 -74.0 44.2 -177.7 -71.7 182.4 -194.6 -50.6 -190.5 255.5 344.6 344.0 -2.4 -263.4 13.6 18.7 122.7 98.7 -85.3 87.8 -26.2 -342.2 102.8 112.9 33.9 2281 2599 2620 2359 2321 2722 2576 2303 2558 2588 2825 2533 2092 $\begin{array}{r}
18.6 \\
-249.4 \\
-289.6 \\
115.9 \\
-20.5 \\
-283.1 \\
-28.9 \\
6.8 \\
234.8 \\
-70.7 \\
-70.7 \\
-70.7 \\
\end{array}$ $\begin{array}{r} 2610.6 & -289.6 \\ 2606.1 & 115.9 \\ 2596.5 & -20.5 \\ 2586.1 & -283.1 \\ 2579.9 & -28.9 \\ 2581.2 & 6.8 \\ 2590.2 & 234.8 \\ 2603.7 & -70.7 \\ 2616.5 & -524.5 \end{array}$ 59 5.9 -305.1 61 457.0 133.0 -75.0 -243.0 64 65 66



Figure 5-147. PDPLOT of the occultation of X09514.



Figure 5-148. NOISEPLOT of the occultation of X09514.

TABLE 5-71 X09514: SUPPLEMENTAL STATISTICAL INFORMATION

VARIAN CE/CO-VARIAN CE MATRIX

DIAM	PREI	POST	TIME	VELO
3.254E ⁻ 14	5.005E 07	-5.270E 07	-5.312E 08	1.290E_09
5.005E 07	4.841E02	1.686E01	-1.663E01	2.919E ⁰¹
-5.270E 07	1.686E01	6.996E02	-2.319E01	4.042E 01
-5.312E08	-1.663E01	-2.319E01	1.380E01	-2.444E_01
1.290E 09	2.919E 01	4.042E 01	-2.444E ⁰¹	5.961E 03

CORRELATION MATRIX

	DIAM	PREI	POST	TIME	VELO
DIAM	1.000000	0.766505	-0.761986	0.128535	-0.124485
PREI	0.766505	1.000000	-0.172303	-0.479860	0.481833
POST	-0.761986	-0.172303	1.000000	-0.713209	0.709453
TIME	0.128535	-0.479860	-0.713209	1.000000	-0.999975
VELO	-0.124485	0.481833	0.709453	-0.999975	1.000000

NUMERICAL RANGES OF THE PARTIAL DERIVATIVES

	MAXIMUM	MIN IMUM	
DIAM	3.055E08	-3.000E08	
PREI	1.365E00	1.974E ⁻ 03	
POST	9.980E 01	-3.649E ⁻ 01	
TIME	2.322E01	-2.477E01	
VELO	1.311E03	-1.336E03	
	DIAM PREI POST TIME VELO	MAXINUM DIAM 3.055508 PREI 1.365500 POST 9.980501 TIME 2.322E01 1.311E03	MAXINUM MINIMUM DIAM 3.055508 -3.000508 PREI 1.365600 1.974E^03 POST 9.9805^01 -3.649E^01 TIME 2.322801 -2.477801 VEL0 1.311E03 -1.336E03



Unfortunately, as had occasionally happened in previously observed events, the signal strength of the WWUB time code was too low to provide a useful time reference at the time of the event. A successful post-event calibration of the SPICA-IV/LODAS clock could not be obtained until 2 hours after the event. Assuming a worst case internal clock drift of 2 seconds per day, the Coordinated Universal Time of geometrical occultation was found to be 03:35:32.8 (+/- 0.1 seconds).

ZC1030 (Epsilon Geminorum)

ZC1030 (Epsilon Geminorum, Mebsuta, BS 2473) was a prime candidate for the occultation program. This 65 supergiant (luminosity class Ib) is on the order of 6000 times as luminous as the sun. Its apparent V magnitude of 3.18 gives Epsilon Geminorum a corresponding distance of appoximately 350 parsecs. To gauge the possibility of a diameter detection for Epsilon Geminorum, Allen (1963) gives diameters of 65-I stars to be in excess of 120 solar diameters. At this distance this linear diameter corresponds to an angular diameter on the order of 3 or 4 milliseconds of arc, well above the detection threshold. On April 7, 1976, Epsilon Geminorum was occulted by Mars. This well-observed planetary occultation has been discussed by Wasserman et al. (1977).

The intermediate bandwidth "y" filter was initially selected for this observation. This star, brighter by far than any other previously observed in the occultation

program, should have produced an observation with an unprecedented signal-to-noise ratio. Though X09514 had been observed under clear skies only two hours earlier, the weather conditions deteriorated rapidly. Twenty minutes before the predicted time of the event, the sky was completely covered with a heavy layer of cirrus clouds, and no stars could be seen with the unaided eye. Just prior to the event, no surface markings could be distinguished on the lunar surface. (The observing log, in fact, referred to the lunar image as "appearing like a milk bottle smeared with mayonnaise"). The star, however, was visible in the photometer viewing optics against a sky background which was chalky white. Though the clouds had attenuated the stellar signal by at least an order of magnitude, the star-plus-sky signal was still 40 percent higher than the sky background.

Under these conditions, due to adverse seeing and transparency effects, the possibility of obtaining data of sufficient quality to yield a stellar diameter was out of the question. Yet, though this fundamentally important piece of astrophysical information was to be lost, at least the time of geometrical occultation could still be obtained. The data acquisition rate was reduced to 100 points per second, while leaving the amplifier time constant at 2 kiloHertz. While this would result in severe undersampling for the purpose of stellar diameter detection, it was necessary under the poor observing conditions in order to extract a usable time of occultation. Five minutes before the predicted time of the





Figure 5-151. INTPLOT of the occultation of 2C1030.

event the selected filter was changed to the "b" filter, as the star-plus-sky/sky signal ratio was a approximately 10 percent higher than with the "y" filter.

The RAWPLOT of the event, Figure 5-150, shows the slowly acquired photometric data. The integration plot, Figure 5-151, is indicative of the variation in atmospheric transparency brought on by the layer of intervening clouds and haze. As can be seen, the change in signal level due to the occultation is unmistakable. No diffraction fringes are contained in the data (as expected) due to the slow data acquisition rate. The data sample corresponding to the time of disappearance was noted strictly by visual inspection. After reduction of the WWVB time code the Coordinated Universal Time of geometrical occultation was found to be 05:50:06.77 (+/- 0.01 seconds).

Under normal circumstances such an observation would have been impossible, and the presence of a layer of heavy cirrus clouds would have prevented any observing attempt. Yet, as was shown by the partial success of this observation, a bright star can yield some information (even if only the time of occultation to low precision) under adverse conditions. Occultation photometry, for chronometric and low precision astrometric work, is guite forgiving.

Summary of the Occultation Observations

The observations which have been discussed in the previous section were all made during the thirteen lunations

spanning the time period from March 1983, to March 1984. During this period, 67 occultation observations were planned. Of these, photometric data were obtained for 24 occultations. The remainder, with the exception of two, were lost due to inclement weather. The two other events mentioned were not observed due to a mechanical failure in the instrumental system, which was remedied in time for the next night's observing.

Of the 24 events for which data were obtained, two (the occultations of ZC3458 and ZC1030) were observed with a data acquisition rate which was too slow to allow a complete analysis to be carried out. If the observing year discussed was typical (in terms of weather), then one might expect approximately one fruitful occultation observation for every three planned for the Rosemary Hill Observatory.

The stars ZC0916-A, X13534, and X13607 were discovered to be "close" occultation binaries. (The B component of ZC0916 was previously known to be a spectroscopic binary). Hence, though based on admittedly small number statistics, 14 percent of the 22 observations which were fully analyzed resulted in the discovery of previously unknown "close" companions. Widely separated, though much fainter, companion stars were found for ZC1221, X18067, ZC0126. Thus, 14 percent of the stars studied were revealed to have fainter "wide" components. Considering both "close" and "wide" systems, 27 percent of the stars studied turned out to have previously unknown companions. Table 5-72 summarizes the

TABLE 5-72

DERIVED	QUANTITIES	FOR	THE	OCCULTATION	BINARIES
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		-				
	STAR	т -	mV-T	mV-1	mV−2 	P. S.
	ZC0916-A	С	4.74	5.05	6.27	3.34
	ZC0916-B	С	5.19	5.28	7.66	8.41
	ZC0916-AB	ω	4.30	4.74	5.19	31.5
	ZC1 221	ω				440
	X13534	С	8.4	8.73	9.86	6.07
	X13607	С	8.2	8.58	9.52	13.15
	ZC0126	ω				38.8
Ν	otes: 1. Typ	e co	des (T)	: W = w	ide, C	= close.

Notes: 1. Type codes (1): W = Wide, C = Close. 2. mV=T is combined mV for both stars. 3. Projected Separation (P. S.) is milliseconds of arc.

STELLAR ANOULAR DIANETERS													
STAR	Diam.	S.E.	S+N/N	Q									
ZC0916-B1	5.89	0.84	12.680	427.6									
ZC1 221	2.70	2.41	9.044	62.4									
X07598	5.45	2.04	4.657	22.4									
ZC1 462	9.27	3.69	5.666	24.9									
ZC2209	12.18	1.86	9.904	43.3									
X01217	2.88	4.98	3.145	1.2									
ZC3158	1.74	1.68	4.578	90.3									
ZC0835	1.08	1.46	11.808	265.9									
Note: S.E.	is the St	andard	Error (One	Sioma)									

TABLE 5-73 STELLAR ANGULAR DIAMETERS

TABLE 5-74

COURDINATED	UNIVERSAL TIME	S OF GEOR	TEIRICAL OCCULIA	110NS
STAR	с. U. т.	S. E.	STAR-2 C.U.T.	S. E.
ZC0916-A ZC0916-B	02:44:33.347 02:44:33.421	0.001	02:44:33.326 02:44:33.413	0.003
ZC1 221	00:44:55.877	0.012	00:44:54.731	0.018
ZC1 222	01:14:46.177	0.004		
×07589	01:35:12.172	0.051		
×07598	02:12:02.030	0.005		
×13534	02:46:29.632	0.003	02:46:29.594	0.008
×13607	04:33:58.827	0.007	04:33:58.773	0.009
ZC1462	05:13:37.199	0.001		
ZC2209	03:12:53.356	0.007		
×18067	02:16:52.406	0.007		
ZC3214	03:35:43.107	0.018		
×13590	01:02:40.701	0.014		
X01217	00:02:07.790	0.004		
×01246	01:20:01.938	0.011		
ZC0126	01:28:01.981	0.010	01:28:01.753	0.010
×01309	03:32:06.631	0.005		
ZC3158	00:06:39.731	0.004		
ZC0835	00:39:28.8	0.1		
X07145	01:28:05.604	0.006		
X07202	02:15:30.366	0.012		
X07247	03:17:07.618	0.004		
×09514	03:35:32.779	0.1		
ZC1030	05:50:06.77	0.01		

Note: S. E. is the Standard Error (One Sigma) in seconds.

projected angular separations found for these double stars. In addition, the derived apparent V magnitudes are given for the "close" doubles.

Stellar angular diameters were determined for the stars ZC0916-B1, ZC1221, X07598, ZC1462, ZC2209, X01217, ZC3158, and ZC0835. However, the uncertainties in the angular diameters determined for ZC1221, ZC3158, ZC0835 and X01217 (all with angular diameters less than 3 milliseconds of arc) were rather large. In each of these cases the formal one sigma uncertainty of the diameter was nearly equal to the determined diameter itself. Indeed, in the case of X01217, the uncertainty in the diameter exceeded the angular diameter by 73 percent. The roughly 9 millisecond of arc diameter found for ZC1462 must be viewed with caution, as previously discussed. It is highly likely that, in this case, the diameter found was spurious.

Perhaps one of the two most interesting observational results obtained in the course of this investigation was the determination of the unexpectedly large angular diameter of ZC2209 (32 Librae). This determination, of approximately 12 milliseconds of arc, is quite good and cannot be easily dismissed.

Table 5-73 lists the angular diameters and their statistal uncertainties for each star for which a diameter was found. Also found on this table (as given separately in the occultation summary tables) are the photometric (S+N)/N ratios and the ratio of the pre-occultation to

post-occultation signal intensities. A "quality index", denoted Q, has been assigned to each of these events. Q is defined as the product of the (S+N)/N and intensity ratios divided by the normalized standard error of the photometric O-C's. The "quality index" is not defined from any rigorous standpoint. Rather, it is designed merely to give, in a single quantity. in indication of the overall photometric conditions for a particular event. Thus, an intercomparison of Q for different events provides a relative scale for judging the quality of the observational data used in the parametric solutions.

The Coordinated Universal Times of geometrical occultation were determined for each of the 24 events. In the case of "close" doubles, the individual times of disappearance of the components were found. These times, which have been reported to the International Lunar Occultation Center (in Japan), are summarized along with their one sigma uncertainties in Table 5-74. These statistical uncertainties arise both from the internal formal error in the DC solution for the time parameter and the error inherent in edge detection of the WWVB time code.

For each individual event, the determined R and L rates depend upon the local slope of the lunar surface. In principle, for any one event the slope can assume any value. However, the likelihood of encountering an extremely large local slope should be quite small, as near-vertical protuberances are rare in comparison to the total lunar

surface area which can be seen along the limb. Clearly, there should be a most likely value for a lunar slope based on the distribution function of individual slopes. This distribution function could be found by measuring the slopes of lunar surface features from the Lunar Orbiter or Apollo photographs. However, this would be a formidable task, as the lunar surface would have to be treated as a grid of no more than a few meters on a side in order to obtain a distribution function useful in the reduction and analysis of occultation observations.

If a detailed lunar slope distribution function were known, it could serve to assess the probability of a proper determination of the slopes for an individual event. In the course of this investigation, 27 local lunar slopes were determined, and the derived distribution function is shown in Figure 5-152. This figure, binned into regions 3 degrees wide, shows the occurrences of the absolute values of the determined slopes.

One would not expect a preferential alignment of lunar surface features, and hence, the mean of all slopes should average to zero. (Indeed, the mean slope angle found for the 27 observations was found to be -1.8 degrees). The mean absolute slope was approximately 15 degrees. The largest slope encountered was just in excess of 33 degrees. This distribution function is in good agreement with the distribution of lunar slopes which have been reported from observations made at McDonald Observatory over the last decade.



NUMBER OF OCCURRENCES

Future Directions for the Occultation Program

In order to meet the goals of the program of occultation observation discussed in Chapter 1, an instrumental system capable of obtaining fast photometric observations of lunar occultations was developed and implemented at the Rosemary Hill Observatory. Further, careful consideration of the methods available for the reduction and subsequent analyses of these observations led to improvements in the conventional numerical and computational procedures. These improvements were incorporated into a set of algorithms, implemented as APL functions, to allow a routine program of analysis to be carried out.

The systems and procedures developed have been proven viable as demonstrated by the successful determination of stellar angular diameters, the discovery of previously unsuspected multiplicity in several stellar systems, and in measuring the times of occultation events to a degree of precision useful to ongoing astrometric programs. This is not to say that there is not room for improvement. Indeed, on many levels the overall occultation program can be expanded and improved upon.

Improvements to the instrumental system at RHO could lead to the successful observation of at least twice as many occultations as are now currently possible. Specifically, all events observed to date were dark limb disappearances. For reappearance events, the star cannot be seen (as the moon is interposed between the star and the telescope prior to the event), and must one rely on the pointing accuracy of the telescope to acquire the unseen star. The 76-centimeter Tinsley reflector, however, lacks the pointing accuracy required to acquire a field typically of less than 15 arc seconds without visual reference.

A 4-inch refracting telescope with a high power eyepiece (to serve as an offset guider) was added to the 76-centimeter telescope. This was done in hope of offsetting to a nearby star, prior to the disappearance of the star under study, and observing the reappearance by guiding on the offset star. This modification, unfortunately, was unsuccessful due to differential flexure in both the telescope superstructure and the secondary mirror support. The mechanical problem of flexure is a difficult one to overcome without major modifications to the telescope structure itself.

A concerted effort to modify the drive system of the RHO 46-centimeter telescope is now beginning. New drive gears along with stepping motors for both the Right Ascension and Declination axes will allow computer control of the telescope for tracking and target acquisition. If this system is successful, the experience gained in its implementation could be used in upgrading the tracking system of the 76-centimeter telescope in a similar manner. With the ability to program differential offsets for target acquisition, the observation of dark limb reappearances becomes feasible.

Bright limb or daylight events could not be observed by the photometric system described in Chapter 2. The

brightness level of the background would literally swamp the stellar signal for even the brightest of sources if observed in the visible region of the spectrum. For a number of years observations of lunar occultations of late stars in the infra-red have been made with oreat success (as in the case of Aldeberan as reported by White and Kreidl, 1984). While such observations are not normally as effective from low altitude sites such as RHO as at high altitude sites, K-Y Chen (1985) has indicated that infrared observations, nevertheless, can be made rather effectively. Recently, P. Chen has been testing a new infrared photometric system on the 76-centimeter telescope. Though the opportunities for infra-red observations of lunar occultations from RHO would be more restrictive than from a high altitude site, on the basis of previous work done at RHO there is reason to believe that such observations would be successful.

Clearly, the next logical step to enhance the occultation program is to move in the direction of simultaneous multicolor observations. This could be done at a minimum expense, and the gain would certainly be worth the small additional effort. The Astromechanics photometer, as previously mentioned, can be used in two channels (i.e., blue and red) simultaneously. A three channel photometer built by Flesch (1975) and used in an investigation of flare stars was recently anodized and could be used for three color observations. Use of the latter would require adding an additional input channel to the SPICA-IV/LODAS for WWWB time code detection as all three 12-bit A-to-D's would be used for acquisition of photometric data. However, the time signal could be detected with only one bit (since it is digital in nature), and this additional input could be added with little impact on the memory allocated for photometric data storage.

Multicolor observations would be advantageous from two viewpoints. First, the simultaneous solution of the residual equations resulting from multicolor observations would be more highly constrained than in the case of one single observations. The time of geometrical occultation must be the same for all colors. This condition, when enforced in the parametric solution, would result in an improved decoupling and, therefore, a more reliable determination of the solution parameters. In addition, while the diffraction fringe spacing scales linearly (in natural units) with wavelength, the effects due to irregularities on the lunar limb do not. Thus, multicolor observations would help differentiate limb effects from diffraction effects for the small number of observations where limb irregularities would otherwise hamper a meaningful solution.

In most cases, for the occultation events which were discussed, the uncertainty in the determined times of geometric occultation did not reflect the statistical uncertainty in the DC solution. Usually, the primary source of the uncertainty in the event timings arose from an inability to detect second transitions in the WWVB time code to a degree commensurable with the precision in the

determination of the relative time of geometrical occultation. This situation, brought on by poor radio reception, could probably be improved with a higher gain VLF antenna. A new antenna configuration should be investigated as one of the first items considered in making continued improvements to the occultation instrumental system.

In terms of the numerical procedures used in determining the solution parameters, the consideration of second order terms in the residual equations, as suggested by Eichhorn and Clary (1974), may soon be computationally feasible. The primary reason for the linearization of the residual equations was to reduce the overall computer time necessary to acheive a convergent solution (see Chapter 3). The new generation of array processors or "supercomputers" may remove this artificial computational constraint. Indeed, at least one manufacturer (Analogics Inc, 1984) has announced an array processing computer, with effective execution speeds on the order of tem million floating point operations per second, specifically designed for APL.

The question of the distribution function of lunar limb slopes is not an idle one. Eichhorn (1977) has noted that the least squares solution process could be improved by the introduction of probabilistic constraints in the adjustments. One might reasonably assume that the distribution function of the pre-occultation and post-occultation intensities is a Gaussian (or possibly Poisson) distribution. The distribution function of the lunar limb slopes must, however,

be determined empirically. As more occultation observations are made and solved, the better the Knowledge of this distribution function becomes. Thus, this is a self-improving process. If one constrains the adjustment of the L-rate on the basis of the observed slope distribution function, then each newly determined slope should be added to the growing baseline data for that distribution. The inclusion of probabilistic constraints, like the consideration of second order terms, is a refinement that deserves investigation and possible future implementation.

The currently discussed observations do leave a number of questions unanswered, as pointed out earlier. As an example, the detected fourth component in the 1 Geminorum system remains a bit of an enigma. Griffin (1984) notes that the probability of non-detection by use of his photoelectric "radial velocity meter" (due to a near zero inclination of the orbital plane) is approximately one part in 800. Yet, the evidence for the existence of the A2 component is quite strong. Photoelectric radial velocity meter measurements, however, are sensitive to the spectral types of the component stars. This instrument requires that a mask be employed whose physical characteristics strongly depend on the anticipated spectral characteristics of the stars under study. The analysis of the spectra of composite stellar systems is enormously complex, and it is quite easy to imagine that a previously unsuspected component could possibly bias the interpretation of those spectra. In any

event, the 1 Geminorum system is certainly one that bears further investigation.

Also lurking in the relm of unanswered questions is the unexpectedly large diameter found for ZC2209. Corroborative (or refuting) observations are most certainly needed, and observations of future occultations of this star must be planned and implemented with care.

No observational program can operate in a figurative vacuum. Lunar occultation observations, in particular, could benefit from collaborative observations from more than one site. Since the position angle of the event will vary, as seen from topocentrically different observing stations, stars which are double (or multiple) could be completely solved in terms of the positions of the components. Multiple observations at different lunar geometries will also make any possible distorting limb effects readily apparent. And, of course, multiple observations must independently yield consistent results, and hence, are essential in the confirmation of those results. Efforts should be made, if indeed the occultation program is to continue, to seek out other institutions carrying out similar programs to mount collaborative efforts.

A vacuum, however, in the literal sense, may indeed provide the ideal conditions to observe lunar occultations. As early as 1970, Nather pointed out that the moon, as seen from an orbiting spacecraft, would be stationary in Right Ascension twice each orbit. At those times the time-scale of

the movement of the diffraction pattern resulting from the occultation, as projected onto the spacecraft, would be slowed enormously. From a suitably equipped orbital observing platform, data acquisition sampling rates on the order of 1 to 10 Hertz (as opposed to 1000 Hertz) could be employed with no degradation (i.e., smearing) of the diffraction pattern. Thus, from photon statistics alone, the (S+N)/N ratios could be improved by a factor of 10 to 30. This, of course, does not include the added gain of eliminating the effects of atmospheric scintillation, seeing, and transparency variation. The observed intensity curve would be adversely affected only by photon arrival statistics.

By the end of the next calendar year, the Hubble Space Telescope (HST) will be operational. The possibility of using a 94-inch orbital telescope, equipped with a high speed photometer, for spaceborne occultation observations is intriguing. There are two problems which would have to be overcome. First, normal HST operation guidelines disallow the pointing of the telescope in the vicinity of the moon due to possible damage to the various light sensitive detectors in the instruments. Second, HST requires two guide stars in different "pickles" (as described by Giacconi, 1982). The process of guide star selection might prove, in most cases, to be almost impossible. If the target star is about to be occulted then at least one and a half of the three fine guidance sensor "pickles" would be in the moon's shadow,

obscuring any possible guide stars. These are technical problems which need careful study. It may indeed turn out that, despite the ideal environment for occultation observations, HST might not be suitable for a majority of candidate events.

One cannot help but reflect on the nearly mind boggling array of spaceborne and "new generation" astronomical instruments which will soon become a reality. The spectrum spanning network of the HST, the Space Infra-Red Telescope Facility, the Advanced X-Ray Astonomy Facility, the Gamma Ray Observatory, as well as UV and EUV telescopes will all be in orbit and operational in the not too distant future. New multi-meter ground-based telescopes with state of the art support equipment are being planned and will be on-line well before the turn of the century. The new 10-meter telescope to be located on the summit of Mauna Kea would be an unparalled achievement in itself, and even now, there is a glimmer of a possibility that two such instruments might be built and operated as an optical interferometer. These new instruments will open up new vistas as of yet unimagined.

Some vocal proponents of spaceborne astronomical facilities have remarked that the days of ground-based observatories are numbered. At the same time, other voices have remarked that ground-based astronomy will be accomplished only by large (greater than approximately 3 meter diameter) telescopes. With "competition" from some of the most sophisticated and finely built opto-electronic

instrumentation ever to be constructed, some ask if it is now folly to plan a program for a small, ground based instrument.

The answer to both the cynics and skeptics alike is an unequivocal no. It is a fact that in terms of scientific return (expressed in the now-in-vogue units of photons per dollar), it is the small, research grade instruments which can be the most productive. Institution oriented astronomical sites, such as the University of Florida's Rosemary Hill Observatory, have the luxury of carrying out long term projects, unthinkable and misplaced on large, heavily oversubscribed instruments.

Specifically, telescopes in the 1 meter class are ideal for carrying out occultation observations in the visible and perhaps even near IR wavelengths. In the case of RHO, its other concurrent research programs require dark time, leaving moonlit nights otherwise unused. What better way to fill the unproductive hours than by adding to the knowledge of fundamental astronomical data in the measurement of stellar diameters, elucidation of the geometry and physical makeup of multiple systems, and compilation of a long term baseline of timing data for astrometric purposes?

It is to this end that this program of the observation and analysis of lunar occultations was established. It is hoped that this program will continue in full force and perhaps similar programs implemented at other small, institution oriented astronomical facilities.

APPENDIX A LISTING OF THE 6502 PROGRAM LODAS (VERSION E07)

This appendix contains the source code for the Lunar Occultation Data Acquisition System software, which is described in Chapter 2. The revision shown here (E07) is currently in use at the Rosemary Hill Observatory.

The assembly of this program was carried out using the ASM6502 cross-assembler, available at the Northeast Regional Data Center (NERDC). The assembly directives seen in this source listing are described by NERDC (1980).

The one statement which is flagged as an error in the compilation actually is not. The ASM6502 cross-assembler has a minor "bug" which caused this valid statement to be flagged. It is assembled correctly, and the erroneous error message does not affect the generated machine code.

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ICE STMT	Constraints 10,000 (10,000)))))))))))))))))))))))))))))))))	"REWRITE THE DISPLAY WITH TIME. ACTION CODE. CHANNEL COUNTS	LY JSR DISBUF ; DUTPUT 015PLAY BUFFER Los #400 ; CLEAP UPDATE 11ME 015PLAY FLAG	JMP MAIN : WEWS GO BACK TO MAIN WWWS	. OUTPUT THE DISPLAY BUFFER.	UF DOX #114 1111ALIZE SECHENI COUNTER TO ONE PAST LAST SECHENI 122 DEX 1 DECREMENT COUNTER COUNTER OUT ON FUSHING 11 ON SIACK 1114	PM BUFF, THE UNGOIN ROUTINE CLEBERS FIELS REFEIRED LOADBUFF, LOAD THE ACCONULATOR BITH CHARACTER TO BE DAD 8990 351 HI BI FOR DIRFCI OTSELV. JSO OUTOOL OUTOOL OUTOOL AREGISTER TO SECRET A	TAX • • RESTORE THE SEGREN CONTERT A REGISTER TAX BKE DISBF2 : IF MORE TO OUTPUT GO OD 17 RTS : ALL DONE	# PUT TWO BYTE ASCII ENTO OUTPUT BUFFER+	CZ TAY LIVE TAY LUE TA AS OFFSET 10 DECIMAL CHAR TARE CZ TAY CLUELY TAGET TEE DECIMAL VALUE CZ STAT TAPUDO T SAVE ASKED DECIMAL IN TERPERARY LOCATION IMPUDD T SAT THORA T GET LEFT HALF DORO OF A REGISTED ACTION IMPUDD	ADC #330 CONVERT TC ASCII	LAA TWPUPO 5 SET FLAM WEAK OUTER PLOSITION LOA TWPUPO 5 GET RACK CRIGIAAL DECIMAL VALUE UT AND #50F 5 LOOK AT SECONO DECIMAL DIGIT	ACC #530 1 CONVERT TO ASCII Sia Douff.x 1 Dut in Buffer POSITION X P12		C3 JSR HX18CD 2 CONVERT HEX TO TWO BYTE DECIMAL 314 TMPUPD 2 PUT TENS AND 016173 1A TENP LCCATION 744	LAR NOWUT PUT ACCUT CHAR FOR THIS DIGIT IN DISPLAY BUFFER LDA TWPUPO I GET TENS AND DIGITS VALUE	INX 55ET NEXT 015PLAY GUFFER POSITION JSR HIDUF ; PUT THESE TWO BYTES IN DISPLAY BUFFER RTS	BENTER COMMENTS ON PRINTER FROM JPOS	NT JSR PNDM 3 DUTPUT THE COMMENT PROMPT > FCC >	FCN \$00 X REGISTER USED AS CHARACTER COUNTER LOX \$900 X REGISTER USED AS CHARACTER COUNTER 1 JSD GETKEV GET A CHARACTER FROM KEVEDARD	DNE COMC2 IF NOT GO TO CONC2 DNE TOMC2 IS THIS THE FIRST CHARACTER?	DER CONCERTE IN SOUTH AND				
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 TO STACK \$2015PACE) + \$80 IND CURSOR1 DUTPUT BLANK TO DISPLAY AT CURRENT PDSITION USE X AS DISPLAY ELEMENT COUNTER \$13=20 GET NEXT CMARACTER Restore Location of Calling Rcutine GET ADDRESS OF CALLING ROUTINE STORE LOW BYTE NEXT POSITION Not done, more spaces needed carlage reven line feed residre registers SAVE THE A REGISTER ON STACK SAVE THE X REGISTER ON STACK CLOCK TIMER, HOURS, MINUTES, SECONDS+ # TOGGLE THE PRINTER # RESTORE OLD ACTION CODE DISPLAY MESSAGE FROM JPD+ PRESS, WALT IF NO ENTRYS +CLEAR DISPLAY, FRDW JPO+ + TOGGLE PRINTER ON/CFF + -----*GET A KEY AND S + TO GGLE PR #\$A0 0UTD01 JSR PHXY JSR KEYCK CMP #SFF JSR PLXY RTS PLXY PNOM1 **TNI PG** UPCL UPCL **SET** SCURCF STMT GETKEY GETK1 TOGPRT CONC.4 2 MDNG COMC2 CLEAR CLEA1 PNOM1 E MONY PNDM -----0100 STHT 1 590 1673 0674 0675 690 VALU 0296 0000 0287 697A 0721 6A13 6A13 E6E1 07CF 6896 0713 00FF 02C8 68AC 0021 00259 0259 0021 0021 0021 0303 697A 02F3 9020 027 6100 0040 EF 78 0208 E A 1 3 27 0f 136A CFD7 9668 1907 FF F9 ACE8 ONdO EE 00 7AE9 2107 E1E6 CF07 F0 21 20 8 NNULNS 200 0400 02C5 02C5 02C5 02C5 02C5 02C5 02C5 028E 02C1 02C4 8

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-	DECHEX #61 THSET 3 THSET 3 THSET 3 ROL ACT	SET STOP	PNDM	\$20	secs:	# 00 READ # 13 FRZSTI	FRZST2 PACK28 DECHEX FRZTMF	F ZT MEC ROL ACT	UT PACK	READ DUTPUT PACK READ DUTPUT PACK	THEFT	#9FF CLCTBL.	ET UP T	TVTBLL.	TVT BLH.	TVIICS.	#160 T1CST1
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	J CONVERSIONS SPIE - 20 TO TAPE BUFFER FOW LAST RYTE DOWNWARD FFEN- MEECFOING RYTE (11	DSEC MINIPUM IS REGUIRED). EeL 1 TO EEGIA CONVERSION MEL 2 TO EEGIA CONVERSION MEL 3 TO EEGIA CONVERSION	V147 FRCM CLCCK CFIP 37 FIRST FEADING LISECOND ROUTINE	IY FIRST FEADING ICK CHIP (READ STATUS REG) Y FLAG 1110 SEOPUDI 1110 SEOPUDI 112 REGUEST FLAG FEMANLE MARF #S INTERUPTS	LYNCHROM12ATION. Erupts at via Iest flag 11 Interupt routine	5. 0 restration microseconds 0 ream clibo processing 1 clart reseecee 1 clart ecister 0 intenunt counter	EGS INCRFMENTING, OD IT. LLISECONO INTERUPI COUNTER Cond Interupis occurred? Interupis
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OF 4. SAMPLE A TO 0'S. SFEC LFD OISPLAY, UPDATE TV. READ CHANNEL 1 A-TO-D HI BYTE STORE TON DISELAY USE STORE TON DISELAY USE STORE TWANTERGE REGISTER HI FOR 2 BYTE AVERAGE STORE CHANNEL 2 A-TO-D LO BYTE STORE TH AVERAGE REGISTER LI FOR 2 BYTE AVERAGE READ CHANNEL 2 A-TO-D HI BYTE Store for distantion Store in Average Register HI for 2 byte Average Read Channel 2-to-d LD hifor 2 byte Average Store in Average Register LI for 2 byte Average READ CHANNEL 3 A-TO-D HI BYTE STORE TON DISLAY USE STORE IN MARAGE REGISTER HI FOR 2 RYTE AVERAGE STORE IN MARAGE REGISTER LI FOR 2 BYTE AVERAGE STORE IN AVERAGE REGISTER LI FOR 2 BYTE AVERAGE 3 GET THE CURRENT STORAGE POINTER LOYTE IN PAGEL IF NO EXIL INTERUPT RCUTINE Correspondence CLCCX Pestore Figisters and exit interupt routine I PLOT POINT LIF NECESSARYI ON VIDEC DISPLAY GET THE DATA SAMPLE RATE COMPRER AGAINST THE SAMPLE COUNTER I E SAMPLO SHOULD BE STOFFD GD TO SAMPLO I IF NOT EXIT AT STAPDO2 SAME PROCEEDURE AS ABOVE. FOR CHANNEL 2 A contract of the contract of SELECT WHICH OF 4 DATA SAMPLING STEPS IS NEEDED. + PHASE 2 OF 4, SAMPLE A-TO-D+S, AVERAGE, STCRE TREINITIALIZE THE SAMPLE COUNTER .CHECK TO SEE IF THIS POINT IS TO DE STORED. STEPO COMPLETED GD TO STPEND -----... LDA ATDCHI STA DSPCHI STA AVGPHI LDA ATDCHI+I STA AVGPLI + LDA ATOCH3 SIA DSPCH3 SIA AVGRH3 LDA ATOCH3+1 SIA AVGRL3 LOY CURSTR CLC CURSTR CLC A TOCH1+1 ACC ATGRL1 ACC ATGRL1 STA ATGRL1 LOA ATGRL1 ADC ATGRL1 ADC ATGRL1 RDA ATGRL1 RDA ATGRL1 RDA ATGRL1 RDA ATGRL1 LCA ATDCH2 SIA OSPCH2 STA AVGRH2 LDA ATOCH2+1 SIA AVGRL2 CMP #\$01 BNE INTXIT JSR ROCLCK ATDCH2+1 AVGRL2 #\$E0 AVGRL2 ATDCH2 AVGRL2 AVGRH2 LDA SAMRTE CMP SAMCTR DEO SAMPLO JMP STPNO2 LOA #00 LOA \$704 LOA \$7755 CAND #\$01 CAND #\$01 CAND #\$01 CAND #\$01 JMF \$5757 JMF \$5457 JMF \$5457 JMF \$5457 CMP #\$02 JMF \$7672 JMF \$76722 JMF \$76722 JMF \$76722 JMF \$76722 JMF \$76722 JMF \$ PHASE 1 JMP STPND1 JSR TV015P SOURCE SLMT . ADDA ADDA TIXIN SAMPLE SAMPLO SAMPL2 STEPOB STEPO STEP1 STHT 000 010 20029 10000000 021 /VFO 0001 0468 0668 0013004444 A200 0014 0024 A201 A201 A 300 0015 0028 A 301 A600 0016 002C A601 0029 06.90 0693 A201 0027 0060 4200 A200 A301 0028 0028 0028 0028 0028 1000 002A OPND 01 03 9706 13 30 9506 28 AD 0046 85 16 85 20 AD 0146 85 29 00A2 14 24 01 42 00A3 15 28 0142 27 27 27 27 24 24 20 9006 01 A3 28 28 28 00 A3 28 4C 93D6 5 24 8 2204 40L4 4 84 NULUG4UD4 000000 00000 *********** ----04E1 04E3 04E5 04E9 04E0 04E1 Ľ 051100511 A 150 0510 051F 0521 0524 0526 0528 0528 0528 0528 0535 0532

RCE STMT	RCR 5 # AVGRN2 FGR AVGRN2	CL AND	CA ACCENT, I CE CRANEL I AN OVTE AVERAGE SA RESCETTAT E SAVE I IN DATE AVERAGE CA AVERTAT I CET CHANEL 2 MI DATE AVERAGE LA AVERAGAT VE CET CHANEL 2 MI DATE AVERAGE SA RESCALTAT I SAVE I PARAMENT SA RESCALTAT I SAVE I PARAMENT	DNE STEPIC I IN DATHA JIAN STANDAGE DDINTER DINTER IN DATA STANDA I INCREMENT THE STORAGE DAGE DDINTERS IN DATA STANDA I INCREMENT THE STORAGE DAGE DDINTERS	(1) Conversion 1. J. Conversion 1. Conver	2 LOA FRFLG 1 GET THE FREEZE FLAG 960 Steptus 1 Get The Freeze Flag 960 Steptus 1 F Glear, Keep on Sampling, GC TD Step28	LA RETER ET AL VE FELST VIE CLI. STEL VUE 100 SETSP FIE VOLGAL VIE CLI. STEL VUE 100 SETSP FIE VOLGAL VANIED O ST SAMPLA 100 SETSP FIE VIE VIE VIE VIE VIE SAMPLA 100 SETSP FIE VIE VIE VIE SAMPLA 101 SETSP FIE VIE VIE VIE SAMPLA 101 SETSP FIE VIE VIE VIE VIE SAMPLA 101 SETSP FIE VIE VIE VIE VIE VIE VIE VIE 101 SETSP FIE VIE VIE VIE VIE VIE VIE VIE 101 SETSP FIE VIE VIE VIE VIE VIE VIE VIE 101 SETSP FIE VIE VIE VIE VIE VIE VIE VIE VIE 101 SETSP FIE VIE VIE VIE VIE VIE VIE VIE VIE VIE 101 SETSP FIE VIE VIE VIE VIE VIE VIE VIE VIE VIE V	N. DA FZELG CLEAN NE PREZE CLAG SLA FZELG GCTN NE FREZE CLAG SLA FZELG GCTN NE FZELG ALGO LA FZELC GSEST NE FZELG FREEZE THA COUNTER LA FACL 351 THE SAMP FLG STA SPALC PH THOOT I PULL REGISTERS AND PTI	LOA ATOCHI I EFEQ. OLANNELL IA L'DO-MI ATOFI 35 NAGMNI I EFEDE NA NEELAGE RECISTER NI FOR 2 NYTE AVERAGE 10A ATOCHI I EFED CANNEL 2 - FOD-OLO DA TOCHI 54 NAGMLI I STORE IN NYERAGE RECISTER LI FOR 2 BYTE AVERAGE	DA ATOCK 7 READ CIMPLE 2A-TO-DA IDTE 31A AVENT 7 STORE IN AVENCE REGISTER HI FOR 2 DATE AVENCE 10A AVENT 7 STORE IN AVENCE REGISTER II FOR 2 DATE AVENCE 11A AVENT 7 STORE IN AVENCE REGISTER II FOR 2 DATE AVENCE 11A AVENT 7 STORE IN AVENCE REGISTER II FOR 2 DATE AVENCE LOA ATOCH 3 FREAD CHANNEL 3 A-TC-O HI BATE
SOURCE				STEP18	STEP LC	STEP 2		SE TSNP	5TEP 28	
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dD	489 499	14 <i>~~</i> /044686 8090909496	915 915 915 915 915 915 915 915 915 915	200 200	401 00 000 000 000 000 000 000 000 000 0	10 10 10	4FR40048040 N06N2001610	848484 888880	808 908 908	40 400 A
S S	0550 0556 0556	0560 0566 0566 0566 0566 0566 0566 0566	0516 0576 0576 0576 0576	59595	10200	059A	0055755500005575000055750000557550000557550000557555000055755500005575550000555500005555000005555500000555550000	4999999 999999 999999 999999 999999 99999	0300	0505

Ī Ŧ STORE THE SHIFLED HI BYTE INTO AVERAGE REGISTER STORE THE SHIFTED HI BYTE INTO AVERAGE REGISTER STORE IN AVERAGE REGISLER HI FOR 2 BYTE AVERAGE Store Chamber 3 and 2 byte average Store in Average Register LI For 2 byte average READ CHANNEL 2 A-10-0 HI DYTE AVERAGE ADOF T STRELJOUS SAMEE FON 2 DYTE AVERAGE READ CHANNEL LATONG HI DYTE ADOT O PREVIOUS SAMER FON 2 DYTE AVAGE SHIFT 12 DIT PLUS CARRY RIGHT 5 DITS AVAGE ADD IC DREVIOL AT DO HI BYTE ADD IC DREVIOLS SAMPLE FOP 2 EVIE AVERACE REAC CHANNEL A VERACE FOP 2 EVIE AVERACE REAC CHANNEL A TOPOG HI WYTE REAC DO PREVIOUS SAMPLE FCR 2 EVIE AVERACE SHIFT 12 BIT PLUS CARPA RICHT 5 HITS READ CHANNEL J 4-TO-D HI BYTE SADE TSTERIOUS SAMPE FOR Z EVTE AVERAGE READ CHANNEL J 4-TO-D HI BYTE READ CHANNEL J 4-TO-D HI BYTE ADO TE PREVIOUS SAMPE FOR TIGHT 5 BITS SHIFT IZ BIT PLUS CARAFY ATGHT 5 BITS 2002317, 1 640X M0 540X CHANNEL 1 10 171E 11 AM 1000011, 1 61C CHANGE 1 1 641E 2 5AMPLE AVERAGE 1000012, 1 54CX MNG SAVE 1 M REF. 2 5AMPLE AVERAGE 1000020, 1 54CX MNG SAVE 1 M REF. 2 5AMPLE AVERAGE 1000020, 1 54CX MNG SAVE 1 M REF. 2 5AMPLE AVERAGE 1 54CX MNG SAVE 1 M REF. 2 5AMPLE AVERAGE 1 54CX MNG SAVE 1 M REF. 2 5AMPLE AVERAGE 1 54CX MNG SAVE 1 M REF. 2 5AMPLE AVERAGE 1 54CX MNG SAVE 1 M REF. 2 5AMPLE AVERAGE 1 54CX MNG SAVE 1 M REF. 2 5AMPLE AVERAGE 1 54CX MNG SAVE 1 M REF. 2 5AMPLE AVERAGE 1 54CX MNG SAVE 1 M REF. 2 5AMPLE AVERAGE 1 54CX MNG SAVE 1 M REF. 2 5AMPLE AVERAGE 1 54CX MNG SAVE 1 M REF. 2 5AMPLE AVERAGE 1 54CX MNG SAVE 1 M REF. 2 5AMPLE AVERAGE 1 54CX MNG SAVE 1 M REF. 2 5AMPLE AVERAGE 1 54CX MNG SAVE 1 M REF. 2 5AMPLE AVERAGE 1 54CX MNG SAVE 1 M REF. 2 5AMPLE AVERAGE 1 54CX MNG SAVE 1 M REF. 2 5AMPLE AVERAGE 1 54CX MNG SAVE 1 M REF. 2 5AMPLE AVERAGE 1 54CX MNG SAVE 1 M REF. 2 5AMPLE AVERAGE 1 54CX MNG S 4. SAMPLE A-TO-O'S. AVFRAGE, SHIFT. STORE GET THE CURRENT STORAGE POINTER STEPO COMPLETED GC TO STPENO GO READ THE CLOCK SIA AVGRH3 LDA ATDCH3+1 STA AVGRL3 : PP ATDCHI+1 AVGRL1 AVGRL1 AVGRL1 AVGRH1 AVGRH1 AVGRL1 ATDCH2+1 AVGRL2 AVGRL2 AVGRL2 ATDCH2 AVGRH2 ATDCH3+1 AVGRL3 AVGRL3 ATDCH3 ATDCH3 AVGRH3 AVGRL 1 AVGRL I AVGRH1 A VGRL 2 A VGRH2 AVGRL 3 JSR RDCLCK ION STPND1 . PHASE 4 CURSTR A VGRL 1 A VGRL I A VGFL 2 AVGRL 2 A VGRL 2 A VGRL 3 AVGRL 2 A VGRL 3 A VGRL 3 A VGRL 3 STHT DURCE TIMOUT STEP 3 002C 4601 0029 08EB 0693 A201 0027 A200 A200 002A 0027 0027 0027 1000 0027 002A A301 0028 A300 A300 0028 0028 0029 0028 0028 0028 A601 0029 0029 0020 0020 0020 0029 0029 9200 0024 00028 00028 00028 00028 00028 00001 00001 00001 00001 00001 0PN0 2C 01 A6 29 69083 9000 0143 28 28 28 28 28 28 28 28 28 0146 29 2046 2046 29 29 5 202 ŝ 53 5 8 10 4 B 20 ¥ 050A 050A 050A 050A DSE2 0624 1449000 3

CHANNEL INCREMENT THE AME STORAGE DUNIER IF NOT RESET TO NERT PAGE GO 10 STE PT CA SE IT CA STORAGE REGION EXMLSTE C IS NOT STORAGE REGION EXMLSTE C IS NOT STORAGE ACC J CHANNEL STORAGE POINTERS RESET ACL J CHANNEL STORAGE POINTERS P iss same interesting the sam POINTERS INCREMENT THE SAMPLE NUMMER Get the Incremented Sample Mumer Compare with the Disclay and if hot an integral value of Disrif tvexit INCREEM STORAGE POINTER IF MOT RESET TO NEXT ACT TO STPENO SEE IF OX STORAGE REGION EXHAUSTEC IS THIS THE LAST PAGET IS NUT OF TO STREAT RESET ALL 3 CHANNEL STORAGE PAGE POINTE RESET ALL 3 CHANNEL STORAGE PAGE POINTE INCREMENT THE STURAGE PAGE POINTERS INCREMENT THE STORAGE PAGE POINTERS PAGE # UPDATE THE STORAGE OVTE POINTER Increment the Step Lag Increment the Sample counter AT CAN CAN ALL 1. J CANNEL 1. J DUTE AVENCE TAL MORECULT. SCIENCE II. MAN DUTE AVENCE TAL MORECULT. SCIENCE I POINT (IF NECESSARY) ON TV OISPLAY. INCREMENT THE LINE PDINTER Get the Line Pointer A REGISTERS. X+ AND •• MSBCH1+1 MSBCH2+1 MSBCH2+1 MSBCH1+1 MSBCH2+1 MSBCH2+1 .RESTORE V. < CURSTR STPFL6 SAMCTR INPTR SPLAY • D1 S 44 SIMI TAAAAAAAAAAAAA YYY ≥¥2 PLAY PLAY PLAY PLAY ŸŽŽ ¥a STPEND STPND1 SOURCE STEP 38 51EP.3C STPNDO STPEN2 T V01 SP RGET NT+001000000000 STWT VALU 0000123 0651 0002 0015 0669 0007 0007 00015 00015 00015 0000 0001 00025 0032 0 N JO 000 0010101000 000 PROPERTY NO 90 0040404040 000 200 000400 0000 1-0-1-0-0 9 K 0664 0666 0666 0666 0667 0672 0672 1690 0690 069F 06A1 06A3 06449 06449 06449 06449 06449 06449 06649 06649 06649 0685 S

10.1 International and the second sec IS THIS THE FNO DF THE SCREEN? IF NOT GO TO TYOSP2 IF LITIS RESET TO TOP OF SCREEN STORE TO LINE POINTER **eKEYBDARD ROUTINE FROM AIM-65 MCNITORe** TRDONEK TDEBKEY DNEK 2 8904 1005P2 8904 LINPTR KEVCK1 #803 #87F \$5024 \$5024 \$305 \$4480 \$4482 \$5008 \$4482 CWD DNC LOAD 1 41 2 SECANSPA SECON A SPL SOURCE T V05P2 TVEX LT KEYCKD KEVCK STHT 0000000 ALU 00 90 00015 0020 00039 ECEF E02A 000F A480 A482 0 7 4 A E008 1400 OPNO 90.00 20 32 0005 EFEC 2AED 8F 82A4 08 ED 200 A SELLON \$ 0000 068900688 ä

	KEVCOO	KEYCKO KEYCO	8 A 4 2 5 e 4 F F	\$A42A ;DNEKEY	K E Y C 1 B 5 A 2 2 5 A 2 2	KEYGK2 KEYG0	SED2C 1DERK1		8 4 4 2 B	K EYCK 4	KEYCK3 SF421.v		# # 10 KEYCKS	#53F Keycke		# # # 40 Kevck7	a nor	REYCK7	KEYCKS	KEYCK8	#\$10 Kevcke	0988		834F	#500 54477	\$F050	DNVERT HEX BYTE TO 2 BYTE BCD。 BYTE 7D NE CONVERTED IN A register. Result: MSG in Y register. LSB in a registere	# SOO 3 CLEAR THE CARRY PIT #900 3 ASSUME THAT VALUE IS <100
STM	PLA PLA	SN8 SN8 SN8			L NO	500	1SR	No.	LOA		LOA DAR	TXA TXA	NA DE	224	PLA	QN Q		0 KD	PLA		BNE	2 del	A P P P	AND	21 A	122		55
SOURCE			RETCOU	KEVCK1			KEVC18 KEVCK2			кетска	KE YCK4				KE VCK5						NETCHO	KEYCK7 KEYCK9		KF VC 11	KEVC10			нхавсо
STMT	1326					200			1320		1355	1356	1361	1 1 6 9	1365	1969	0451	1372	546	110	13790	1362		1386	1306	1391	465	1396
VALU	0140	072E	A428 00FF	A424 E005	0759 A420	0710	602C		A 4 28	0760	0766	0797	0110	0035		0040	000	0000	2610	0198	00100	0900	00000	004F	0000 A477	F 05 0		0000
DPND	8	55	2844	2444 05ED	2844	:50	2CED		2844	03	FA 21F4		28	35		22	J.	222	5 5	8	010	0.0	00	4	77.44	5010		00
90	400	5823			8898		200		8 Q Q	-	288		Nº.	0 0 U		200	4 0	200						6 0 8		0 0 9		
L OC	4610		1420	1410		0155	8520	0760	0763	1910	0760	0770	0775	8110 8110	0110	1010	0705	0787		0610	5610	1610	9610	0140	07A5	0748		0740

STMT	100 21220 1444 2220 202244.17 202 2120 1444 2220 202244.17 202 2120 1444 2220 202244.17 202 2120 1454 2220 202244.17 202 2120 202 2220 202 2220 202 2220 202 220 220	ро жео зоотку очествить реа ул. и на поста со заче репорятие ул. и на поста со заче репорятие ул. и на поста со заче до ти у ебсисацие на поста поста поста поста и и и и в поста поста поста поста поста и и и и и и на поста поста поста поста поста и и и и и и и и на поста поста поста поста поста и и и и и и и и и и и и и и и и и и и	LC 2000 LC 2010 LC 201	UN TEMPTAR I RESTORE THE X REGISTER (OUTODI POINTER) Oclear Display and restore Old Action code.	ON CLEAR I CLEAR THE DISPLAY DA DLAGT I CLEAR THE DISPLAY COCE In Action I store back display and action code in dougfay I put in the display duffer	*SETUP VIDED DISPLAY PARAMETERS*	an Defan i Clam The Diselv Shown i Feinf Pagnet Fis diselar chance enter Ci diselar	CC 810MMEL	C6 \$20 CC A:	(0.800 SP READ I READ KEYBDARD DP 9000 I IS FIRST CHARACTER A CARRIAGE RETURN? DE VESTOR I FI TI SC OD DIVESED ST OUTPUT IF HOT OUTPUT IT	DU 993] MAS CANNUEL ISLECTED DU 992E1 MAS SO GO 71 192E1 DO 192E1 MAS SO GO AND USEL DO 192E1 MAS CO AND USE CO ED DO 192E1 MAS CO AND USE CO AND USE DO 192E1 MAS CO AND USE CO A	VE IVSET 1 IF NDT, ILLEGAL CHANNEL, REPROMPT VD #40F 567 Nev viine of channel.	COMMENT & GET MEN VALUE UP CHANNEL 50 4001 1 DECEMBER BY I 15 CANNEL 3 AVE AS DATA CHANNEL CFFSET	R CLEMP I CLEAR JHE DISPLAY SP Porm I Print Prompt FDR DISPLAY CHANEL ENTRY C DISPLAY I PRINT PROMPT FDR DISPLAY
	0000000000	9	ຸມ ພ	<i>μα</i> •• •• •	- 2000 -	•••		E E		1007	08080	5-4		
SDURG	GTEZO	GTELO	HBSAV		P CL. AC		1 v 5 E T					TVSFT		TVSETE
1 M I S	00000000000000000000000000000000000000			0100	1929	1426	1420	10.01		\$14 \$16 \$16 \$16 \$16 \$16 \$16 \$16 \$16 \$16 \$16			8448 1440	254
VALU	0000 0789 0764 0701 0701	00C8 07C5 0064	001A		02030		0203			693C 0000 0000 0000	0000	0006	1000	0203
DPND	00440	C507 64	14		20102		0100 0100 4 > L 5 S = 020	VIKZZWJ	<	3CE9			100	602 802 1 5
QD	000000000		898 898 898 898 898 898 898 898 898 898		000000		0040500-00	000-0000	341 S	SNULN	202020	5 6		00400
Ľ	0785 0783 0783 0787 0789 0789	0795 0775 0775 0775	0100	0106	0704		0709 0705 0761 0761 0761 0763	0769 0769 0769 0769 0769 0760	0765 0765	0772	00000	0000	0000	0915 0915 0915 0915

																	-								
			READ KEYBDARD	IF IT IS GD TVSET3 TATA AND FOLDENS	WAS CHANNEL & SELECTEO? IF SO GO TO TVSET2 MASNEL 2 SFLECTEO?	IF SD GD ID TVSETZ WAS CHANNEL 3 SELECTED7 IF NDT 111 EGAL ECHANNEL	GET HEX VALUE OF CHANNEL	OECREWENT BY 3 Save as data channel cffsei	TEMP WILL BE USED IN IMIS SUBADUTIVE	CLEAR THE DISPLAY	PHANT PHUNPY FUR CHART PATE ENTRY						GET FIRST OF THREE DIGITS I VALUE <256) WAS A CARRIAGE RETHON ENTROCOS	IF SD EXT	WAS FIRST DIGIT A ZERGY	WAS FIRST DIGT A LT	MAS FUST DIGIT & 27 MAS FIRST DIGIT & 27 ILLEGAL FIRST DIGIT, REPRDMPT	SET TO 200 IDECIMALI AND GD TD TVSETS	SET TO LOO (DECIMAL & Store in temp	INPUT NEXT 2 DIGITS Convert 10 Heridecimal	ADD #x400 DEC#MAL AS NEEDED
			** **		*****	*** ***	• ••	****	**	••••	•						** **		• •••	• •••	*****	••	****	** **	••
17	3 \$20 CHANNEL	5 \$2D 8:	2 200 2 200	DUTPUT	1 TVSET2	0833 1VSET2	# \$0F	CHANL B	0 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2 2	CLEAR	1 CATA		\$20 BATE1 /01 -				READ	TVSETX	0544	10.0	rvser.	#200 1 1 1 1 1 1	#100	P ACK 2 DECHEK	TEMP
s	55	20	0.55	JSF JSF		BN C	-NA	Ser .		SI			55			- 5	See 1	960			Se y	No.	AT S	SR	28
SOURCE							TVSET2		TVSE 13									-					TVSET4	TVSET6	
1 M 1 S	2 4 5 5 4 5 5 5 5 5 5	1450	1459	463		1467	1410		976	0.41	14.90		1441			1041	484	1400	0.04	0.04	1992	500	1498	2005	1503
VALU			E93C	697 A 5	0032	0033	000	1000	0000	0203							E93C	0000	00.00	1500	0032	0000	0004	0307	6100
OPNO	¢,⊲≻ vi∢zzwj	±0.++	30.69	16 7AE9	800	500	0Ľ		001	0302	-0-	(= «	α	< = w =	· \ a -	••	3069	TAFO	20		200	800	**	0103	6
9	00+00000-0000	0 ~ ~	800	200	201	208	00	0.0	4 0 4 0	0.0	8 4 -		00		1.04	10	201		02	02	28	-	99	200	e e 6
L DC	00019 00019 00019 00019 00019 00019 00021 00021 00021 00021	0824	0828	0820 082F	0836	0630	00036	1480	0845	0840	0841	0652	00554	0856	4220	0650	085F	0000	0.000	0900	0673	5180	9419	0190	1999

STRIM, CELL SET UP, Y FELITER IS PACE OFSET STRIM, CELL SET UN'T CELL STRIM I CELL SET STREE FIRST WITE OF UPE PACE? ERST WITE OF UP ACES ALL VEGOV PACE? CLEARED? 3 INDEX POINTER TO DATA WORD {W.DI.D2.H.W....1 1 SELECT WORD TO OF READ VESABLE MEMDRY RUNS FROM \$1000 TO \$7FFF USE DESTRUCTABLE PAGE ZERO TEMP REGISTERS SAVE THE DISPLAY RATE GO CLEAR THE VIDEO DISPLAY DRAW TICS DN THE NDW CLEARED SCREEN Restore the OLD Action Code I SAVE THE A AND X REGISTERS DN STACK 1 SEND "GD" CDMMAND TD CLCCK CHIP
1 RESET SECS.XXX=00.000 INDEX TO CLOCK BUFFER Get Each Value Select Data WORD to be written Sette Data WORD to be written Get Wext Value 1 STORE DATA IN CLOCK BUFFER 1 GET NEXT DATA WORD FROM CLOCK *** BRITE TO CLOCK (START OUTBOARD TIMERI*** # RESET CLDCK FDR READING I IF NDT. KEEP ON GOING # SET CLOCK FOR BRITING
SET FCR 10 SECCNDS
SELECT SECDNOS WORD INITIALIZE CLOCK FDR BRITING · INITIALIZE CLOCK FDR READING+ VIDED DISPLAYS PREAD FROM CLOCK+ TVCLEA 1 14 LISS 4410 SIX 50 SIX 5 PHA TXA PHA LOX #507 1 51X VIAORA 201 VIAORA SIA WILSEC.X 061 STA DISRTE JSR TVCLEA JSR TICSET JSR ROLACT RTS LCA PAFF STA VIADR8 LDA #89F BNE RDCL12 SOURCE STMT ARCL IN PHA TVSFTX ROCL IN ROCL 12 GDCLCK GOCL C2 PDCLC2 ROCL CK 87489×840-2025 8000 FR ----VALU 0031 0692 03E1 03E1 00000 00000 00000 8000 AF02 0018 AF0C 6100 00000 00FF AF 02 009F D8B 3 00002 0010 0002 AF01 AF01 AF01 AF01 AF01 AF01 AF01 AF01 AF01 00002 00002 00002 00002 00002 00002 00002 00002 00002 05 31 20 9208 20 6103 20 CF07 OPND 48 49 00 49 10 80 024f 68 00 4F 68 00 4F 68 00 4F 48 49 FF 80 02 AF A9 9F 00 EF 20 8408 42 02 42 02 86 014f 86 014f 85 09 85 09 85 09 86 004f 80 014f 20 003f 20 003f 20 003f 20 003f 01 AF 01 AF 00 AF 09 90 040 NHON4 00000 ŭ

	1 RESTORE REGISTERS E Clock Frich Reyddard Commance	I CLEAR THE DISPLAY PRINT AND DISPLAY CLOCK START PRIMPT			I GET KEYBDARD ENIRY ISIT A CARAGE RETURN7 ISITA STREM RTS IFITISTREM RTS GD MRITE TO THE CLOCK	I TEMP MILL BE USED IN THIS SUMPOUTHE Clear Hite Display I paint puompt fer data rate entry			GET FIRST OF THREE OIGITS I VALUE (236) Vas Acartage Return entereo?	DUTPUT FIRST DIGIT ENTERED I WAS FIRST DIGIT & ZERCT I WAS FIRST DIGIT & ZERCT	WAS FIRST DIGHT & 17 1 FS SO GOT PO DRSE4 1 WAS FIRST DIGHT & 27 1 LLEGAL FIRST DIGHT, REPRCMPT	SET TC 200 IDECIMAL) AND GD TC DRSETS S SET TD 100 IDECIMALI	\$ STORE IN TEMP \$ INPUT NEXT 2 DIGITS
	RDCLC2	CLEAR PNOM START=A		5 2 0 E X 1 1 = CR	\$00 REAU #13 TIMGD2 GDCLCK	FEND TENP CLEAR PNOM 0ATA	\$20 RATE	#57PT:	READ READ READ	0UTPUT #\$30 DRSET6	085ET 4	#200 DRSET5	T EMP P ACK 2
STM	PLA TAX PLA PLA	1 SR 1 SR		FCC	128 128 128 128 128 128 128 128 128 128	158 JSR JSR	FCG	FC8	S S S S S S S S S S S S S S S S S S S	5100	See a	LOA BNE	STA STA
SOURCE		11 MGD			114602	OR SE T						DR SET 4	ORSE 15 ORSE 16
STMT	1576 1578 1578 1579 1561 1561	15654		1568	000000	1599	1601	1603	1605	1010	0000	1017	1620
VALU	08F 0	0203 0268			E93C 000D 0521 0521	0000 0019 0203 0268			0000	E97A 0030 0558	0032	0000	0019 D3C4
CPNO	5	0302 6802 5	< a + s < z >	u×≈⊢t⊌α	3CE9 00 03 C408	00 0302 68C2 68C2	< α < ⊢ u		3CE9	7AE9 30	1000	800 800	19 C4D3
9	494 99 99 4 99	0000	4000 440	098040 m N	0000000	880084 88008 800800	00-45	00000000	8000	000	5208	6 0 6 4 6 0 6 4	000
L OC	0859 0850 0850 0850 0850	0902	10000000000000000000000000000000000000	0000 0000 0000 0000 0000 0000 0000 0000 0000	0916 0917 0916 0916	0924 0924 0924 0926 0920 0920 0920	1600	50000000000000000000000000000000000000	0660	4460	0940	0953	0959

DURCE STWT	FCC MAILE	FCC COMPLETED	FCB #400 JSR CLENUP : CLEAR THEN RE-ENABLE 1/10 SECCHD INTERUPTS AT VIA JM MAIN : =>> GC TO MAIN =>>>	* * * * * * * * * * * * * * * * * * *	LENUP LOA #199 : RESET MALF MS. INTERUPI COUNTER Lox #10 : Clear Interupis defore re-emelinc Six Misma : N-1-A. Adoress register	LOA MARMA IV-LAA UNIA MEGISTAL LDA MARZ IEMALE I/10 SECONO INTERUPTS AT VIA STA VIALER I RESTORE OLO ACTION CODE JSR POLACT I RESTORE OLO ACTION CODE	RTS FRANSFER CONTIGUOUS BLCCK OF DATA TO TAPE®	31APE ČCA #354 ; SET OUTPUT DEVICE CODE=+1+ (TAPE) Šta outfla	.TRANSFER ENDING AND STARTING ADDRESSES TO ADDR AND SI.	LOA TPEND CETLCBTE OF ENDING ADORESS Stadda Store Offe of ENDING ADORESS USED NY WRIAPE Stadda Store Offe off Stadthad ADORESS USED NY WRIAPE Stadthad Stadthad ADORESS USED NY WRIAPE	LOJ TOFENOL I GET MI BYTE OF FROMA DODRESSEUT THATE LOA POSTAL I STORE MI DYTE OF ENOLVADORESSEUT AFFEND LOA POSTALI I STORE MI DYTE OF ENOLVADORESSEUT AFFEND STA SILI I STORE MI DYTE OF ENOLVA ADDRESS FOR TAREW	JSR NAMO \$ GET THE FILE NAME	UN 401 SET INVUT FLAGA I I TPE ULUEUT) JAT TATZ I SET INVUT FLAGA NO MAITE FOR USER RESPONCE JAT CHARLE I DEMA TLE FOR UNUTI I DI AFE NY SUCKS JAT CHOW I CARRIAGE RETURN LINE FEED ID EISKLAY	LOA #00 i CLEAR THE RECOFD CCUNT STA S21 STA S24	LOV #01 5 OUTPUT ONE MDFE BYTE JSR NXTADD 5 ADD Y TO ADDR+1 ANO ADDR	JZ JSR CRLF ; CR/LF TO ANY DUPUT DEVICE ITAPE)	CALCULATE NUMBER OF BYTES YET TO BE DUMPED.	JSP CLCRK ; CLEAR CHECKSUM Lea aodr ; choing aodress - current aodress sec
201					CLE CLE			191								DU2		
STWT	1690	1691	1693	1691	1700	17054	1705	1712		1110	1721	524	1729		9611	5671	142	
VALU			094E		0000 0010 AF 01	0082 AF 0E 07CF		0054 A413		0100 9410 9410	0016 0010 0010 0010 0010 0010 0010 0010	EBCF	0001 E8C5 E56F EAL3	000000000000000000000000000000000000000	0001 E2CD	E9F0		684D A41C
ONIO	<cu =c-+u<="" td=""><td></td><td>AED9</td><td></td><td>C7 00 01 AF</td><td>82 0EAF CF07</td><td></td><td>1344</td><td></td><td>10</td><td></td><td>CFEB</td><td>CSE0 CSE0 LJEA</td><td>000</td><td>01 C0E2</td><td>F0E9</td><td></td><td>4068 8044</td></cu>		AED9		C7 00 01 AF	82 0EAF CF07		1344		10		CFEB	CSE0 CSE0 LJEA	000	01 C0E2	F0E9		4068 8044
dD	40400F00404	0 M L 0 0 U M 4 0 4 0	80V		0 4 0 4 C 4	22 8 N	0	4 0 B		80 S C C C C C C C C C C C C C C C C C C	4040	50	N000 V N N N	800 800	80 N	8		000
L OC	49990 99900 99000 99000 99000 99000 99000 99000 99000 99000 99000 99000 99000 90000 90000 90000 90000 90000 90000 90000 90000 90000 90000 90000 90000 90000 9000000	00000000000000000000000000000000000000	0940		0984 0982 0984	V0600	0.002	0903		80000 80000	0000	0900	090F 09E1 09E1	096A 096C	0.96.2	0.9F 7		09FA 09F0 0A00

SOURCE STMT	580.51	PHA	ANE DUG I NUMBER OF DYTES HI	* SEE IF 24 OR WORE OVIES LEFT.	PLA I NUMBER OF BYTES HI MAS ZERO	ACC 000 IS THE NUMBER OF EVICE > 241 ACC 000 INC. ONLY DUTPUT REMAINED THES ACS 007 TYES. 24 OVTER IN NUMBER SACONCESS	DUG PLA	U7 LOA #\$24	UB PHA JSR SEMI î OUTPUT ";", NUMBER CF NYTES, ANO PUSH A REGISTER	STA COUNT & SET NUMBER CF BYTES JSR OUTCK & OUTPUT NUMBER OF BYTER	JSR DUTCK : OUTFUI AODRESS	SUTPUT DATA.	U9 JSR DUTCKS 1 GET CHARACTER SPECIFIED NY 51 Loa #00 ; 1 CLEAR DISKLAY POINTER Sta choor ; 1 CLEAR DISKLAY POINTER	JSR A0051 1 INCREMENT 51+1. AND 51 DEC COUNT 1 DECREMENT BLYE COUNT DEC OUD 1 NOT OCKE WITH THIS RECORD	.OUTPUT THE CHECKSUM.	LOA CKSUM+1 JSR DUTCHT 1 CHECKSUM LO RYTE LCA CKSUM	JER DUTCKI I CHECKSUM HI BYTE JER INCSZ I INCREMENT VERTICAL COUNT	10.2 JSR 16460 I NUTLAST RECORD. LET MONITCE CLEAN UP 715 ALL OUTPUT LAST RECORD. LET MONITCE CLEAN UP	* *CONVERT ONE BYTE HEX 100-OF1 TO ASCII 130 TO AT1	HXAS CMP #30A BCC LT10 1 1F <10 GO TO LT10	ADC #537 \$ AOD 37 FUR 541 ETC.	10 CLC ADC #530 1 ADD 30 FOR 530 ETC.	<pre># MOVE LEFT MALE WORD OF A REGISTER INTO RIGHT HALF WORD. THEN ZERD OUT THE LEFT HALF WORD.</pre>	WOFA AND #3FO # ZERO RIGHT HALF WORD	LSR I SHIFT LEFT MALF WORD TO RIGHT MALF WORD LSR	R.SR RIS	OECIMAL CHARACTER TABLE USED IN CLOCK AND DISPLAY ROUTINES.	TRL FCB \$00
1 2	9		0.0	200	4 10 V		5 50	5	3	- e c	- 0 - N	n e i	2					no		180		5		LHW				5
1 51			11							000				221	00	170	611	044	1 79	1011	1799	0001	0000	1 009	1010		1015	1.1.0
VAL	141	144	DALE		0452	CALT		0 0 2 4	E9BA	A419 E530	E538 A41A E538		E531 0000 A415	6550 4419 0430		A41F E53B A41E	E 5 6 6	EAEB		000A 045E	0037	0030		0 OL 0				
OPNO	1444	1044	5		24	01		24	9AE9	1944 3865 1844	30E5 1 A A 4 3 0 E 5		31 E 5	19 A 4		F 44 19 F 5	1001	0E4			~	0		•				
g	08	4 8 8	8		200	200	69	6 V	000 •N0	000	8 9 0 8 9 0		0 × 0			000	204	200 200		00 000	- 	500		1 8 4 1				2
Γo	DA01	0405		0400	DALO	ALA D	0V10	OA17	0419 0419	0A1E 0A21 0A24	0427 0424 0420		0430	9738				455			054	ASF		40%	1000	468	0 0 00	

4D VALU STMT SOURCE STMT		1910 I T CHANACIEN TABLE FOR VIDEC 11C MANKS **	1920 TVIICS FOO SFFF 1921 FOO SFFF	1022 FOB SFFF 1923 FOB SFFF 128 FOB SFFF	1926 F DB \$FFFF	1928 FOUS SFFFF 1928 FOUS SFFFF 1929 FOUS SFFFF	1930 F08 SFFFF 1931 F08 SFFFF		1011 1011 1011 1011 1011 1011 1011 101			1942 FDB \$0000 1943 FDB \$0000	1945 FDB \$0000		1944 FDB 30000 1951 FDB 30000 1951 FDB 30000	1951 FDB \$0000 1954 FDB \$0000 1955 FDB \$0000	1956 FDB 80000 1957 FDB 40100	1959 FDB \$0100 1960 FDB \$0080 1961 FDB \$0080
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÷		02.4	131				142	11	242		194	555		80.00	198	****		123
SOURCE ST		FCD	28	100	200	001	000	100	101	000	101	101	100	000	222	2220	200	FCB
STMT	500-NN45522400-NA4502200-NA450 	1 647	0000	1651	629	924	1656	1961	1864	0000	1868	1970	1873	1875	0101	1002	1986	1000
VALU																		
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r oc	00000000000000000000000000000000000000	0 48.7	0489	DA80	DABE	0640	0492 0493	0495	0497 0498	× 6×0	06¥0	DA9E		DAA3	DAA5 DAA5	CAAS CAAS CAAS	DAAF	0480

	2 5555505555555555555555555555555555555	
	5 4 4 5 5 5 5 5 5 5 5 5 5 5 5 5	
	D YY X	
	□ DONAP 4 AF = 400 EMp J 35 4 0F V 4 4 0 Am = 40 8 25 0 0 N N F 4 1 0 0 0 0 2 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0	」↓ = = = = = = = = = = = = = = = = = = =
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r oc		00000000000000000000000000000000000000

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UNDU .	
do	00000000000000000000000000000000000000
1.00	
	00000000000000000000000000000000000000
SBURCE STM	<u>20020092920000000000000000000000000000</u>
1 STMT	8588-111195195-00000000000000000000000000000
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CPN	
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START.

SOURCE STM	555595565555
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V AL U	
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L UC	
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1 SC	こうち ひのし そうち ひっしょう ちょうしょう ちょう かうしょう ゆうか かつ ローミ うゆび かつつ ミシリル ちゅつ ニ シッチ ちゅつ ニ シッチ ちゅう 二 シッチ ロック ローマリ
534	
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LOC	
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SOURCE	
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f, OC	

11 36 0573 0574 0589 0605 0611 1423 0782 1132 1134 1 0747 1420 1428 1452 1478 1584 1598 0248 0275 0472 D\$13 D517 D521 0525 0529 0533 0545 0324 D532 0962 1422 0560 0571 0726 0756 0450 0471 0485 0559 0996 1270 1265 1002 1009 1120 1222 0508 0528 0680 1730 0639 0646 16 10 REFERENCES VALU **DBMAS**

1498 1503 1513 1515 1510 1521 1522 1597 1620 1625 1415 0290 0392 0960 9701 0410 0498 0499 0914 0916 0917 0709 0703 0926 C920 0921 1556 1562 1572 1702 1557 1563 1573 1703 0729 0255 0327 0511 0617 0620 0716 0738 0038 0053 1262 0836 0051 1260 0842 0538 0527 0936 1715 1720 1718 1722 1505 1552 0251 0326 0520 1438 1705 0500 0923 0924 1460 19591 0127 0607 0142 0142 0825 55500 55500 55500 55500 55500 55500 1543 REFERENCES 0764 1467 0290 DEFN VALU 9 < m b h b m m 2010 n 000 n 0 UUFCH YMBOL

I STATEMENTS FLAGGED IN THIS ASSEMBLY

APPENDIX B LISTING OF THE 6502 PROGRAM OCCTRANS

This appendix contains an assembly language source code listing for the 6502 machine language program called OCCTRANS. This program is discussed in Chapter 4. This program was assembled with the ASM6502 cross-assembler mentioned in the introduction to Appendix A.

Market and a second state and a	Storbitic OF CE, AND MALTING LF. MAITING 15 OUNE IN AND SET LODP. INTERAPPIS ARE NOT USED. OTH. CTS. RTS. AND OSP SF2-22 STORMES ARE GOUREO.	<<< PAGE ZERD LOCATIONS >>>	GNG \$00 BYTE FCD \$00 ; LEFT MALF WORD OF ASCII RYTE UTTE FCD \$00 ; RIGHT MALF WORD OF ASCII BYTE	016 186 016 1860 1 CODOS P-REGISTER U3 (FILE NAME SPECIFICATION)	CENB FCB \$60 ; COODS SERVICE CALL PROCESSOR ENABLE RYTE	<pre></pre>	IAO EQU SAADO 3 ACIA COMMAND, CONTROL AND STATUS IAI Equ saadi 3 Acia data	<< <aim-65 monitor="" routines="">>></aim-65>	EQU 4E704 1 PRINT AND DISPLAY A DUESTION WARK HEK EQU 4E907 3 SCAN KEYBDARD, IF KEYDOWN PUT ASCII IN A-REG ELSE 4FF	LOW EQU SEA15 UNITOUT A CALVER TO PRINTER JUSTALY LOW EQU SEA15 UNITY CALVE TO PRINTER DISPLAY A EQU SEA46 CONVERTS 2 HEX #*5. IN <a2. and="" ascii.="" output<="" th="" to=""><th>DI EQU SFEGG I INPUT I CHAR FROM KEYBOARD, PUT ASCII IN A-REG</th><th>ORG \$0300 1 RAM SPACE \$0300</th><th>**** COLD START PROGRAM INITIALIZATION ****</th><th>TART ŠEI I DISAGLE INTERRUPTS DN THE MICROPROCESSOR RUS CLO i Clear decimal node</th></a2.>	DI EQU SFEGG I INPUT I CHAR FROM KEYBOARD, PUT ASCII IN A-REG	ORG \$0300 1 RAM SPACE \$0300	**** COLD START PROGRAM INITIALIZATION ****	TART ŠEI I DISAGLE INTERRUPTS DN THE MICROPROCESSOR RUS CLO i Clear decimal node
-~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~	* 5 % ~ 0	00-	19	100	2000	225	4 4 4 4 6 6 4		0.00	8529 1823	122		00	22 CS
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			000	000	000				E 704		FEG	0.00		
			00	0002	80									7.6
			0000	0086	0000							0300		0000

I REAO A CHARACTER FROM THE KEVROARD I SI TA CTRL-I (FROMMIT DATA FILE) I SO BAIT UNTIL FILE TRANSMISSION COMPLETED CONVERT HEX TO ASCII GET ASCII REPRESENIATION OF HI PART OF BYTE TRANSMIT IT TO THE HARRIS COMPUER GET ASCII REPRESENTATION OF LO PART OF BYTE TRANSMIT IT TO THE HARRIS ENABLE THE CODOS SERVICE CALL PROCESSOR ACLL > CHARTER DUTPUT TO DISPLAT LLINE END INDICATION) ACLA CHARACTER FROM KEYODARD DD A CALLE DD A CALLE CHUECK FOR USER ENTRY CHECK FOR USER ENTRY GET A CHARACTER FROM THE ACIA 15 IT A LINE FEED IRECORD TERMINATOR ? 15 NDT GET NEXT ACIA CHARACTER Ξ 15 IT A CTRL-F (SPECIFY FILE NAME I IF SO GET FILE NAME IF SO GET FILE NAME IF NOT THEN CHRCK KET CLOSUPE AGAIN IF NOT THEN CHRCK KET CLOSUPE AGAIN SEED CHANNEL NUMBER FOR CODDS NDT A RREAK ISSUANCE OF SVC #15 ASSIGN CHANNEL #2 TD NAME OF ILE HANGLE EFRADE IF FILE NOT FDUND CHAPACTER COUNTER NDT A BREAX 1530ANCE OF SVC #903 INPUT NYTE FROM ASSIGNED CHANNEL HANOLE END-OF-FILE CONSITION FILE MAME BUFFER POINTER EIN EXT CHARATER OF FILE MAME STORE IN BUFFER OF FILE MAME STORE IN BUFFER OF FILE MAME VERACY OF ENTERPO VASA CA ENTERPO VASA CA ENTERPO DO A CALF "SET UP ACIA FOR BAUD RATE, STOP BITS, PARITY. FILE NAME TO BE STORED AT \$0200 β INCREMENT CHARACTER COUNT PRINT PROMPT *FN2* . CHECK KEYPDARD COMMAND. DD A CP/LF -----.... ---------------... LOA #00 51A PREGU3 LDA #802 51A PREGU3+1 #00 RE01 \$0200.X LDA #180 51A SVCENN : .CHECK KE1 158 RE01 CMP #114 # 806 F NAME # 800 CMDCHK RDACIA #100A READLP #15E 0UTPUT CRLOW UUTPUT CRLOW UUTPUT CRLOW # \$00 GET NAM C PL OW LOA #503 574 ACIA0 LDA #511 574 ACIA0 LOA #580 LOA #580 #1446 0UTPUT #14E 0UTPUT 0M #815 NOF ILE HKA5C H187TE XM17 L007TE XM17 CRLOW 2014 \$\$00 #803 E0F 45 XGXXLWG ND ON D - Standard 19891 ž SOURCE CMDCHK CMDCK2 ACLAIN REAOLP GETNAM CCDUNT SVCIN FNAME 0011210012100117 01139 V M.U 0003 0011 0011 0010 0010 0010 0080 00EE 60000 FE 96 0000 0356 EA LE 01100 EA13 0.30D 0.00A 0.32F 0056 6978 6813 6978 0320 0046 E97A 004E E97A E704 0005 86.00 0000 0.382 46 78E9 78E9 19E7 00 00 0 Å 0 55 7469 7469 7469 7469 2003 00 F5 13EA 0404 00 C 000 01 OPPINO 20 13EA 00 96FE 0002 0013 1961 8582 84 8101 1 ŝ 8 ş ä 200 0200 828 NAPECSO đ 0000 4240 ****** o si Ŷ 40000 4 N 4 N N NONO VOID 4000 00000 00004 80 0305 0300 0315 0310 0326 0326 0328 1125 0330 0347 0345 0345 01500 0366 0366 0366 0366 036A 0366 0366 011000110 32C 336 5

	A40 : HAVE 72 IOECHMAL) CHMMACTERS BEEN SENT GETUAT : GET AND SEND NET BYTE GETUAT : GET AND SEND NET BYTE SEND : SEND A CR TO THE MARRIS IEND OF LINE) XMIT	KOACH INNOVE ON RECEIPT OF CA BY MARRIS MATTIS IN ON SPICING FACEIPT OF CA BY MARRIS WATTIS IN ON SPICE ON MARGANCE RELOW IN A MATTIS ON A CAPACITY OCAY INTO Y ANTIS COORDS COUNT INTO Y ANTIS COORDS	COOJS FILE ERROR CONDITION - FILE NOT FOUND.	OUTPAT I PRINT ******* OUTPAT OUTPAT OUTPAT OUTPAT	CRLOW I DO A CR/LF CMCCHK I CHECK FOR USER COMMAND	END-UF-FILE ENCOUNTERED.	GARS : PRIMT ROF GUTPUT GUTPUT GUTPUT GUTPUT	#05 1 CODOS CHANNEL MANAGER 1 NOT A OREAK CODOS SERVICE CALL #16 1 FREE THE ASSIGNED CHANNEL	CRLOV : DO A CR/LF CMOCHK : CHECK FOR USER COMMAND	TRANSMIT AN ASCIT BYTE IN THE A-REGISTER #	ACIAO I SANE THE APRESITED ON THE STACK ACIAO I GET THE ACLA STAUDS REGISTER IS EMPTY 402 CHECK TO SEE IF TRANSITI REGISTER IS EMPTY 402 CHECK TO SEE IF TRANSITI REGISTER IS EMPTY 902 THE ADD TAULTURIL II IS 0564 M ANIT 3 MILLURIL II IS	ACIAI I GET THE HEY DYTE ACIAI I GET THE HEY DYTE ACIAI I PUI IT IN THE ACIA TRANSMIT REGISTER	REAO AN ASCII BYTE FROM THE ACIA RECEIVE REGISTER +	AGAU I GET ACIA STATUS REGISTER AGIA I CHECK RECELVE BIT ROACIA I FEWPT WAIT UNTL. CHARACTER RECEIVED	ACIA: 1 GET THE CHARACTER FROM ACIA DATA REGISTER	APAN 1 15 17 11 MAS A LA Rockii 16 17 MAS Ekit Without Duiput Output 1 Output Received Character to Display	WAIT IN A DELAY LOOP (5005 MICROSECONDSI +	REF I OLLAY LOOP COUNTER 18FF I 6-MICAOSCOND INSTRUCTION RFFF IS ROW 18FFF
SIM	LOAD	Stanser.		San and a second		÷	SA SA SA SA	ŏžĐ	S &	•	102478	444	•	4 D B	8 2 1	252	•	888 888
SOURCE		SHITA		NOF 3 LE			2010							BDACLA L		B ROEXIT A	****	DELAY2 R
IMIS	2410 2410 2410 2410	015100152	0150	0152	0110 0110 0110	0172	0175 0175 0175 0175 0175 0175	1910	1010	0107	01 00 01 89 01 90 01 91 01 92 01 92	0195 0195 0197 0197 0197	0199	0201 0203 0203	02 05	0209	0213	0215
VALU	004500360	0300 0004 0388 6413 0351 0351		002A 697A 697A 697A 697A 697A 697A 697A 697	EA13 0310		6978 6978 6978 6978	5000	EA13 0310		AA00 0002 03CC 03CC	10 44		0000	A 4 01 007F	0.3f 0 E 97 A		6645 8645 8645
DN40 4	8 48 0 69 0 60 0 60	5003 59 5103 5103 5103 5103		2A 7AE9 7AE9 7AE9 7AE9	1364		7469 7469 7469	6	1 003		00AA 02 02 17 1103	01 V V		00 A A 01 69	01AA 7F 0A	0.J		1111
õ	0004N	NU0004		ANANANAN	* *		4 N 4 N 4 N 4	89	20		4000 2000 2000 2000	000		4 0 0 V	4 % Å	500 200		8 9 9 9 6 6 9 7
Ľ	037E 037F 0381 0381 0383	0388 0388 0386 0385 0385 0385		0000 0000 0000 0000 0000 0000 0000 0000 0000	0340		0384 0384 0389 0389 0386	0300	0305		0305	0309		0360	0369	0360		0353

	LOOP COUNTER K WITH 2-WICROSEC INSTRJCTION	DDD5 CHANNEL NUMBER	AND SAVE IN HIBYTE/LOBYTE .	BYTE DN STACK Hu Right Hu		O ASCII AND OUTPUT	TYTE FROM STACK T HW D ASCLI AND DUIPUT D BYTE
	DECREMENT ZERD CHEC	RESTORE C	TD ASCII	SAVE HEX GET LEFT SHIFT TD		CONVERT T SAVE AS P	GET HEX D GET RIGHT CONVERT T SAVE AS L
	**	-	ΗEX				
	8FFFF #800	DEL AY 2 # 105	CDNVERT	#87.0		NDUT	#\$0F NOUT LOBYTE
STMT	ROR CPXX	BNE R15	•	HA NA	583	STA .	AND JSR 518 R15
SOURCE				HXA SC			
STMT	0213	0220 0221 0222	0223	0226	0229	02.32	0235
VAL U	FFF	0005		0 0F 0		6451 0000	000F E A 5 1 0 0 0 1
UPND	6 F F F	62 05		f o		00 00	0F 51 L A 01
ď	9 4 0 9 4 0	883		997	::::	50 50	60000 00000 00000
LOC	031.0	0.401 0.401		4040	0000	040F0	01140

SYMBUL VALU DEFN REFERENCES

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																		74.1	2												
																			2												
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																			010												
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601		1	6	1	2														11.7												
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000			010		012														5				5			5				5	
0003	0205		0000		0110		6143						0233		02.30		0237		0111	0076			0103			0100				0142	
1.00		156	860	112	2		155	220	136	197	147	125	130	136	Ξ	2	232		108	074	110		203	208	105	003	0.83		153	140	1 92
9	0	•	•	•	•		•	•	•	•	•	•	•	•	•	•	•		•	0	•		•	•	•	•	0		0	·	Č
0000	1000	5113	00 92	6003	0062	1100	0214	0215	6173	+110	0134	0121	0043	0226	**00	01 60	0004	0003	1900	0047	0059	0000	0201	0210	0103	0065	00500	1000	0151	01.88	018%
< 0		<	c	•	'n	•	_	n	N	~		ç	0	4		2	-	•	4	2	4	2	<u>o</u>	•		2	w	•	38	0	2
0.30	A A O	036	100	0.32	EAL	000	0.35	P.CO	950	4034	036	0.35	000	040	000	610	EAS	EAA	5 4 3	000	670	6.90	000	10.0	0.32	ŝ	00	001	0.36	õ	5
z,		ĩ	ž	2×	,	ä	*	72		ų	2	H	2		-	2	l.		h	5		×	4	1	CLP.		98	z	Ŷ		
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~ *	•	-	v	-	-	-	-	-	-	*	-	-	-	-	-	-	-	-	-	-			-			1	1		1	2	

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ND STATEMENTS FLAGGED

APPENDIX C LISTING OF THE APL WORKSPACE OCCPREP

```
V Y+BIT12 X
[1] AUNPACKS 12 BIT BIT-PACKED LODAS DATA
[2] Y+21φ(φ12..66666666667×ρX)ρφ(8ρ2)ΤΧ
  V IDENT
[1]
    1 1
[2]
    * *** LUNAR OCCULTATION DATA PREPARATION WORKSPACE **
    +1
F 3 1
    .
           VERSION: A02/HARRIS'
[4]
    1
           REVISION DATE 18 JUNE 1984'
F 5 1
           NOTE: □IO+0'
[6]
[7]
     A GLENN SCHNEIDER DEPARTMENT OF ASTRONOMY
[8]
     A UNIVERSITY OF FLORIDA GAINESVILLE, FL 32611
   V READ:Y
[1] AREADS IN A LODAS DATA FILE FROM A SPICA-IV/CODOS DIS
   K
F 2 ]
    ATHE SPICA-IV MUST RUN THE PROGRAM "HARRISTERM"
[3]
    ATHE DOWNLOADED FILE IS IN THE TEXT VARIABLE "DATA"
[4]
    DATA+10
    L1:' '
[5]
[6]
    →(V/(2 3p' ∈ T L EN D') ∧ .= 3 † Y+U)/L2
[7]
    →L1, pDATA+DATA,Y,(72-pY)p'*'
[8] L2 : ERRS + ' * ' + . = DATA
    Π
   V TRANSLATE X:D:O:M:A:C:F:Y:DIO
[1]
    APREPARES A DOWNLOADED LODAS FILE FOR APL USE
[2] [10+0
[3]
    'TIMESTAMP: ',(TS X),' DATA SET LENGTH: ',(*pX).' BY
    TES.'
[4]
    D+256+Y+16 161$((.5×pX),2)p'0123456789ABCDEF'120+X
    F+1+L(-2048+D[2 7]+.×256 1)+1.5
[5]
[6]
    Q+1+, 0' ', 9 2p(<u>HEADER</u>+512+20+X)[18+118]
[7] M+12 3p'JAN FEBMARAPRMAYJUN JULAUG SEPOCTN OVDEC'
[8] D+DATE+Q[6 7],' ',M[-1+*Q[3 4];],' 19',2+Q
```

. V APPENDIX D LISTING OF THE APL WORKSPACE: OCCRED

V B+W BBDY I. [1] AB=NORMALIZED RADIATED BLACKBODY POWER CONTRIBUTION [2] AL[1]=ASSUMED EFFECTIVE COLOR TEMPERATURE OF THE STAR [3] AL[2]=NUMBER OF TERMS USED IN THE NUMERICAL INTERGATI ON Γ41 A WEWAVELENGTH LIST B+B++/B++/.000037412+(W+5)×*-1+1.43879+L[1]×W+.000000 $01 \times W \circ .+ (0.1 L[2]) - L[2] \div 2$ Δ ∇ R+CM X:V [1] ACORRELATION MATRIX FOR THE MATRIX <X> [2] $\nabla + (\nabla \circ \cdot \times \nabla + 1 \quad 1 \otimes R + (\otimes R) + \cdot \times R + X - (\circ X) \circ (+ \neq X) + 1 + \circ X) + \cdot 5$ [3] APATCH FOR DIVISION SCALING PROBLEM ON THE HARRIS Γ47 V[1:1]+R[1:1] [5] R+R+V Δ V DC P; DI; CVEC; ITER; N; I; X; V; D; R; Y; S [1] ADIFFERENTIAL CORRECTIONS PROCEDURE FOR LUNAR OCCULTA TION S [2] AFIVE PARAMETER MODEL, FIXED LIMB DARKENING AND TEMPE RATURE [3] SOLS+1 14pP, SSE+0pDI+(5p0), 1.02, (4p1.0003), 1 Ē41 CVEC+ 1 \$\$ + 15 + ITER+0 Ē 5 Ī R+R++/R+(+FILT)[;2]×(+FILT)[;1]BBDY P[4],(1+-+02 2p+F ILT) +5 Γ67 S++/2×P[11]LDARKEN GRID P[5] [7] L0:X+(2+P[3])00 [8] +5+<u>SOL</u>+P,0p<u>SSE</u>+<u>SSE</u>,[]+(Y+<u>OBS</u>-<u>COMP</u>+P[8]+(-/P[7 8])×R WI DE P)+.*2 Γ97 DI[6]+1.02++/(P[6]<1 2 3×0+180×3600000)/.2 .12 .06 [10] L1:P[I]+DI[I]×1E¹²[P[I+1+CVEC+1¢CVEC] D+-COMP-P[8]+(-/P[7 8])×R WIDE P A +((5,(P[12]>ITER)/3)∧.≠1+pX+X,D+-/(DI[I],1)×1E⁻12[([12] P+SOL)[1])/L1 [13] +(5≠1+pX+X,D+-/(DI[I],1)×1E⁻12[(P+SOL)[I])/L1 [14] →(P[12]<ITER+ITER+1)/L2</pre> [15] P[N]+(LIMS[1;15]LP[N]+P[14]×YBX)[LIMS[2;15]

```
「16]
       +L0,,<u>SOLS</u>+<u>SOLS</u>,[1]P
[17]
       L2: SOLUTION HISTORY: '
F187
       +SOLS+(0, 1P[12]),(($SSE)111+pSSE),SSE,SOLS
[19]
       'STAN DARD ERRORS: ', FER+(1 1 QCOV+(((,Y-X+.×Y图X)+.*2)+
      -/\rho X) × H IN V(Q X) + . × DI + X) * . 5
[20]
      <u>PDER</u>+X
      Π
     V P2 DC2 P; DI; CVEC; ITER; N; I; X; V; D; R; Y; AD; S
 [1]
      ADIFFERENTIAL CORRECTIONS PROCEDURE FOR LUNAR OCCULTA
      TION S
 [2]
      AN INE PARAMETER MODEL, FIXED LIMB DARKENING AND TEMPE
      RATURE
 [3]
       SOLS2+SOLS+1 14pP,SSE+0pDI+DI2+(5p0),1.02 1.0003,1.00
      1,(201.0003),1
 Γ4]
      CVEC+ 1 0N+5+15+1TER+0
 [5]
      R+R++/R+(#FILT)[;2]×(#FILT)[;1]BBDY P[4],(1+-/02 2p#F
      ILT) ÷5
 Γ61
      R2+R2++/R2+(#FILT)[;2]×(#FILT)[;1]BBDY P2[4],(1+-+++
      20PFILT)+5
 [7]
      S++/2×P[11]LDARKEN GRID P[5]
 F 8 1
       L0:X+X2+(2+P[3])\rho0
 [9]
       SSE+SSE, [+(Y+OBS-COMP+P[8]+(P2[7]×R2 WIDE P2)+P[7]×R
     WIDE P)+.*2
[10]
      +5+<u>SOL</u>2+P2,0p[+5+<u>SOL</u>+P
[11]
       DI[6]+1.02++/(P[6]<1 2 3×0+180×3600000)/.2 .12 .06
       DI2[6]+1.02++/(P2[6]<1 2 3×0+180×3600000)/.2 .12 .06
       L1:P[I]+DI[I]×1E 12[P[I+1+CVEC+1¢CVEC]
[13]
[14]
       D+-COMP-P[8]+(P2[7]×R WIDE P2)+P[7]×R WIDE P
[15]
       +(5>1+pX+X,D+-/(DI[I],1)×1E<sup>-</sup>12[(P+SOL)[I])/L1
[16]
       AA P[N]+(LIMS[1;15][P[N]+P[14]×YBX)[LIMS[2;15]
[17]
       CVEC + 1 \Phi N
[18]
       L1B:P2[I]+DI2[I]×1E 12[P2[I+1+CVEC+1¢CVEC]
[19]
      →(I=8)/L1B
[20]
       D \leftarrow COMP - P[8] + (P2[7] \times R WIDE P2) + P[7] \times R WIDE P
[21]
      →(9>1+pX+X,D+-/(DI2[I],1)×1E<sup>-</sup>12[(P2+SOL2)[I])/L1B AA
[22]
       AA →(4>1+pX2+X2,D+-/(DI2[I],1)×1E<sup>-</sup>12[(P2+SOL2)[I])/L1
     B
[23]
      P[N]+(LIMS[1;15][P[N]+5+AD+P[14]×YBX)[LIMS[2;15] AA
[24]
       P2[5+S]+(LIMS2[1:S][P2[5+S]+-4+AD)[LIMS2[2:S+1 2 4 5]
       A A
ſ25]
      AA P2[5+S]+(LIMS2[1:S][P2[5+S]+P[14]×YBX2)[LIMS2[2:S+
     1 2 4 5 ]
[26]
      +(P[12] < ITER+ITER+1)/L2</pre>
      +L0,,<u>SOLS+SOLS</u>,[1]P,0p<u>SOLS</u>2+<u>SOLS</u>2,[1]P2
[28]
      L2: 'SOLUTION HISTORY:
[29]
      +<u>SOLS</u>+(0, 1P[12]),((<u>$SSE</u>)111+p<u>SSE</u>),<u>SSE</u>,<u>SOLS</u>,<u>SOLS</u>2[;5+S
[30]
      'STAN DARD ERRORS: ', * ER+(1 1 QCOV+(((, Y-X+.×Y图X)+.*2)*
     -/ pX) ×H IN V2 ( &X) + . × DI + X+X, X2) *.5
     PDER+X
[31]
```

V Z+FFT X;T;I;G;N;Q;□IO
```
Γ1]
                   AFAST FOURIER TRANSFORM OF X: FROM RGS
   F 2 ]
                     AX VECTOR OF REALS OR X[:0]=REAL X[:1]=IMAGINARY
   Ē3]
                    □I0+0
   Γ4]
                   +(2=00X)/L0
   [5]
                     x+x.[.5]0
   [6]
                   L0:T+210(Go2)T1I+2*G+[2@1+oX
   [7]
                   X+((N+I+2), 2, 2)o((I, 2)+X+I+, 5)[T;]
   Γ81
                   T+02 10.00(N+T)+2×N
   Γ9]
                   Z+(+/[1]X),[.5]-/[I+1]X
[10]
                   LA: Z+((N+N+2), (2 \times I+I \times 2), 2) \cap Z
[11]
                   X + (0.1.0) + Z
[12]
                   Z + (0, -1, 0) + Z
[13] Z+(Z+O),Z-O+(+/O×X),[1,5]-/(O+(N,I,2))(I,2)+T)×\phiX
[14] →(1<N)/LA
[15] Z+((2×T.1)oZ)[2+0(Go2)T12*G:]
                  Δ
              V Z+FFTI X:T:I:G:N:O:□IO
   [1] AFAST INVERSE FOURIER TRANSFORM OF X: FROM RGS
   [2]
                 AX IS VECTOR OF REALS. OR X[:0]=REAL X[:1]=IMAGINARY
   [3]
                 \Pi T 0 + 0
   [4]
                  →(2=00X)/L0
   [5]
                  X+X.[0.5] 0
   [6]
                 L0:T+210(Go2)T1I+2*G+[2@1+oX
   [7]
                  X + ((N + I + 2) \cdot 2 \cdot 2) \circ ((I \cdot 2) + X + I + \cdot 5) [T : ]
   [8]
                   T+02 10.00(N+T)+ 2×N
   [9]
                  Z + (+/[1]X) \cdot [.5] - /[I+1]X
[10]
                  LA: Z+((N+N+2), (2\times I+I\times 2), 2) \circ Z
[11]
                   x \neq (0.1.0) \neq Z
[12]
                   Z + (0 - I 0) + Z
[13]
                   Z + (Z+0), Z - 0 + (+ / 0 \times X), [1,5] - / (0 + (N, I,2)) + T) \times \phi X
F141
                   \rightarrow (1 < N) / T.A
[15]
                Z + ((2 \times I_{1}) \rho Z) [2 \perp \Theta (G \rho 2) \top (2 \times G :]
                 Δ
              V W←FNOS P:T
   [1]
                W \leftarrow (P[10] \times P[9] - 1P[3]) \circ + (P[2] \times 1000 \times 30P[6] + 2) + 2 \times T + - 5 -
                T = 12 \times T + P[5] \times (\pm .000000001 500 \times . \times 2 \pm P) \times .5
                 Δ
              ▼ F+N FRESNEL W:□IO
   [1]
                AW IS VECTOR OF FRESNEL NUMBERS
                    AN IS NUMBER OF TERMS IN NUMERICAL INTEGRATION
                 F + .5 \times + / (.5 + (1 1 \circ . \times W) \times (2 + N) \times + / 1 2 \circ .00 \cdot 5 \times ((W + N) \circ . \times 1 + 2 \times 1 + 
                 1.5×N)*2)*2
                 Δ
             V F FTSMOOTH X:□IO:T
   [1]
                    AFOURIER SMOOTHING OF OBSERVATIONAL DATA
   [2]
                    AF IS THE CUTOFF FREQUENCY IN HERTZ
  [3]
                     AX[0] IS THE BIN NUMBER OF OCCULTATION
  [4]
                    AX[1] IS LENGTH OF SMOOTHED DATA SET
   [5]
                    ACH1 MUST BE RESIDENT IN THE WS
  F 6 1
                   □I0+0
```

```
OBS+FFT 1024+(-512+514[x[0]|3595)+CH1)*2
    OBS+FFTI 1024 2+(F,2)+OBS
۲8 T
   OBS+x[1]+(512-L.5×x[1])+OBS
[9]
   Π
   ▼ Y+GRID N:R:R2:X
[1] ANORMALIZED RADIAL GRID AREA DISTRIBUTION (2D) SPHERI
   CAL STAR
     AN = SOUARE ROOT OF NUMBER OF GRID POINTS PER QUADRANT
F 2 ]
[3] X+(1R).[.5] 1+1R+1+pY+(2pN)p0
[4] L1:R2+(X×((R*2)-X*2)*.5)+R×R×10X+R
    Y[;1+R]+Np-/R2-((X×((R*2)-X*2)*.5)+R×R×-10(X+(0 -1+pX
E 5 1
    ) + X) + R+R-1),0
[6]
   →1.1×11≠1+0 R2
[7] Y+Y+4×+/+/Y+(1-1N)ΦΦY×(1N)°.≤1N
    Ω
   ▼ Y+HINV X:C
[1] APATCH FOR PROBLEM WITH MONADIC E IN HARRIS APL
    C+10*.5×10⊕X[1:1]
F2]
    x[:1]+x[:1]+C
F31
F 4 1
    x[1:]+x[1:]*C
Ē 5 ]
    Y+⊞X
[6]
    Y[1:]+Y[1:]+C
[7] Y[:1]+Y[:1]+C
    Π
   V Y+HINV2 X:C
     APATCH FOR PROBLEM WITH MONADIC E IN HARRIS APL
[1]
F27
    C+10*.5×10⊛[/[/X
[3]
    X[:1 6]+X[:1 6]+C
٢4٦
    x[1 6;]+x[1 6;]+C
Ē 5 Ī
     Y≁₽X
٢61
    Y[1 6;]+Y[1 6;]+C
[7] Y[;1 6]+Y[;1 6]+C
    77
   V IDENT
     1 1
[2]
      *** LUNAR OCCULTATION DATA REDUCTION WORKSPACE ****
           VERSION CO3/HARRIS BATCH'
F37
     1
           REVISION DATE: 29 APRIL 1985'
[4]
     .
Ē 5 1
     .
           NOTE: THIS WORKSPACE IS □IO+1'
[6]
[7]
     .
      A GLENN SCHNEIDER DEPARTMENT OF ASTRONOMY
      A UNIVERSITY OF FLORIDA GAINESVILLE, FL 32611
[8]
    Δ
  ▼ INPUT2 X;PR;C
ſ1]
    AX IS RAW OBSERVING DATA (I.E. CH1)
     ACONVERSATIONAL ENTRY OF INITIAL DC2 PARAMETERS
ſ2]
[3]
     NAME+32+1.001+'NAME OR CATALOG NUMBER OF STAR: '
Γ4]
     DATE+27+0,000+'U. T. DATE OF OBSERVATION: '
    COMN+28+0,000+'ANY COMMENTS TO BE PRINTED: '
٢51
```

F 6]	PV←□.0o□+'ENTER THE MONOCHROMATIC WAVELENGTH, OR ENTE
203	R O TO USE FILTER TABLE'
[7]	<pre></pre>
[8]	→(<u>PV</u> ≠0)/10
[9]	'ENTER THE NAME OF THE FILTER TABLE TO BE USED: '
[10]	FILT+C
[11]	IO: <u>PV+PV</u> , U, OpU+'ENTER THE DISTANCE TO THE LUNAR LIMB
	IN KILOMETERS'
[12]	RIN+O,ODO+, ENTER STARIING BIN NOWPER OF DAIR TO BE 03
C 1 2 7	DU-DU DODLIENTED NUMBER OF DATA POINTS TO BE USED!
	PY + PY = 0, $OBC + CRIER ROMBER OF DATA FOR TO BE CONSOBC + (-1 + DV) + (BIN-1) + Y$
	PV+PV. D.OOT+'ENTER THE EFFECTIVE COLOR TEMPERATURE OF
	THE STAR'
F161	PV+PV
	RID PTS./QUADRANT'
[17]	PV+PV, 0+C+206264806.2, 0p +'ENTER ESTIMATED DIAMETER I
	N ARC-MILLISECONDS'
[18]	LIMS+[]+C,Op[]+PR+' ENTER UPPER AND LOWER LIMIT'
[19]	<u>PV+PV</u> , 0,00 +'ENTER INTENSITY (COUNTS) FOR FIRST STAR'
[20]	LIMS+LIMS, 0,00 +PR
[21]	PV+PV, 0,000+'ENTER BACKGROUND SKY-LIGHT CONTRIBUTION'
[22]	LIMS+LIMS, 0,00 +PR
[23]	<u>PV+PV,U-BIN,OPU+'ENTER ESTIMATED BIN NOMBER OF GEOMET</u>
F 0 / 7	RICAL OCCULTATION FOR STAR I'
[24]	LIMS+LIMS, (ITPV) I IND, OPD+ FLOS OK MIROS HOW M
[25]	$PV + PV$, $(PV[2]x^{3}OO + 206265)$, $0 \circ O + !FN TER SHADOW VELOCITY$
	IN A-SEC/SEC'
[26]	LIMS+LIMS, (PV[2]×30[+206265),00]+PR
[27]	PV+PV, 0,000+'ENTER ASSUMED LIMB DARKENING CO-EFFICIEN
	T *
[28]	LIMS+06 2pLIMS,0,0p0+PR
[29]	<u>PV</u> +1E ⁻ 12 [[] <u>PV</u> ,□,0ρ□+'ENTER MAXIMUM NUMBER OF ITERATIONS
	FOR DC PROCEEDURE'
	$\underline{PV} + \underline{PV}$, $ 1 + -f2 2 + \underline{PFILT}$
[31]	PV+PV,U,UDU+'ENTER FRACTIONAL ADJUSTMENT TO BE USED (
[32]	
	PV2[4]+D 0-D+IENTER TEMPERATURE OF THE SECOND STAR!
[34]	$PV2[6]+\Pi \pm C.00\Pi \pm !ENTER DIAMETER ESTIMATE FOR SECOND ST$
20.3	AR'
[35]	LIMS2[;1]+[+C,0p[+PR
[36]	PV2[7]+0,000+'ENTER INTENSITY (COUNTS) FOR SECOND STA
	R'
[37]	LIMS2[;2]+[,0p[+PR
[38]	<u>PV2[9]+U-BIN</u> ,0pU+'ENTER TIME OF GEOMETRICAL OCCULTATI
C 2 0 7	UN (BIN) FOR STAR Z'
L 2 9]	LINDZL;4JT(IT PV2/ *. TI IAU, UPU+ PLUS OR MINUS HOW MA
	V AIDDIGIGON DOI
	,
	V INPUT X; PR; C
[1]	AX IS RAW OBSERVING DATA (I.E. CH1)

F 2 7 ACONVERSATIONAL ENTRY OF INITIAL DC PARAMETERS [3] NAME+32+ . Oo + 'NAME OR CATALOG NUMBER OF STAR: ' Γ47 DATE+27+0.000+'U. T. DATE OF OBSERVATION: ' Γ51 COMN+28+F.OoF+'ANY COMMENTS TO BE PRINTED: ' [6] PV+0.00 +'ENTER THE MONOCHROMATIC WAVELENGTH, OR ENTE R O TO USE FILTER TABLE' •(PV≠0)/'FILT+''MONO'''.0pMONO+1 2pPV.1 [7] Ē8] +(PV≠0)/I0 [9] 'ENTER THE NAME OF THE FILTER TABLE TO BE USED: " [10] FILT+ [11] 10:PV+PV.□.00□+'ENTER THE DISTANCE TO THE LUNAR LIMB IN KILOMETERS' [12] BIN+□.00□+'ENTER STARTING BIN NUMBER OF DATA TO BE US ED [13] PV+PV.□.00□+'ENTER NUMBER OF DATA POINTS TO BE USED' [14] OBS + (-1 + PV) + (BIN - 1) + X[15] PV+PV,□,00□+'ENTER THE EFFECTIVE COLOR TEMPERATURE OF THE STAR' [16] PV+PV.□.00□+'ENTER THE SOUARE ROOT OF THE NUMBER OF G RID PTS./OUADRANT' [17] PV+PV, 0+C+206264806.2.000+'ENTER ESTIMATED DIAMETER I N ARC-MILLISECONDS' [18] LIMS+[]+C.Op[+PR+' ENTER UPPER AND LOWER LIMIT' [19] PV+PV, 0,000+'ENTER PRE-EVENT LEVEL OF STAR+SKY SIGNAL [20] LIMS+LIMS. 0.00T+PR [21] PV+PV,□,00□+'ENTER POST-EVENT LEVEL OF LIMB+SKY SIGNA 1.1 [22] LIMS+LIMS, 0,00 +PR [23] PV+PV, □-BIN,00 □+'ENTER ESTIMATED BIN NUMBER OF GEOMET RICAL OCCULTATION ' [24] LIMS+LIMS, $(-1+PV) \circ + 1 - 1 \times 1 \circ 0 \circ 1 \leftrightarrow ?$ PLUS OR MINUS HOW M ANY MILLISECONDS? * PV+PV, (PV[2]×30[+206265),00[+'ENTER SHADOW VELOCITY IN A-SEC/SEC' [26] LIMS+LIMS, (PV[2]× 30[+206265), 0p[+PR [27] PY+PY, 0,000+'ENTER ASSUMED LIMB DARKENING CO-EFFICIEN тţ [28] LIMS+Q6 2pLIMS, 0,0p0+PR [29] PV+1E⁻12 [PV,□,0p□+'ENTER MAXIMUM NUMBER OF ITERATIONS FOR DC PROCEEDURE' [30] PV+PV. 11+-/2 2+ FILT [31] PY+PY, 0,000+'ENTER FRACTIONAL ADJUSTMENT TO BE USED (0 TO 1)' ▼ Y+X LDARKEN G:T [1] AY=NORMALIZED GRID DISTRIBUTION OF LIMB DARKENED STA R AG=STELLAR GRID QUADRANT TO BE LIMB DARKENED (GENERA TED BY: GRID) [3] AX=LIMB DARKEN ING COEFFICIENT

[4] ARESULT OBTAINED FROM A LINEAR LIMB DARKENING LAW

 $[5] \quad Y+Y+4x+/+/Y+Gx(T\circ,\leq T)\times(1-T)\phi\phi(\rho G)\rho(1-X)+X\times20^{-1}O(T-.5)$

```
+0T+11+0G
     Π
    ▼ Y+NPOL V:T:I:□CT
     AINTERPOLATION OF FRESNEL INTENSITY VALUE FOR FRESNE
 [1]
     L NUMBER <V>
     Y + FREN[1] + (- \neq FREN[^1] 0 \circ + I + (1 + .5 \times \rho FREN) - 100 \times T]) \times 100 \times V
 F 2 ]
     -T+.01×L.5+V×100
     Δ
    V OUTPUT2:S:T:PM:CT:W:O
 [1]
     AFORMATTED OUTPUT FOR DC2 SOLUTION
 [2]
      (200' '), 'OCCULTATION OF: ',NAME
      (20p' '), DATE, ' U. T.'
 F 3 T
 Γ41
      (200 ' ').COMN
      1 1
 [5]
 [6]
     IN PUT PARAMETERS:
      1 1
 [7]
      'CENTRAL WAVELENGTH OF PASSBAND: ', (Ov+/×/@FILT).' ANG
 [8]
     STROMS *
 [9]
      'LUNAR LIMB DISTANCE (KM):
                                          ',6 0 TSOL[2]
[10]
     'EFFECTIVE TEMPERATURE STAR-1 : ', VSOL[4]
     *EFFECTIVE TEMPERATURE STAR-2 : *. *SOL2[4]
[11]
     'LIMB DARKENING COEFFICIENT STAR-1: '. FOL[11]
[12]
[13]
      'LIMB DARKENING COEFFICIENT STAR-2: ', # SOL2[11]
      1 1
[14]
[15]
      MODEL PARAMETERS:
      1 1
[16]
[17]
      'NUMBER OF DATA POINTS: ', VSOL[3]
      'NUMBER OF GRID POINTS: ', T(SOL[5]×2)*2
F187
      'NUMBER OF SPECTRAL REGIONS: ', #1+p#FILT
[19]
F207
      WIDTH OF SPECTRAL REGIONS: '.(V-/$(*FILT)[1,2]]to*F
     ILT;1]),' ANGSTROMS'
[21]
      'NUMBER OF ITERATIONS: '. FSOL[12]
[22]
     1 1
[23]
     * SOLUTIONS FOR STAR-1:*
     1 1
F247
[25]
      PM+ ' +/- '
[26]
      'STELLAR DIAMETER (ARC-MILLISECONDS): '.(6 27SOL[6]×2
     06264806), PM.6 2 ¥ ER[1]×206264806
     'BIN OF GEOMETRICAL OCCULTATION:
[27]
                                               '.(7 1¥SOL[9]+B
     IN), PM, 4 1 # ER[4]
[28]
      'SIGNAL LEVEL OF STAR-1 (COUNTS)
                                               '.(6 1¥SOL[7]).
     PM.5 1 #ER[2]
[29]
     ST+×1-T+SOL[10] + PV[10]
[30]
     'OBSERVED LUNAR SHADOW VELOCITY: ',(7 5*SOL[10]),
     PM.(7 5 # ER[5]). * KM/SEC*
[31]
     'PREDICTED LUNAR SHADOW VELOCITY:
                                             *.(7 5*PV[10]).*
      KM/SEC'
[32]
     T+'LOCAL SLOPE OF LUNAR LIMB:
                                              '.'- +'[ST+2]..
     (6 27.5×+/S+(180+01)× 20T+T[+T), DEG'
     T, PM, 5 3 # 1.5 × - / 0 [ 0 + ( + 0 + 180 ) × (SOL [ 10 ] • . + 1 ] × ER[ 5 ] ) + P
[33]
     <u>v[10]</u>
[34] 2 1p' '
```

[35] * SOLUTIONS FOR STAR-2:* . . [36] F 37 T 'STELLAR DIAMETER (ARC-MILLISECONDS): ',(6 27SOL2[6]× 206264806) PM.6 2FER[6]×206264806 [38] 'BIN OF GEOMETRICAL OCCULTATION: ',(7 1 ¥SOL2[9]+ BIN), PM, 4 17ER[8] [39] 'SIGNAL LEVEL OF STAR-2 (COUNTS) ', (6 1 v SOL2[7]) ,PM,4 17<u>ER[7]</u> POST-EVENT LIMB+SKY SIGNAL LEVEL: '.(6 1 #SOL[8]). F40] PM.4 1 VER[3] F417 'OBSERVED LUNAR SHADOW VELOCITY: '.(7 5▼SOL2[10]) .PM.(7 5 # ER[9]). * KM/SEC* F42] ST+×1-T+SOL2[10]+PV[10] [43] T+ LOCAL SLOPE OF LUNAR LIMB: ','- +'[ST+2]., (6 2 ¥.5×+/S+(180 ±01)× 20T+T[±T). DEG * Γ44] T.PM.6 $3 \forall [.5 \times -/0] | 0 \leftrightarrow (\div 0 \div 180) \times (SOL2[10] \circ . +1 ~ 1 \times ER[9]) \div$ PV[10] [45] 2 10' ' [46] . . OTHER PARAMETERS AND DERIVED SOLUTIONS: * 1 1 [47] F487 'POST-EVENT LIMB+SKY SIGNAL LEVEL: '.(6 1¥<u>SOL[8]</u>), PM.4 1 # ER[3] [49] 'TEMPORAL SEPARATION OF THE STARS: '.(6 1+|SOL[9]-S OL2[9]), PM. (4 1 F(ER[4 8]+.*2)*.5), * MS.* 'WEIGHTED MEAN L-RATE: ', (6 4 W+((SOL[10], SOL2[10])+. [50] *ER[5 9]) ++/ + ER[5 9]) . * KM/SEC* [51] 'SPATIAL SEPARATION FROM R-RATES WHICH ARE: ' [52] T+' PREDICTED ('.(*Q).' ARC SEC/SEC): '.6 2*(|SOL2[9]-SOL[9])×0+206265×30+/PV[10 2] [53] T. MILLISECONDS OF ARC' T+' WEIGHTED MEAN (',(6 4¥Q),' ARC SEC/SEC): ',6 2₹([54] |SOL2[9]-SOL[9])×0+206265×30W ± PV[2] [55] T. ' MILLISECONDS OF ARC' [56] 'BRIGHTNESS RATIO (BRIGHTER/FAINTER): ',6 4T++/(SOL[7],SOL2[7])[#SOL[7],SOL2[7]] [57] 'MAGNITUDE DIFFERENCE: *.6 4▼(100*.2)⊕T [58] 2 10' ' [59] . PHOTOMETRIC STATISTICS: * 1 1 [60] F617 'SUM OF SQUARES OF RESIDUALS: ', VS+(OBS-COMP)+.*2 [62] 'SIGMA (STANDARD ERROR): '.▼S+(S+ 1+pOBS)*.5 'NORMALIZED STANDARD ERROR: [63] ', ▼S+S+SOL[7]+SOL2[7] [64] 'PHOTOMETRIC (SIGNAL+NOISE)/NOISE RATIO: ', *(1+S) *S '(INTENSITY CHANGE)/BACKGROUND: ', *S+(SOL[7]+SOL2[[65] 7]) + SOL[8] [66] 'CHANGE IN MAGNITUDE: *, ¥-(+100++5)⊕1+S 2 10' ' [67] [68] 1 SUPPLEMENTAL STATISTICS: * 1 1 [69] [70] 'VARIANCE/CO-VARIANCE MATRIX:' [71] CT+9 11ρ1+□+' DIAMETER1 INTENSITY1 SKY INTS OCC TI ME1 L-RATE1 DIAMETER2 INTENSITY2 OCC TIME2 L-RAT 'MD-DE11.4' DFMT COV [72] Ē73Ī 1 1

```
[74] (13p' '), CT
[75] ((110/\CT=' ') OCT), 'ME-EF11.6' [FMT CM COV
[76]
    RESTAR
     Δ
    V OUTPUT:S:T:PM:CT
 [1] APRODUCES A FORMATTED REPORT ON THE STELLAR SOLUTION
 [2]
     (20p' '), 'OCCULTATION OF: ', NAME
 [3] (20p' '), DATE, ' U. T.'
 [4]
     (20p' '), COMN
 [5]
      1 1
 [6]
      ' INPUT PARAMETERS: '
 [7]
     . .
 [8]
     'CENTRAL WAVELENGTH OF PASSBAND:', (0 +/×/ FILT),' ANG
     STROMS!
 [9]
     'LUNAR LIMB DISTANCE (KM):
                                         ',6 0 # SOL[2]
     'EFFECTIVE TEMPERATURE (KELVIN): ', VSOL[4]
[10]
                                        '. # SOL[11]
[11]
      'LIMB DARKENING COEFFICIENT:
[12]
[13]
     'MODEL PARAMETERS:'
     1 1
[14]
[15]
      'NUMBER OF DATA POINTS: '. #SOL[3]
[16]
     'NUMBER OF GRID POINTS: '. F(SOL[5]×2)*2
     'NUMBER OF SPECTRAL REGIONS: ', #1+p#FILT
[17]
[18]
     WIDTH OF SPECTRAL REGIONS: '. (*-/$(*FILT)[1,2]]+0*F
     ILT;1]),' ANGSTROMS'
[19]
      'NUMBER OF ITERATIONS: ', * SOL[12]
[20]
     1 1
[21]
     ' SOLUTIONS: *
[22]
     1 1
[23]
     PM+' +/-'
[24]
      'STELLAR DIAMETER (ARC-MILLISECONDS): '.(6 2 SOL[6]×2
     06264806), PM, 6 2 FER[1]×206264806
[25]
     'BIN OF GEOMETRICAL OCCULTATION:
                                             '.(7 1 #SOL[9]+B
     IN), PM, 4 1 FER[4]
[26]
      'PRE-EVENT STAR+SKY SIGNAL LEVEL:
                                             '.(6 1 #SOL[7]).
     PM,4 17ER[2]
[27]
     'POST-EVENT LIMB+SKY SIGNAL LEVEL:
                                             ',(6 1*<u>SOL[8]</u>),
     PM,4 1 # ER[3]
[28]
     ST+×1-T+SOL[10] + PV[10]
[29]
     'OBSERVED LUNAR SHADOW VELOCITY:
                                            ',(7 5¥<u>SOL</u>[10]),
     PM, (7 5 FER[5]), ' KM/SEC'
[30]
      'PREDICTED LUNAR SHADOW VELOCITY: '.(7 5*PV[10]).'
      KM/SEC'
[31] T+'LOCAL SLOPE OF LUNAR LIMB:
                                              ','- +'[ST+2]..
     (6 2 .5 ×+/S+(180 +01) × 20T+T[+T), ' DEG'
[32]
     T, PM, 5 37 |.5×-/QL |Q+(+0+180)×(SOL[10] .+1 1×ER[5])+P
     ¥[10]
[33]
      2 10' '
[34]
      'VARIANCE-COVARIANCE MATRIX:'
[35]
     'MO-OF12.5' OFMT COV
[36]
     2 1 p ' '
F371
     'CORRELATION MATRIX:'
[38]
      60+,CT+5 10p' DIAMETER
                                  *+LIMB LIMB+SKY TIME OC
```

C VELOCITY ' [39] ((1⊥\oh)(CT=' '))¢CT),'MŪ-ŪF10.5' □FMT CM <u>COV</u> [40] 2 lp' ' 'SUM OF SQUARES OF RESIDUALS: ', FS+(OBS-COMP)+.*2 [41] 'SIGMA (STANDARD ERROR): ', #S+(S+ 1+pOBS)*.5 [42] [43] 'NORMALIZED STANDARD ERROR: ', FS+S+-/SOL[7 8] 'PHOTOMETRIC (SIGNAL+NOISE)/NOISE RATIO: '. #(1+S)#S F447 '(INTENSITY CHANGE)/BACKGROUND: ', #S+(-/SOL[7 8])* F457 SOL[8] '.▼-(+100*+5)@1+S [46] 'CHANGE IN MAGNITUDE: [47] RESTAB Δ V RESTAB:0:N APRODUCES A TABLE OF OBS, COMP, AND RESIDUALS F11 4 10' ' [2] [3] (29+1 1), 'TABLE 5-XXX' [4] ' ',NAME,': OBSERVATIONS, COMPUTED VALUES, AND RESID HALS FROM BIN '. FBIN [5] '' [6] 720*=* [7] 'NUM OBS COMP RESID NUM OBS COMP RESID NU M OBS COMP RESID' [8] '--- ------- ---- ----- ------ ---- ----- ------! [9] Q+((3×N+[(pOBS)+3),4)+(-1+1pOBS),OBS,COMP,[1.5]OBS-CO MP $[10] Q+Q[1N;],((N,0)+(-N,0)+Q),(2\times N,0)+Q$ '3(BI3.X1.BI4,X1,BF6.1,X1,M2-EBF6.1,X3)' [FMT Q [11] [12] 72p'=' V ▼ START [1] AOVERALL CONTROL FUNCTION FOR DC EXECUTION [2] 'REDUCTION RUN FOR THE STAR: '.NAME [3] DC PV [4] DC PV+(-3+3+, SOLS[SOLH+SOLS[:2]11:]),10,PV[13].5 [5] DC(-3+3+.SOLS[SOLF+SOLS[:2]11:]).0.PV[13].5 [6] OUTPUT V V START2 [1] **AOVERALL CONTROL FUNCTION FOR DC2 EXECUTION** 'REDUCTION RUN FOR THE STAR ',NAME [2] [3] PV2 DC2 PV V ▼ Y+R WIDE P:I:T [1] ACOMPUTATION OF WIDEBAND FRESNEL PATTERN [2] AP IS THE CURRENT PARAMETER VECTOR (PV) [3] AR=CONVOLVED SYSTEM RESPONCE AND SOURCE SPECTRAL ENER GY CURVES [4] I+1+Y+P[3]00 [5] AS++/2×P[11]LDARKEN GRID P[5]

[6] W0:Y+Y+R[I]×(NPOL FNOS P+(<u>sFILT</u>)[I+I+1;1],1+P)+.×(\$T) ,T+\$ +(IF1+psFILT)/W0 v

NAME: VFILTER SHAPE: 53 2	NAME: BFILTER SHAPE: 42 2	NAME: NARROWV SHAPE: 5 2
4450 0.0001294 4500 0.0002588	3600 0.0000953 3650 0.0004767	5295 0.0621118 5345 0.2251553 5205 0.4687578
4600 0.0005175 4650 0.0006469	3750 0.0023834 3800 0.0044809	5445 0.2251553 5495 0.0388199
4700 0.0007763 4750 0.0011644 4800 0.0016820	3850 0.0115359 3900 0.0207837 3950 0.0397559	NAME: NARROWB
4850 0.0023289 4900 0.0047872	4000 0.0469063 4050 0.0489084	SHAPE: 5 2
4950 0.0112563 5000 0.0208306	4100 0.0499571 4150 0.0505291 4200 0.0511965	4600 0.0658106 4650 0.2247191 4700 0.4558587
5100 0.0380386 5150 0.0450252	4250 0.0513872 4300 0.0514825	4750 0.2247191 4800 0.0288925
5200 0.0513650 5250 0.0548583	4350 0.0510058 4400 0.0497664	
5300 0.0555052 5350 0.0553759 5400 0.0545996	4500 0.0475850 4500 0.0451902 4550 0.0427114	
5450 0.0534351 5500 0.0507181	4600 0.0400419 4650 0.0376585	
5600 0.0452840 5650 0.0426963	4750 0.0331776 4800 0.0295548	
5700 0.0398499 5750 0.0364860 5800 0.0337689	4850 0.0271713 4900 0.0244065 4950 0.0219277	
5850 0.0307931 5900 0.0282055	5000 0.0192583 5050 0.0164935	
5950 0.0252297 6000 0.0225126 6050 0.0196076	5100 0.0137287 5150 0.0111545 5200 0.0085804	
6100 0.0166904 6150 0.0142321	5250 0.0058156 5300 0.0034322	
6200 0.0122914 6250 0.0104800 6300 0.0087980	5350 0.0020974 5400 0.0015254 5450 0.0008580	
6350 0.0071161 6400 0.0054341	5500 0.0004767 5550 0.0002860	
6450 0.0041403 6500 0.0029758 6550 0.0024583	5600 0.0001907 5650 0.0000953	
6600 0.0020701 6650 0.0016820		
6750 0.0012938 6800 0.0010351		
6850 0.0007763 6900 0.0006469		
7000 0.0002588		

7050 0.0001294

APPENDIX E LISTING OF THE APL WORKSPACE OCCPLOTS

```
V Y+N BIN D;□IO;T
 [1]
      ABIN DATA FOR INTO CLASSES FOR HISTOGRAM
 [2]
      AD IS DATA VECTOR: N IS NUMBER OF CLASSES
 [3] [10+1
       Y+(T+D[1]-.5×1+T),[1.5]+/(<sup>-</sup>1+1N) •.=+/((1+D)+T++\Np(([
 F41
      /D)-|/D) +N) ∘. <D+D[▲D]
      Δ
    V D CENTER S:R
 [1] ACENTER TEXT LABEL
 [2] R+S[0]+.5×S[1]-(pD)×8*53
 [3] ((1.oD)oD)LABELD 1 3oR, 2+S
      Π
    V Y+N ESMOOTH D:□IO:T:V
 [1] AN-POINT WEIGHTED SMOOTHING OF THE VECTOR D
 [2]
     A(1+N) IS NUMBER OF TERMS TO USE NUMERICALLY IN E SMO
      OTHING
 [3]
     e(1+N) IS THE 1/E FOLDING SPACING WRT THE DATA SPACIN
     G
 Γ41
      \Pi T 0 + 0
 [5]
      Y \leftarrow ((V \leftarrow 1.5 + .5 \times 1 + N) p + Y), Y \leftarrow (T + + / T \leftarrow * - (+ 1 + N) \times (1 + N) + . \times (1 + 1 + N)
     N) Φ(1+N, pD) pD
      Y+(-3×V)+Y
 F 6 1
 [7] Y \leftarrow (\rho D) \rho Y, (2 \times V) \rho (+/(-1 + N) + D) + 1 + N
      V
    ▼ Z+FFT X:T:I:G:N:Q:□IO
 [1]
      AFAST FOURIER TRANSFORM OF X; FROM RGS
       AX VECTOR OF REALS OR X[:0]=REAL X[:1]=IMAGINARY
 [2]
       □I0+0
 Ē4Ī
       \rightarrow (2=00X)/LO
       X+X,[.5]0
 [6]
       L0:T+210(Gp2)T1I+2*G+[2@1+pX
      X+((N+I+2),2 2)p((I,2)+X+I+.5)[T;]
 [8]
      T+02 10.00(N+T)+2×N
 [9]
      Z+(+/[1]X),[.5]-/[I+1]X
[10]
      LA: Z+((N+N+2),(2×I+I×2),2)oZ
[11]
      X + (0, I, 0) + Z
```

```
[12]
      Z+(0,-I,0)+Z
[13]
      Z+(Z+Q),Z-Q+(+/Q×X),[1.5]-/(Q+(N,I,2)p(I,2)+T)×ΦX
[14]
      \rightarrow (1 < N)/LA
[15]
     Z+((2×I,1)pZ)[2⊥⊖(Gp2)T12*G;]
     Δ
    V Z+FFTI X;T;I;G;N;Q;□IO
       AFAST INVERSE FOURIER TRANSFORM OF X: FROM RGS
 [1]
       AX IS VECTOR OF REALS, OR X[;0]=REAL X[;1]=IMAGINARY
 [2]
 [3]
      \Pi I 0 + 0
 Γ41
      \rightarrow (2=00X)/L0
      x+x,[0.5] 0
 F 5 1
 [6]
      L0:T+210(Gp2)T1I+2*G+[2@1+pX
      X+((N+I*2), 2, 2)\rho((I,2)+X*I*.5)[T;]
 [8]
      T + b_2 = 1 \circ .00 (N + T) = 2 \times N
 [9]
      z+(+/[1]x),[.5]-/[1+1]x
      LA: Z+((N+N+2), (2\times I+I\times 2), 2) \rho Z
[10]
[11]
      X+(0,I,0)+Z
      Z + (0, -I, 0) + Z
[12]
[13]
      Z+(Z+O), Z-O+(+/O\times X), [1,5]-/(O+(N,I,2))(I,2)+T)\times \phi X
[14]
      +(1 <N)/LA
      Z + ((2 \times I, 1) \rho Z) [2 \perp \Theta (G \rho 2) T 1 2 * G;]
[15]
     V FITPLOT X: IO; ZRO; SPA; XPOS; STV; NUMN; N
 [1]
      APLOT OF THE FITTED OCCULTATION CURVE
 [2]
       AX=STARTING BIN NUMBER: OBS, COMP, SOL MUST BE GLOBAL
 [3]
       □I0+0
 [4]
      x+x,(x+sol[2]),(x+sol[8]),sol[9],sol[5]\times 206264810
 [5]
      START
      PLOTSET -.3 0
 [6]
      AXISD 2 1.75 7 0 .35 .15
AXISD 2 1.75 7 0 1.4 .05
 F87
      AXISD 2 1.75 4.75 90 .475 .15
 Γ91
       AXISD 2,(1.75+4.75+20),4.275 90,(4.75+10), .08
F107
[11]
       AXISD 9 1.75 4.75 90 .475 .15
       AXISD 9,(1.75+4.75+20).4.275 90 .475 .08
[12]
[13]
       ZRO + 2 + (-/X[2 0]) \times 7 + -/X[1 0]
[14]
       SPA+10×*X[3]×(-/X[1 0])*7
[15]
       AXISD ZR0.6.5.(ZR0-2).180.SPA..15
[16]
       AXISD ZR0,6.5,(9-ZR0),0,SPA, .15
[17]
      ('ZF3.1' [FMT .1×111)LABELD 1.55,(1.75+.475×111),[.5]
F187
      (*LI4* □FMT(<u>|/OBS</u>)+((<u>[/OBS</u>)-<u>|/OBS</u>)×.1×111)LABELD 9.05
      ,(1.75+.475×111),[.5]0
[19]
       ('LI4' [FMT X[0]+(15)×.2×1+-/X[1 0])LABELD(2+1.4×15)
      ,1.55,[.5]500
[20]
      XPOS+ZRO-SPA×L(ZRO-2) * SPA
[21]
      STV+-10×| (ZRO-2) + SPA
[22]
      NUMM+([(ZRO-2) + SPA)+[(9-ZRO) + SPA
[23]
       ('MM-MLI4' DFMT STV+10×1NUMM)LABELD((XPOS+NUMM++\0,NU
      MMpSPA).[.5]6.55).0
       'TIME IN MILLISECONDS FROM BEGINNING OF DATA WINDOW'
[24]
      CENTER 1 9 1.3 0
```

```
'ANGULAR SEPARATION IN MILLISECONDS OF ARC' CENTER 1
[25]
     9 6.75 0
      (1 20p'NORMALIZED INTENSITY')LABELD 1 3p1.4 2.8 90
[26]
      (1 6p'COUNTS')LABELD 1 3p9.8 4.9 270
[27]
     N \neq 0
[28]
      L1:AXISD ZRO,(2+.475×N),(.475+2),270
[29]
[30]
     +(10≠N+N+1)/L1
      1 SPLINE(2+7×(1pOBS)*1+pOBS),1.75+4.75×(OBS-L/OBS)*[
[31]
     /OBS-L/OBS
     1 SPLINE(2+7×(10 COMP) * 1+0 COMP), 1.75+4.75×(COMP-L/OBS)
[32]
     ) + [/OBS-1/OBS
      AXISD 8.5,(1.75+4.75×(SOL[6]-L/OBS)+([/OBS)-L/OBS),.5
[33]
      AXISD 2.(1.75+4.75×(SOL[7]-L/OBS)*([/OBS)-L/OBS),.5 0
[34]
[35]
      AXISD ZR0,6.32,(X[4]×.1×SPA),0
[36]
      PLOTEN D
     Δ
    V FITPLOT2 X: DIO: ZRO: SPA: XPOS: STV: NUMN:N
 [1]
      APLOT OF THE FITTED TWO-STAR OCCULTATION CURVE
 [2]
      AX=STARTING BIN NUMBER: OBS, COMP, SOL MUST BE GLOBAL
 [3]
      □I0+0
      X+X,(X+SOL[2]),(SOL[8]+X2+X),SOL[9],SOL[5]×206264810
 [4]
      START
 [6]
     PLOTSET .3 0
 [7]
      AXISD 2 1.75 7 0 .35 .15
                             -.05
 [8]
      AXISD 2 1.75 7 0 1.4
      AXISD 2 1.75 4.75 90 .475 -.15
 [9]
[10]
     AXISD 2,(1.75+4.75+20),4.275 90,(4.75+10), .08
[11]
      AXISD 9 1.75 4.75 90 .475 .15
      AXISD 9,(1.75+4.75+20),4.275 90 .475 .08
[12]
      ZRO+2+(-/X[2 0])\times7+-/X[1 0]
[13]
      SPA+10×*X[3]×(-/X[1 0])*7
[14]
[15]
      AXISD ZR0.6.5.(ZR0-2).180.SPA..15
[16]
      AXISD ZR0,6.5,(9-ZR0),0,SPA, .15
[17]
      ('ZF3.1' [FMT .1×111)LABELD 1.55,(1.75+.475×111),[.5]
     0
[18]
     ('LI4' □FMT(<u>L</u>/<u>OBS</u>)+((<u>Γ</u>/<u>OBS</u>)-<u>L</u>/<u>OBS</u>)×.1×111)LABELD 9.05
     ,(1.75+.475×111),[.5]0
[19]
     ('LI4' □FMT X[0]+(15)×.2×<sup>-</sup>1+-/X[1 0])LABELD(2+1.4×15)
     .1.55.[.5]500
F207
      XPOS+ZRO-SPA×L(ZRO-2) + SPA
[21]
     STV+-10×L(ZRO-2)*SPA
[22]
     NUMM+(L(ZRO-2)*SPA)+L(9-ZRO)*SPA
      ('ME-ELI4' FMT STV+10×1NUMM)LABELD((XPOS+NUMM++\0,NU
[23]
     MMpSPA),[.5]6.55),0
[24]
      'TIME IN MILLISECONDS FROM BEGINNING OF DATA WINDOW'
     CENTER 1 9 1.3 0
[25]
      'ANGULAR SEPARATION IN MILLISECONDS OF ARC' CENTER 1
     9 6.75 0
[26]
     (1 20p 'NORMALIZED INTENSITY')LABELD 1 3p1.4 2.8 90
      (1 6p'COUNTS')LABELD 1 3p9.8 4.9 270
[28]
      N \neq 0
[29]
      L1:AXISD ZRO, (2+.475×N), (.475+2), 270
```

```
\lceil 30 \rceil \rightarrow (10 \neq N \neq N + 1) / L1
[31] 1 SPLINE(2+7×(100BS)+ 1+00BS), 1.75+4.75×(0BS-L/0BS)+[
     /OBS-L/OBS
     1 SPLINE(2+7×(10 COMP)+ 1+0 COMP), 1.75+4.75×(COMP-L/OBS
F 32 ]
     ) + [/OBS-L/OBS
     AXTSD 8.5.(1.75+4.75×((SOL[7]+SOL2[6]+SOL[6])-L/OBS)+
[33]
     ([/OBS)-1/OBS)...5 0
     AXISD 2,(1.75+4.75×(SOL[7]-L/OBS)+([/OBS)-L/OBS),.5 0
[34]
[35] AXISD ZR0,6.32,(X[4]×.1×SPA).0
[36] x + x2, (x2 + sol2[2]), (x2 + sol2[8]), sol2[9], sol2[5] \times 206264
     810
[37] N+0×1+ZRO+2+(-/X[2 0])×7+-/X[1 0]
[38]
     SPA+10×*X[3]×(-/X[1 0])*7
[39] L2:AXISD ZRO.(2+.475×N).(.475+2).270
\lceil 40 \rceil \rightarrow (10 \neq N + N + 1) / L2
[41] AXISD ZR0,6.32,(X[4]×.1×SPA),0
[42] PLOTEND
     V
    V R←FT Y:M:N:□IO
 [1] ADISCRETE FOURIER TRANSFORMATION OF EVENLY SPACED REA
     T. DATA
 [2]
     \Pi T 0 + 0
 [3] M+(12×N)×O+N+(pY)+2
 [4] R+1 2ρΦ, (+/Y)+2×N
[5] L1:R+R.[0](*N)×+/(1 10.×Y)×2 10.0(1+pR)×M
 [6]
    +(N><sup>-</sup>1+1+pR)/L1
    R[N:]+.5 0×. 1 2+R
     Ω
    V FTPREP X:□IO:T
 [1]
      ASETS UP GLOBAL FOR POWER SPECTRA ANALYSIS
F 2 1
      AX IS CENTER OF TRANSFORM WINDOW FROM DATA SET
 [3]
      □I0+0
[4]
      FTOBS+.5×+/(FFT FTDATA+1024+(~512+514[X|3595)+CH1)*2
[5]
     [6]
[7]
      FTOUT+.5×(FFT FTOUT)+.*2
     T+[1+SOL[6]+(COMP[0]-SOL[6]) \times T+[/T+T-L/T++++1+L.5\times1]
۲8 T
     024-0COMP
 [9]
     FTMOD+T.COMP
[10]
      FTMOD + FTMOD, 1 + SOL[7] + (-SOL[7] - 1 + FTMOD) \times \Phi T + [/T + T - ]/T +
     + \+ \ 11+ 1024- p FTMOD
     FIMOD+.5×(FFT FIMOD)+.*2
[11]
    V IDENT
 [1]
      .
          ***LUNAR OCCULTATION PLOTTING WORKSPACE****
[2]
      1
              VERSION: B01/HARRIS'
[3]
      1
              WSID: DISPLOT'
[4]
      .
              LAST REVISION: 10 APRIL 85'
Ē 5 Ī
      .
[6]
      1
              NOTE: THIS WORKSPACE IS DIO+0'
     Ω
```

```
▼ INTPLOT X:□IO
      AINTEGRATION PLOT OF THE DATA VECTOR X
 F11
 [2]
      \Pi I0 + 0
 Ē31
      START
 Ē41
      AXISD 2 1.75 7 0 .4375 .15
 [5]
      AXISD 2 1.75 7 0 .875 .05
      AXISD 2 1.75 4.75 90 .475 .15
 Γ61
      AXISD 2.(1.75+4.75+20),4.275 90,(4.75+10), .08
      ('ZF3.1' [FMT .1×111)LABELD 1.55,(1.75+.475×111),[.5]
 [8]
     0
     ('LI4' [FMT((pX) +8)×18)LABELD(2+.875×18),1.55,[.5]8p0
 [9]
      'TIME IN MILLISECONDS FROM BEGINNING OF DATA WINDOW'
ſ10]
     CENTER 1 9 1.3 0
     (1 200'NORMALIZED INTENSITY')LABELD 1 301.25 3.22 90
[11]
[12]
      x \leftarrow (1/x) - x \leftarrow x \leftarrow (+/x) \neq \rho x \leftarrow x [2 3] \cdot 2 \neq x
[13]
      1 SPLINE (2+7×(1pX)*1+pX),1.75+4.75×X*[/X
[14]
     PLOTEND
     V
    V F+X INVFT R:□IO
 [1] AINVERSE FOURIER TRANSFORMATION FROM COEFFICIENT MATR
     тΧ
 [2]
      □I0+0
 [3]
     X+2 0 102 10.0X0.×1+11+pR
 [4]
     F \leftarrow (1 \ 1 + R) + + / + / ((\rho X) \rho 1 \ 0 + R) \times X
    V NOISEPLOT X: OC: PD: R: MEAN: SIGMA; Y; BR; N
 [1]
      ANOISE FIGURE FOR LUNAR OCCULTATION
 [2]
       AX IS BIN NUMBER OF START OF DATA USED IN SOLUTION
 [3]
      AMUST HAVE SOL OBS AND COMP AS GLOBAL VARIABLES
 [4]
      OBSA = SOL[2] + (X-1) + CH1
      AX[1]=FIRST BIN OF DATA X[2]=LENGTH OF DATA
 [5]
      OC+((X-1)+CH1)-SOL[6]
 [6]
      OC \leftarrow OC, (OBS - COMP), ((-1 + SOL[2] + X) \neq CH1) - SOL[7]
 [8]
      PD+BR+$50 BIN OC+2+OC+CH1
[9]
      PD[0:]+PD[0:]-PD[0:0]
[10]
      PD + PD \times Q(\phi_0 PD)_07 \quad 4.75 \pm [/PD
[11]
      X+2+1+<sup>-</sup>1+,(7×(151)+50)•.×1 1
[12]
      Y+1.75+.PD[1:] •.×1 1
[13]
      START
      PLOTSET -.1 0
[14]
[15]
      AXISD 2 1.75 7 0 .14_.15
[16]
      AXISD 2 1.75 7 0 .7
                             -.05
[17]
      AXISD 2 1.75 4.75 90 .475 .15
[18]
      ('ZI4' □FMT .1×(111)×0 1/R+([/BR)-|/BR)LABELD 1.38.(1
     .75+.475×111),[.5]0
[19]
      ('MI-ELF5.1' [FMT 100×(1+L/BR)+.1×(110)×1+R)LABELD(2+
     .7×110),1.55,[.5]10p0
[20]
      1 SPLINE X.[.5]Y
[21]
      'NOISE LEVEL AS A PERCENTAGE OF COMPUTED INTENSITY' C
     ENTER 1 9 1.3 0
[22] (1 20p'NUMBER OF OCCURANCES')LABELD 1 3p1.2 2.8 90
```

[23] MEAN + (+/2 + 0C) + 4094[24] SIGMA+((0C*2)+.*.5)*4093 [25] N+0 L1:AXISD(2+7×(MEAN-1+L/BR)+1+R),(2+.475×N),.2375 90 [26] [27] \rightarrow (10 \neq N \neq N+1)/L1 AXISD T+(2+7×((MEAN-.5×SIGMA)-1+L/BR)+1+R),2.35,T,0,([28] T+7×SIGMA+1+R),.15 [29] AXISD T×(5p1), 1 (1 20p20+, 'MEAN ',, 'ME-EF8.5' [FMT MEAN)LABELD 1 3p2. [30] 5 6 0 (1 200201, 'SIGMA ', 'F8.6' [FMT SIGMA)LABELD 1 302.5 [31] 5.5 0 [32] PLOTEN D V ▼ PDPLOT S:X:N:□IO [1] APLOT OF THE NORMALIZED FIRST DERIVATIVES OF INTENSIT Y [2] AS IS STARTING BIN NUMBER [3] $\Pi T 0 \neq 0$ [4] $x + .75 \times x + (\rho x) \rho [/x + x - (\rho x) \rho]/x + PDER, COMP$ [5] X+X+(pX)p1.82+.78×16 [6] X+X,[.5]2+((7*p<u>COMP</u>)×1p<u>COMP</u>)∘.×6p1 Ē 7 Ī N+0 F 8 1 START [9] L1:1 SPLINE, QQX[;;N] [10] AXISD 2,(1.82+N×.78),.75 90 .75 -.1 [11] AXISD 2,(1.82+N×.78),.75 90 .75 [12] AXISD 9,(1.82+N×.78),.75 90 .75 .1 [13] AXISD 9,(1.82+N×.78),.75 90 .75 •.1 AXISD(2+7×SOL[8] # pCOMP),(2.56+N×.78),.04 90 F147 AXISD 2,(2.21+N×.78).7 0 [15] +(6≠N+N+1)/L1 [16] [17] AXISD(2+7×SOL[8] + pCOMP),1.75 .07 90 AXISD 2 1.75 7 0 1.4 .05 [18] [19] AXISD 2 1.75 7 0 1.4 .05 [20] 'TIME IN MILLISECONDS FROM BEGINNING OF DATA WINDOW' CENTER 1 9 1.3 0 [21] (06 40'COMPVELOTIMEPREIPOSTDIAM')LABELD 1.35.(2.02+.7 8×16).[.5]600 [22] (1 180'PARTIAL DERIVATIVES')LABELD 1 301.2 2.8 90 [23] ('LI4' [FMT S+(15)×.2× 1+0 COMP)LABELD(2+1.4×15).1.55. [.5]5p0 [24] PLOTEN D Δ ▼ PDPLOT2 S:X:N:□IO F17 APLOT OF NORMALIZED FIRST DERIVATIVES OF INTENSITY [2] AFOR TWO-STAR (DC2) SOLUTION F 3 1 AS IS STARTING BIN NUMBER Γ41 □I0+0 Ē5Ī $X + .52 \times X + (\rho X) \rho [/X + X - (\rho X) \rho [/X + (1 1 0 1 1 1 1 1 1 / PDER]],$ COMP [6] X+X+(pX)p1.82+.52×19

```
[7]
      x+x,[.5]2+((7*pCOMP)×1pCOMP) •.×9p1
 [8]
      N \neq 0
 [9]
      START
F101
      L1:1 SPLINE. QOX[::N]
Ē11]
      AXISD 2.(1.82+N×.52)..5 90 .5
                                      .1
F127
      AXISD 2.(1.82+N×.52)..5 90 .5 .1
                                      .1
      AXISD 9,(1.82+N×.52),.5 90 .5
[14]
      AXISD 9,(1.82+N×.52),.5 90 .5
                                      .1
F157
      AXISD(2+7×SOL2[8]*pOBS),(2.56+N×.52),.04 90
      AXISD(2+7×SOL[8]*pOBS),(2.56+N×.52),.04 90
[16]
[17]
      AXISD 2.(2.08+N×.52).7 0
[18]
      +(9≠N+N+1)/L1
      AXISD(2+7×SOL[8]+00BS),1.75 .07 90
[19]
F207
      AXISD(2+7×SOL2[8]+00BS),1.75 .07 90
      AXISD 2 1.75 7 0 1.4 .05
[21]
[22]
      AXISD 2 1.75 7 0 1.4
                             .05
[23]
      'TIME IN MILLISECONDS FROM BEGINNING OF DATA WINDOW'
     CENTER 1 9 1.3 0
     (09 40'COMPVEL2TIM2INT2DIA2VEL1TIM1INT1DIA1')LABELD 1
[24]
     .35,(2.02+.52×19),[.5]9p0
F257
     (1 18p'PARTIAL DERIVATIVES')LABELD 1 3p1.2 2.8 90
     ('LI4' □FMT S+(15)×.2× 1+pCOMP)LABELD(2+1.4×15),1.55,
[26]
     [.5]5n0
[27]
     PLOTEND
     Δ
    V POWERPLOT:R:X:I:DIO
 [1] APOWER SPECTRA OF MODEL CURVE, STAR+SKY, AND OCCULTAT
     TON
 [2]
      AEXECUTE FTPREP IN BEFORE POWERPLOT
 [3]
      □I0+0
 Γ41
      START
 [5]
      PLOTSET .3 0
 [6]
      AXISD 2 1 5 90 1 .1
 [7]
      AXISD 2 1.5 4 90 1 .05
 [8]
      AXISD 4.66 1 5 90 1 .1
 [9]
      AXISD 4.66 1.5 4 90 1 .05
[10]
      AXISD 7.32 1 5 90 1 .1
[11]
      AXISD 7.32 1.5 4 90 1 .05
[12]
      X+(500+1+FTOBS),(500+1+FTOUT),500+1+FTMOD
[13]
             -1+[10⊕([/x),L/x
      R+10*0
[14]
      X+(10@X) +10@R[0]
「15]
      I+1+ 1+, (5×+500)×(1501) •.×1 1
[16]
      1 SPLINE(2+2.5×,(500+X) •.×1 1),1+¢I
[17]
      1 SPLINE(4.66+2.5×, (500+1000+X) •.×1 1), 1+¢I
F187
      1 SPLINE(7.32+2.5×,(1000+x).×1 1),1+¢I
[19]
      AXIS 2 1 2.5 0.(2.5*-/10@R)..1
F207
      AXIS 4.66 1 2.5 0.(2.5+-/10@R)..1
[21]
      AXIS 7.32 1 2.5 0.(2.5*-/10@R)..1
[22]
      AXIS 2 6 2.5 0, (2.5+-/10@R), .1
[23]
      AXIS 4.66 6 2.5 0,(2.5+-/10@R),
[24]
      AXIS 7.32 6 2.5 0, (2.5#-/10@R),
                                       .1
[25]
      ('LI3' [FMT 100×15)LAXISD 1.8 6 1 270
[26]
      (1 18p'FREQUENCY IN HERTZ')LABELD 1 3p1.5 4.95 271
```

```
[27] (1 11p*OCCULTATION*)LAXISD 4.36 3 0 271
[28] (1 10p'STAR + SKY')LAXISD 7.02 3 0 271
[29] (1 110'MODEL CURVE')LAXISD 9.68 3 0 271
[30] PLOTEND
     V
    V RAWPLOT X: DIO; ZRO; SPA; NEG; POS; XPOS; STV; NUMM; N
 [1]
      APLOT OF THE RAW DATA, X (EVERY 4 POINTS)
 [2]
      \Pi T0 + 0
 [3]
      START
      PLOTSET -.4 0
 [4]
     AXISD 2 1.75 7 0 .4375 .15
AXISD 2 1.75 7 0 .875 .05
 [5]
 [6]
                               .05
      AXISD 2 1.75 4.75 90 .475 .15
 [7]
      AXISD 2.(1.75+4.75+20),4.275 90,(4.75+10), .08
 [8]
 [9]
      AXISD 9 1.75 4.75 90 .475 .15
[10]
      AXISD 9,(1.75+4.75+20),4.275 90 .475 .08
     ('ZF3.1' [FMT .1×111)LABELD 1.55,(1.75+.475×111),[.5]
[11]
[12]
      ('LI4' ∏FMT 4096×.1×111)LABELD 9.05,(1.75+.475×111),[
     .5]0
      (*LI4* [FMT 512×18)LABELD(2+.875×18),1.55,[.5]8p0
      'TIME IN MILLISECONDS FROM BEGINNING OF DATA WINDOW'
[14]
     CENTER 1 9 1.3 0
[15]
      (1 20p'NORMALIZED INTENSITY')LABELD 1 3pl.4 2.8 90
[16]
     (1 6p'COUNTS')LABELD 1 3p9.8 4.9 270
     1 SPLINE(8192p0 0 1 0)/(2+7×(14096)+4095).1.75+4.75×X
[17]
     *4095
[18]
     PLOTEND
     V
    V Y+N SMOOTH D:□IO
 [1]
     AN-POINT UNWEIGHTED SMOOTHING OF THE VECTOR D
 [2] DIO+0
 [3]
      Y+(\rho D)\rho((\lfloor .5 \times N) + D), ((-N)+(+/(1N)\phi(N,\rho D)\rho D) + N), (-[.5 \times N)
     ) + D
```

V

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BIOGRAPHICAL SKETCH

Glenn H Schneider (no "." after the H, please note) was born to Elaine and Ira G. Schneider on October 12, 1955, an otherwise un-noteworthy day, except perhaps for the fact that it marked the 463rd anniversary of the discovery of America, as commonly reckoned. At the tender age of three he received his first telescope (Sears, 40X!), and learned early in life the meaning of light pollution.

He attended public school 97 and J.H.S. 135 (also known as the Frank D. Whalen Junior High School, though no one ever knew who Frank D. Whalen was) in New York City. Thanks to the closing of the school during the summer of 1969 he was permitted to retain possession of the school library's copies of <u>The Larousse Encyclopedia of Astronomy</u> and <u>Norton's Star Atlas</u>. Later that year he became a member of the Amateur Observers' Society of New York. These two events forever influenced his life and started him on the path to a career in astronomy.

He graduated from The Bronx High School of Science in June of 1972. Shortly thereafter, he organized an international solar eclipse expedition, in hopes of observing his second total solar eclipse. The clouds which hung ominously over the Gaspé Peninsula that day served to reinforce his devotion to shadow chasing.

In June of 1976, he earned a Bachelor of Science degree in physics from the New York Institute of Technology. Before entering the graduate program in astronomy at the University of Florida in September of 1977, he worked as an APL technical analyst for Warner Computer Systems, Inc.

After eight years, two trips to the South Pole, seemingly endless commuting between Gainesville, Paris, Noordwijc and Bristol, statistically improbable stretches of cloudy weather during scheduled observing runs, and being thwarted at every turn by wayward computers (micro and macro), he expects, finally, to receive the degree of Doctor of Philosophy in August of 1985. He has accepted a position at the Space Telescope Science Institute, working for the Computer Sciences Corporation as a Science and Mission Operations Astronomer.

He has observed every total solar eclipse since March, 1970, and intends to Keep chasing the moon's shadow well into the future. He has a definite preference for Dr. Brown's Celery Tonic (now called Cel-Ray Soda, alas), and egg creams made with Fox's U-Bet chocolate syrup and Good-Health seltzer.

At the moment he is single, but this malady is expected to be cured on June 7, 1985.

I certify that I have read this study and that in my opinion it conforms to acceptable standards of scholarly presentation and is fully adequate, in scope and quality, as a dissertation for the degree of Doctor of Philosophy.

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Hummin Erchan

Heinrich Eichhorn Professor of Astronomy

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Héywood C. Smith' Associate Professor of Astronomy I certify that I have read this study and that in my opinion it conforms to acceptable standards of scholarly presentation and is fully adequate, in scope and quality, as a dissertation for the degree of Doctor of Philosophy.

Ralph G. Selfridge Professor of Computer and Information Sciences

This dissertation was submitted to the Graduate Faculty of the Department of Astronomy in the College of Liberal Arts and Sciences and to the Graduate School and was accepted as partial fulfillment of the requirements for the degree of Doctor of Philosophy.

August 1985

Dean, Graduate School

F 2 7 ACONVERSATIONAL ENTRY OF INITIAL DC PARAMETERS [3] NAME+32+ . Oo + 'NAME OR CATALOG NUMBER OF STAR: ' Γ47 DATE+27+0.000+'U. T. DATE OF OBSERVATION: ' Γ51 COMN+28+F.OoF+'ANY COMMENTS TO BE PRINTED: ' [6] PV+0.00 +'ENTER THE MONOCHROMATIC WAVELENGTH, OR ENTE R O TO USE FILTER TABLE' •(PV≠0)/'FILT+''MONO'''.0pMONO+1 2pPV.1 [7] Ē8] +(PV≠0)/I0 [9] 'ENTER THE NAME OF THE FILTER TABLE TO BE USED: " [10] FILT+ [11] 10:PV+PV.□.00□+'ENTER THE DISTANCE TO THE LUNAR LIMB IN KILOMETERS' [12] BIN+□.00□+'ENTER STARTING BIN NUMBER OF DATA TO BE US ED [13] PV+PV.□.00□+'ENTER NUMBER OF DATA POINTS TO BE USED' [14] OBS + (-1 + PV) + (BIN - 1) + X[15] PV+PV,□,00□+'ENTER THE EFFECTIVE COLOR TEMPERATURE OF THE STAR' [16] PV+PV.□.00□+'ENTER THE SOUARE ROOT OF THE NUMBER OF G RID PTS./OUADRANT' [17] PV+PV, 0+C+206264806.2.000+'ENTER ESTIMATED DIAMETER I N ARC-MILLISECONDS' [18] LIMS+[]+C.Op[+PR+' ENTER UPPER AND LOWER LIMIT' [19] PV+PV, 0,000+'ENTER PRE-EVENT LEVEL OF STAR+SKY SIGNAL [20] LIMS+LIMS. 0.00T+PR [21] PV+PV,□,00□+'ENTER POST-EVENT LEVEL OF LIMB+SKY SIGNA 1.1 [22] LIMS+LIMS, 0,00 +PR [23] PV+PV, □-BIN,00 □+'ENTER ESTIMATED BIN NUMBER OF GEOMET RICAL OCCULTATION ' [24] LIMS+LIMS, $(-1+PV) \circ + 1 - 1 \times 1 \circ 0 \circ 1 \leftrightarrow ?$ PLUS OR MINUS HOW M ANY MILLISECONDS? * PV+PV, (PV[2]×30[+206265),00[+'ENTER SHADOW VELOCITY IN A-SEC/SEC' [26] LIMS+LIMS, (PV[2]× 30[+206265), 0p[+PR [27] PY+PY, 0,000+'ENTER ASSUMED LIMB DARKENING CO-EFFICIEN тţ [28] LIMS+Q6 2pLIMS, 0,0p0+PR [29] PV+1E⁻12 [PV,□,0p□+'ENTER MAXIMUM NUMBER OF ITERATIONS FOR DC PROCEEDURE' [30] PV+PV. 11+-/2 2+ FILT [31] PY+PY, 0,000+'ENTER FRACTIONAL ADJUSTMENT TO BE USED (0 TO 1)' ▼ Y+X LDARKEN G:T [1] AY=NORMALIZED GRID DISTRIBUTION OF LIMB DARKENED STA R AG=STELLAR GRID QUADRANT TO BE LIMB DARKENED (GENERA TED BY: GRID) [3] AX=LIMB DARKEN ING COEFFICIENT

[4] ARESULT OBTAINED FROM A LINEAR LIMB DARKENING LAW

 $[5] \quad Y+Y+4x+/+/Y+Gx(T\circ,\leq T)\times(1-T)\phi\phi(\rho G)\rho(1-X)+X\times20^{-1}O(T-.5)$

```
+0T+11+0G
     Π
    ▼ Y+NPOL V:T:I:□CT
     AINTERPOLATION OF FRESNEL INTENSITY VALUE FOR FRESNE
 [1]
     L NUMBER <V>
     Y + FREN[1] + (- \neq FREN[^1] 0 \circ + I + (1 + .5 \times \rho FREN) - 100 \times T]) \times 100 \times V
 F 2 ]
     -T+.01×L.5+V×100
     Δ
    V OUTPUT2:S:T:PM:CT:W:O
 [1]
     AFORMATTED OUTPUT FOR DC2 SOLUTION
 [2]
      (200' '), 'OCCULTATION OF: ',NAME
      (20p' '), DATE, ' U. T.'
 F 3 T
 Γ41
      (200 ' ').COMN
      1 1
 [5]
 [6]
     IN PUT PARAMETERS:
      1 1
 [7]
      'CENTRAL WAVELENGTH OF PASSBAND: ', (Ov+/×/@FILT).' ANG
 [8]
     STROMS *
 [9]
      'LUNAR LIMB DISTANCE (KM):
                                          ',6 0 TSOL[2]
[10]
     'EFFECTIVE TEMPERATURE STAR-1 : ', VSOL[4]
     *EFFECTIVE TEMPERATURE STAR-2 : *. *SOL2[4]
[11]
     'LIMB DARKENING COEFFICIENT STAR-1: '. FOL[11]
[12]
[13]
      'LIMB DARKENING COEFFICIENT STAR-2: ', # SOL2[11]
      1 1
[14]
[15]
      MODEL PARAMETERS:
      1 1
[16]
[17]
      'NUMBER OF DATA POINTS: ', VSOL[3]
      'NUMBER OF GRID POINTS: ', T(SOL[5]×2)*2
F187
      'NUMBER OF SPECTRAL REGIONS: ', #1+p#FILT
[19]
F207
      WIDTH OF SPECTRAL REGIONS: '.(V-/$(*FILT)[1,2]]to*F
     ILT;1]),' ANGSTROMS'
[21]
      'NUMBER OF ITERATIONS: '. FSOL[12]
[22]
     1 1
[23]
     * SOLUTIONS FOR STAR-1:*
     1 1
F247
[25]
      PM+ ' +/- '
[26]
      'STELLAR DIAMETER (ARC-MILLISECONDS): '.(6 27SOL[6]×2
     06264806), PM.6 2 ¥ ER[1]×206264806
     'BIN OF GEOMETRICAL OCCULTATION:
[27]
                                               '.(7 1¥SOL[9]+B
     IN), PM, 4 1 # ER[4]
[28]
      'SIGNAL LEVEL OF STAR-1 (COUNTS)
                                               '.(6 1¥SOL[7]).
     PM.5 1 #ER[2]
[29]
     ST+×1-T+SOL[10] + PV[10]
[30]
     'OBSERVED LUNAR SHADOW VELOCITY: ',(7 5*SOL[10]),
     PM.(7 5 # ER[5]). * KM/SEC*
[31]
     'PREDICTED LUNAR SHADOW VELOCITY:
                                             *.(7 5*PV[10]).*
      KM/SEC'
[32]
     T+'LOCAL SLOPE OF LUNAR LIMB:
                                              '.'- +'[ST+2]..
     (6 27.5×+/S+(180+01)× 20T+T[+T), DEG'
     T, PM, 5 3 # 1.5 × - / 0 [ 0 + ( + 0 + 180 ) × (SOL [ 10 ] • . + 1 ] × ER[ 5 ] ) + P
[33]
     <u>v[10]</u>
[34] 2 1p' '
```

[35] * SOLUTIONS FOR STAR-2:* . . [36] F 37 T 'STELLAR DIAMETER (ARC-MILLISECONDS): ',(6 27SOL2[6]× 206264806) PM.6 2FER[6]×206264806 [38] 'BIN OF GEOMETRICAL OCCULTATION: ',(7 1 ¥SOL2[9]+ BIN), PM, 4 17ER[8] [39] 'SIGNAL LEVEL OF STAR-2 (COUNTS) ', (6 1 v SOL2[7]) ,PM,4 17<u>ER[7]</u> POST-EVENT LIMB+SKY SIGNAL LEVEL: '.(6 1 #SOL[8]). F40] PM.4 1 VER[3] F417 'OBSERVED LUNAR SHADOW VELOCITY: *.(7 5 #SOL2[10]) .PM.(7 5 # ER[9]). * KM/SEC* F42] ST+×1-T+SOL2[10]+PV[10] [43] T+ LOCAL SLOPE OF LUNAR LIMB: ','- +'[ST+2]., (6 2 ¥.5×+/S+(180 ±01)× 20T+T[±T). DEG * Γ44] T.PM.6 $3 \forall [.5 \times -/0] | 0 \leftrightarrow (\div 0 \div 180) \times (SOL2[10] \circ . +1 ~ 1 \times ER[9]) \div$ PV[10] [45] 2 10' ' [46] . . OTHER PARAMETERS AND DERIVED SOLUTIONS: * 1 1 [47] F487 'POST-EVENT LIMB+SKY SIGNAL LEVEL: '.(6 1¥<u>SOL[8]</u>), PM.4 1 # ER[3] [49] 'TEMPORAL SEPARATION OF THE STARS: '.(6 1+|SOL[9]-S OL2[9]), PM. (4 1 F(ER[4 8]+.*2)*.5), * MS.* 'WEIGHTED MEAN L-RATE: ', (6 4 W+((SOL[10], SOL2[10])+. [50] *ER[5 9]) ++/ + ER[5 9]) . * KM/SEC* [51] 'SPATIAL SEPARATION FROM R-RATES WHICH ARE: ' [52] T+' PREDICTED ('.(*Q).' ARC SEC/SEC): '.6 2*(|SOL2[9]-SOL[9])×0+206265×30+/PV[10 2] [53] T. MILLISECONDS OF ARC' T+' WEIGHTED MEAN (',(6 4¥Q),' ARC SEC/SEC): ',6 2₹([54] |SOL2[9]-SOL[9])×0+206265×30W ± PV[2] [55] T. ' MILLISECONDS OF ARC' [56] 'BRIGHTNESS RATIO (BRIGHTER/FAINTER): ',6 4T++/(SOL[7],SOL2[7])[#SOL[7],SOL2[7]] [57] 'MAGNITUDE DIFFERENCE: *.6 4▼(100*.2)⊕T [58] 2 10' ' [59] . PHOTOMETRIC STATISTICS: * 1 1 [60] F617 'SUM OF SQUARES OF RESIDUALS: ', VS+(OBS-COMP)+.*2 [62] 'SIGMA (STANDARD ERROR): '.▼S+(S+ 1+pOBS)*.5 'NORMALIZED STANDARD ERROR: [63] ', ▼S+S+SOL[7]+SOL2[7] [64] 'PHOTOMETRIC (SIGNAL+NOISE)/NOISE RATIO: ', *(1+S) *S '(INTENSITY CHANGE)/BACKGROUND: ', *S+(SOL[7]+SOL2[[65] 7]) + SOL[8] [66] 'CHANGE IN MAGNITUDE: *, ¥-(+100++5)⊕1+S 2 10' ' [67] [68] 1 SUPPLEMENTAL STATISTICS: * 1 1 [69] [70] 'VARIANCE/CO-VARIANCE MATRIX:' [71] CT+9 11ρ1+□+' DIAMETER1 INTENSITY1 SKY INTS OCC TI ME1 L-RATE1 DIAMETER2 INTENSITY2 OCC TIME2 L-RAT 'MD-DE11.4' DFMT COV [72] Ē73Ī 1 1

```
[74] (13p' '), CT
[75] ((110/\CT=' ') OCT), 'ME-EF11.6' [FMT CM COV
[76]
    RESTAR
     Δ
    V OUTPUT:S:T:PM:CT
 [1] APRODUCES A FORMATTED REPORT ON THE STELLAR SOLUTION
 [2]
     (20p' '), 'OCCULTATION OF: ', NAME
 [3] (20p' '), DATE, ' U. T.'
 [4]
     (20p' '), COMN
 [5]
      1 1
 [6]
      ' INPUT PARAMETERS: '
 [7]
     . .
 [8]
     'CENTRAL WAVELENGTH OF PASSBAND:', (0 +/×/ FILT),' ANG
     STROMS!
 [9]
     'LUNAR LIMB DISTANCE (KM):
                                         ',6 0 # SOL[2]
     'EFFECTIVE TEMPERATURE (KELVIN): ', VSOL[4]
[10]
                                        '. # SOL[11]
[11]
      'LIMB DARKENING COEFFICIENT:
[12]
[13]
     'MODEL PARAMETERS:'
     1 1
[14]
[15]
      'NUMBER OF DATA POINTS: '. #SOL[3]
[16]
     'NUMBER OF GRID POINTS: '. F(SOL[5]×2)*2
     'NUMBER OF SPECTRAL REGIONS: ', #1+p#FILT
[17]
[18]
     WIDTH OF SPECTRAL REGIONS: '. (*-/$(*FILT)[1,2]]+0*F
     ILT;1]),' ANGSTROMS'
[19]
      'NUMBER OF ITERATIONS: ', * SOL[12]
[20]
     1 1
[21]
     ' SOLUTIONS: *
[22]
     1 1
[23]
     PM+' +/-'
[24]
      'STELLAR DIAMETER (ARC-MILLISECONDS): '.(6 2 SOL[6]×2
     06264806), PM, 6 2 FER[1]×206264806
[25]
     'BIN OF GEOMETRICAL OCCULTATION:
                                             '.(7 1 #SOL[9]+B
     IN), PM, 4 1 FER[4]
[26]
      'PRE-EVENT STAR+SKY SIGNAL LEVEL:
                                             '.(6 1 #SOL[7]).
     PM,4 17ER[2]
[27]
     'POST-EVENT LIMB+SKY SIGNAL LEVEL:
                                             ',(6 1*<u>SOL[8]</u>),
     PM,4 1 # ER[3]
[28]
     ST+×1-T+SOL[10] + PV[10]
[29]
     'OBSERVED LUNAR SHADOW VELOCITY:
                                            ',(7 5¥<u>SOL</u>[10]),
     PM, (7 5 FER[5]), ' KM/SEC'
[30]
      'PREDICTED LUNAR SHADOW VELOCITY: '.(7 5*PV[10]).'
      KM/SEC'
[31] T+'LOCAL SLOPE OF LUNAR LIMB:
                                              ','- +'[ST+2]..
     (6 2 .5 ×+/S+(180 +01) × 20T+T[+T), ' DEG'
F 32 ]
     T, PM, 5 37 |.5×-/QL |Q+(+0+180)×(SOL[10] .+1 1×ER[5])+P
     ¥[10]
[33]
      2 10' '
[34]
      'VARIANCE-COVARIANCE MATRIX:'
[35]
     'MO-OF12.5' OFMT COV
[36]
     2 1 p ' '
F371
     'CORRELATION MATRIX:'
[38]
      60+,CT+5 10p' DIAMETER
                                  *+LIMB LIMB+SKY TIME OC
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C VELOCITY ' [39] ((1⊥\oh)(CT=' '))¢CT),'MŪ-ŪF10.5' □FMT CM <u>COV</u> [40] 2 lp' ' 'SUM OF SQUARES OF RESIDUALS: ', FS+(OBS-COMP)+.*2 [41] 'SIGMA (STANDARD ERROR): ', #S+(S+ 1+pOBS)*.5 [42] [43] 'NORMALIZED STANDARD ERROR: ', FS+S+-/SOL[7 8] 'PHOTOMETRIC (SIGNAL+NOISE)/NOISE RATIO: '. #(1+S)#S F447 '(INTENSITY CHANGE)/BACKGROUND: ', #S+(-/SOL[7 8])* F457 SOL[8] '.▼-(+100*+5)@1+S [46] 'CHANGE IN MAGNITUDE: [47] RESTAB Δ V RESTAB:0:N APRODUCES A TABLE OF OBS, COMP, AND RESIDUALS F11 4 10' ' [2] [3] (29+1 1), 'TABLE 5-XXX' [4] ' ',NAME,': OBSERVATIONS, COMPUTED VALUES, AND RESID HALS FROM BIN '. FBIN [5] '' [6] 720*=* [7] 'NUM OBS COMP RESID NUM OBS COMP RESID NU M OBS COMP RESID' [8] '--- ------- ---- ----- ------ ---- ----- ------! [9] Q+((3×N+[(pOBS)+3),4)+(-1+1pOBS),OBS,COMP,[1.5]OBS-CO MP $[10] Q+Q[1N;],((N,0)+(-N,0)+Q),(2\times N,0)+Q$ '3(BI3.X1.BI4,X1,BF6.1,X1,M2-EBF6.1,X3)' [FMT Q [11] [12] 72p'=' V ▼ START [1] AOVERALL CONTROL FUNCTION FOR DC EXECUTION [2] 'REDUCTION RUN FOR THE STAR: '.NAME [3] DC PV [4] DC PV+(-3+3+, SOLS[SOLH+SOLS[:2]11:]), 10, PV[13]. [5] DC(-3+3+.SOLS[SOLF+SOLS[:2]11:]).0.PV[13].5 [6] OUTPUT V V START2 [1] **AOVERALL CONTROL FUNCTION FOR DC2 EXECUTION** 'REDUCTION RUN FOR THE STAR ',NAME [2] [3] PV2 DC2 PV V ▼ Y+R WIDE P:I:T [1] ACOMPUTATION OF WIDEBAND FRESNEL PATTERN [2] AP IS THE CURRENT PARAMETER VECTOR (PV) [3] AR=CONVOLVED SYSTEM RESPONCE AND SOURCE SPECTRAL ENER GY CURVES [4] I+1+Y+P[3]00 [5] AS++/2×P[11]LDARKEN GRID P[5]

[6] W0:Y+Y+R[I]×(NPOL FNOS P+(<u>sFILT</u>)[I+I+1;1],1+P)+.×(\$T) ,T+\$ +(IF1+psFILT)/W0 v

NAME: VFILTER SHAPE: 53 2	NAME: BFILTER SHAPE: 42 2	NAME: NARROWV SHAPE: 5 2
4450 0.0001294 4500 0.0002588	3600 0.0000953 3650 0.0004767	5295 0.0621118 5345 0.2251553 5205 0.4687578
4600 0.0005175 4650 0.0006469	3750 0.0023834 3800 0.0044809	5445 0.2251553 5495 0.0388199
4700 0.0007763 4750 0.0011644 4800 0.0016820	3850 0.0115359 3900 0.0207837 3950 0.0397559	NAME: NARROWB
4850 0.0023289 4900 0.0047872	4000 0.0469063 4050 0.0489084	SHAPE: 5 2
4950 0.0112563 5000 0.0208306	4100 0.0499571 4150 0.0505291 4200 0.0511965	4600 0.0658106 4650 0.2247191 4700 0.4558587
5100 0.0380386 5150 0.0450252	4250 0.0513872 4300 0.0514825	4750 0.2247191 4800 0.0288925
5200 0.0513650 5250 0.0548583	4350 0.0510058 4400 0.0497664	
5300 0.0555052 5350 0.0553759 5400 0.0545996	4500 0.0475850 4500 0.0451902 4550 0.0427114	
5450 0.0534351 5500 0.0507181	4600 0.0400419 4650 0.0376585	
5600 0.0452840 5650 0.0426963	4750 0.0331776 4800 0.0295548	
5700 0.0398499 5750 0.0364860 5800 0.0337689	4850 0.0271713 4900 0.0244065 4950 0.0219277	
5850 0.0307931 5900 0.0282055	5000 0.0192583 5050 0.0164935	
5950 0.0252297 6000 0.0225126 6050 0.0196076	5100 0.0137287 5150 0.0111545 5200 0.0085804	
6100 0.0166904 6150 0.0142321	5250 0.0058156 5300 0.0034322	
6200 0.0122914 6250 0.0104800 6300 0.0087980	5350 0.0020974 5400 0.0015254 5450 0.0008580	
6350 0.0071161 6400 0.0054341	5500 0.0004767 5550 0.0002860	
6450 0.0041403 6500 0.0029758 6550 0.0024583	5600 0.0001907 5650 0.0000953	
6600 0.0020701 6650 0.0016820		
6750 0.0012938 6800 0.0010351		
6850 0.0007763 6900 0.0006469		
7000 0.0002588		

7050 0.0001294

APPENDIX E LISTING OF THE APL WORKSPACE OCCPLOTS

```
V Y+N BIN D;□IO;T
 [1]
      ABIN DATA FOR INTO CLASSES FOR HISTOGRAM
 [2]
      AD IS DATA VECTOR: N IS NUMBER OF CLASSES
 [3] [10+1
       Y+(T+D[1]-.5×1+T),[1.5]+/(<sup>-</sup>1+1N) •.=+/((1+D)+T++\Np(([
 F41
      /D)-|/D) +N) ∘. <D+D[▲D]
      Δ
    V D CENTER S:R
 [1] ACENTER TEXT LABEL
 [2] R+S[0]+.5×S[1]-(pD)×8*53
 [3] ((1.oD)oD)LABELD 1 3oR, 2+S
      Π
    V Y+N ESMOOTH D:□IO:T:V
 [1] AN-POINT WEIGHTED SMOOTHING OF THE VECTOR D
 [2]
     A(1+N) IS NUMBER OF TERMS TO USE NUMERICALLY IN E SMO
      OTHING
 [3]
     e(1+N) IS THE 1/E FOLDING SPACING WRT THE DATA SPACIN
     G
 Γ41
      \Pi T 0 + 0
 [5]
      Y \leftarrow ((V \leftarrow 1.5 + .5 \times 1 + N) p + Y), Y \leftarrow (T + + / T \leftarrow * - (+ 1 + N) \times (1 + N) + . \times (1 + 1 + N)
     N) Φ(1+N, pD) pD
      Y+(-3×V)+Y
 F 6 1
 [7] Y \leftarrow (\rho D) \rho Y, (2 \times V) \rho (+/(-1 + N) + D) + 1 + N
      V
    ▼ Z+FFT X:T:I:G:N:Q:□IO
 [1]
      AFAST FOURIER TRANSFORM OF X; FROM RGS
       AX VECTOR OF REALS OR X[:0]=REAL X[:1]=IMAGINARY
 [2]
       □I0+0
 Ē4Ī
       \rightarrow (2=00X)/LO
       X+X,[.5]0
 [6]
       L0:T+210(Gp2)T1I+2*G+[2@1+pX
      X+((N+I+2),2 2)p((I,2)+X+I+.5)[T;]
 [8]
      T+02 10.00(N+T)+2×N
 [9]
      Z+(+/[1]X),[.5]-/[I+1]X
[10]
      LA: Z+((N+N+2),(2×I+I×2),2)oZ
[11]
      X + (0, I, 0) + Z
```
```
[12]
      Z+(0,-I,0)+Z
[13]
       Z+(Z+Q), Z-Q+(+/Q\times X), [1.5]-/(Q+(N,I,2)p(I,2)+T)\times \phi X
[14]
      \rightarrow (1 < N)/LA
[15]
      Z+((2×I,1)pZ)[2⊥⊖(Gp2)T12*G;]
      Δ
    V Z+FFTI X;T;I;G;N;Q;□IO
       AFAST INVERSE FOURIER TRANSFORM OF X: FROM RGS
 [1]
       AX IS VECTOR OF REALS, OR X[;0]=REAL X[;1]=IMAGINARY
 [2]
 [3]
       \Pi I 0 + 0
 Γ41
       \rightarrow (2=00X)/L0
       x+x,[0.5] 0
 F 5 1
 [6]
       L0:T+210(Gp2)T1I+2*G+[2@1+pX
       X+((N+I+2), 2, 2)\rho((I,2)+X+I+.5)[T;]
 [8]
      T + b_2 = 1 \circ .00 (N + T) = 2 \times N
 [9]
       z+(+/[1]x),[.5]-/[1+1]x
       LA: Z+((N+N+2), (2\times I+I\times 2), 2) \rho Z
[10]
[11]
       X+(0,I,0)+Z
       Z + (0, -I, 0) + Z
[12]
[13]
       Z+(Z+O), Z-O+(+/O\times X), [1,5]-/(O+(N,I,2))(I,2)+T)\times \phi X
[14]
      +(1 <N)/LA
      Z + ((2 \times I, 1) \rho Z) [2 \perp \Theta (G \rho 2) T 1 2 * G;]
[15]
     V FITPLOT X: IO; ZRO; SPA; XPOS; STV; NUMN; N
 [1]
       APLOT OF THE FITTED OCCULTATION CURVE
 [2]
       AX=STARTING BIN NUMBER: OBS, COMP, SOL MUST BE GLOBAL
 [3]
       □I0+0
 [4]
       x+x,(x+sol[2]),(x+sol[8]),sol[9],sol[5]\times 206264810
 [5]
       START
       PLOTSET -.3 0
 [6]
       AXISD 2 1.75 7 0 .35 .15
AXISD 2 1.75 7 0 1.4 .05
 F81
       AXISD 2 1.75 4.75 90 .475 .15
 Γ91
       AXISD 2,(1.75+4.75+20),4.275 90,(4.75+10), .08
F107
[11]
       AXISD 9 1.75 4.75 90 .475 .15
       AXISD 9,(1.75+4.75+20).4.275 90 .475 .08
[12]
[13]
       ZRO + 2 + (-/X[2 0]) \times 7 + -/X[1 0]
[14]
       SPA+10×*X[3]×(-/X[1 0])*7
[15]
       AXISD ZR0.6.5.(ZR0-2).180.SPA..15
[16]
       AXISD ZR0,6.5,(9-ZR0),0,SPA, .15
[17]
       ('ZF3.1' [FMT .1×111)LABELD 1.55,(1.75+.475×111),[.5]
F187
      (*LI4* □FMT(<u>|/OBS</u>)+((<u>[/OBS</u>)-<u>|/OBS</u>)×.1×111)LABELD 9.05
      ,(1.75+.475×111),[.5]0
[19]
       ('LI4' [FMT X[0]+(15)×.2×1+-/X[1 0])LABELD(2+1.4×15)
      ,1.55,[.5]500
[20]
      XPOS+ZRO-SPA×L(ZRO-2) * SPA
[21]
      STV+-10×| (ZRO-2) + SPA
[22]
      NUMM+([(ZRO-2) + SPA)+[(9-ZRO) + SPA
[23]
       ('MM-MLI4' DFMT STV+10×1NUMM)LABELD((XPOS+NUMM++\0,NU
      MMpSPA).[.5]6.55).0
       'TIME IN MILLISECONDS FROM BEGINNING OF DATA WINDOW'
[24]
      CENTER 1 9 1.3 0
```

```
'ANGULAR SEPARATION IN MILLISECONDS OF ARC' CENTER 1
[25]
     9 6.75 0
      (1 20p'NORMALIZED INTENSITY')LABELD 1 3p1.4 2.8 90
[26]
      (1 6p'COUNTS')LABELD 1 3p9.8 4.9 270
[27]
     N \neq 0
[28]
      L1:AXISD ZRO,(2+.475×N),(.475+2),270
[29]
[30]
     +(10≠N+N+1)/L1
      1 SPLINE(2+7×(1pOBS)*1+pOBS),1.75+4.75×(OBS-L/OBS)*[
[31]
     /OBS-L/OBS
     1 SPLINE(2+7×(10 COMP) * 1+0 COMP), 1.75+4.75×(COMP-L/OBS)
[32]
     ) + [/OBS-1/OBS
      AXISD 8.5,(1.75+4.75×(SOL[6]-L/OBS)+([/OBS)-L/OBS),.5
[33]
      AXISD 2.(1.75+4.75×(SOL[7]-L/OBS)*([/OBS)-L/OBS),.5 0
[34]
[35]
      AXISD ZR0,6.32,(X[4]×.1×SPA),0
[36]
      PLOTEN D
     Δ
    V FITPLOT2 X: DIO: ZRO: SPA: XPOS: STV: NUMN:N
 [1]
      APLOT OF THE FITTED TWO-STAR OCCULTATION CURVE
 [2]
      AX=STARTING BIN NUMBER: OBS, COMP, SOL MUST BE GLOBAL
 [3]
      □I0+0
      X+X,(X+SOL[2]),(SOL[8]+X2+X),SOL[9],SOL[5]×206264810
 [4]
      START
 [6]
     PLOTSET .3 0
 [7]
      AXISD 2 1.75 7 0 .35 .15
                             -.05
 [8]
      AXISD 2 1.75 7 0 1.4
      AXISD 2 1.75 4.75 90 .475 -.15
 [9]
[10]
     AXISD 2,(1.75+4.75+20),4.275 90,(4.75+10), .08
[11]
      AXISD 9 1.75 4.75 90 .475 .15
      AXISD 9,(1.75+4.75+20),4.275 90 .475 .08
[12]
      ZRO+2+(-/X[2 0])\times7+-/X[1 0]
[13]
      SPA+10×*X[3]×(-/X[1 0])*7
[14]
[15]
      AXISD ZR0.6.5.(ZR0-2).180.SPA..15
[16]
      AXISD ZR0,6.5,(9-ZR0),0,SPA, .15
[17]
      ('ZF3.1' [FMT .1×111)LABELD 1.55,(1.75+.475×111),[.5]
     0
[18]
     ('LI4' □FMT(<u>L</u>/<u>OBS</u>)+((<u>Γ</u>/<u>OBS</u>)-<u>L</u>/<u>OBS</u>)×.1×111)LABELD 9.05
     ,(1.75+.475×111),[.5]0
[19]
     ('LI4' □FMT X[0]+(15)×.2×<sup>-</sup>1+-/X[1 0])LABELD(2+1.4×15)
     .1.55.[.5]500
F207
      XPOS+ZRO-SPA×L(ZRO-2) + SPA
[21]
     STV+-10×L(ZRO-2)*SPA
[22]
     NUMM+(L(ZRO-2)*SPA)+L(9-ZRO)*SPA
      ('ME-ELI4' FMT STV+10×1NUMM)LABELD((XPOS+NUMM++\0,NU
[23]
     MMpSPA),[.5]6.55),0
[24]
      'TIME IN MILLISECONDS FROM BEGINNING OF DATA WINDOW'
     CENTER 1 9 1.3 0
[25]
      'ANGULAR SEPARATION IN MILLISECONDS OF ARC' CENTER 1
     9 6.75 0
[26]
     (1 20p 'NORMALIZED INTENSITY')LABELD 1 3p1.4 2.8 90
      (1 6p'COUNTS')LABELD 1 3p9.8 4.9 270
[28]
      N \neq 0
[29]
      L1:AXISD ZRO, (2+.475×N), (.475+2), 270
```

```
\lceil 30 \rceil \rightarrow (10 \neq N \neq N + 1) / L1
[31] 1 SPLINE(2+7×(100BS)+ 1+00BS), 1.75+4.75×(0BS-L/0BS)+[
     /OBS-L/OBS
     1 SPLINE(2+7×(10 COMP)+ 1+0 COMP), 1.75+4.75×(COMP-L/OBS
F 32 ]
     ) + [/OBS-L/OBS
     AXTSD 8.5.(1.75+4.75×((SOL[7]+SOL2[6]+SOL[6])-L/OBS)+
[33]
     ([/OBS)-1/OBS)..5 0
     AXISD 2,(1.75+4.75×(SOL[7]-L/OBS)+([/OBS)-L/OBS),.5 0
[34]
[35] AXISD ZR0,6.32,(X[4]×.1×SPA).0
[36] x + x2, (x2 + sol2[2]), (x2 + sol2[8]), sol2[9], sol2[5] \times 206264
     810
[37] N+0×1+ZRO+2+(-/X[2 0])×7+-/X[1 0]
[38]
     SPA+10×*X[3]×(-/X[1 0])*7
[39] L2:AXISD ZRO.(2+.475×N).(.475+2).270
\lceil 40 \rceil \rightarrow (10 \neq N + N + 1) / L2
[41] AXISD ZR0,6.32,(X[4]×.1×SPA),0
[42] PLOTEND
     V
    V R←FT Y:M:N:□IO
 [1] ADISCRETE FOURIER TRANSFORMATION OF EVENLY SPACED REA
     T. DATA
 [2]
     \Pi T 0 + 0
 [3] M+(12×N)×O+N+(pY)+2
 [4] R+1 2ρΦ, (+/Y)+2×N
[5] L1:R+R.[0](*N)×+/(1 10.×Y)×2 10.0(1+pR)×M
 [6]
    +(N><sup>-</sup>1+1+pR)/L1
    R[N:]+.5 0×. 1 2+R
     Ω
    V FTPREP X:□IO:T
 [1]
      ASETS UP GLOBAL FOR POWER SPECTRA ANALYSIS
F 2 1
      AX IS CENTER OF TRANSFORM WINDOW FROM DATA SET
 [3]
      □I0+0
[4]
      FTOBS+.5×+/(FFT FTDATA+1024+(~512+514[X|3595)+CH1)*2
[5]
     [6]
[7]
      FTOUT+.5×(FFT FTOUT)+.*2
     T+[1+SOL[6]+(COMP[0]-SOL[6]) \times T+[/T+T-L/T++++1+L.5\times1]
۲8 T
     024-0COMP
 [9]
     FTMOD+T.COMP
[10]
      FTMOD + FTMOD, 1 + SOL[7] + (-SOL[7] - 1 + FTMOD) \times \Phi T + [/T + T - ]/T +
     + \+ \ 11+ 1024- p FTMOD
     FIMOD+.5×(FFT FIMOD)+.*2
[11]
    V IDENT
 [1]
      .
          ***LUNAR OCCULTATION PLOTTING WORKSPACE****
[2]
      1
              VERSION: B01/HARRIS'
[3]
      1
              WSID: DISPLOT'
[4]
      .
              LAST REVISION: 10 APRIL 85'
Ē 5 Ī
      .
[6]
      1
              NOTE: THIS WORKSPACE IS DIO+0'
     Ω
```

```
▼ INTPLOT X:□IO
      AINTEGRATION PLOT OF THE DATA VECTOR X
 F11
 [2]
      \Pi I0 + 0
 Ē31
      START
 Ē41
      AXISD 2 1.75 7 0 .4375 .15
 [5]
      AXISD 2 1.75 7 0 .875 .05
      AXISD 2 1.75 4.75 90 .475 .15
 Γ61
      AXISD 2.(1.75+4.75+20),4.275 90,(4.75+10), .08
      ('ZF3.1' [FMT .1×111)LABELD 1.55,(1.75+.475×111),[.5]
 [8]
     0
     ('LI4' [FMT((pX) +8)×18)LABELD(2+.875×18),1.55,[.5]8p0
 [9]
      'TIME IN MILLISECONDS FROM BEGINNING OF DATA WINDOW'
ſ10]
     CENTER 1 9 1.3 0
     (1 200'NORMALIZED INTENSITY')LABELD 1 301.25 3.22 90
[11]
[12]
      x \leftarrow (1/x) - x \leftarrow x \leftarrow (+/x) \neq \rho x \leftarrow x [2 3] \cdot 2 \neq x
[13]
      1 SPLINE (2+7×(1pX)*1+pX),1.75+4.75×X*[/X
[14]
     PLOTEND
     V
    V F+X INVFT R:□IO
 [1] AINVERSE FOURIER TRANSFORMATION FROM COEFFICIENT MATR
     тΧ
 [2]
      □I0+0
 [3]
     X+2 0 102 10.0X0.×1+11+pR
 [4]
     F \leftarrow (1 \ 1 + R) + + / + / ((\rho X) \rho 1 \ 0 + R) \times X
    V NOISEPLOT X: OC: PD: R: MEAN: SIGMA; Y; BR; N
 [1]
      ANOISE FIGURE FOR LUNAR OCCULTATION
 [2]
       AX IS BIN NUMBER OF START OF DATA USED IN SOLUTION
 [3]
      AMUST HAVE SOL OBS AND COMP AS GLOBAL VARIABLES
 [4]
      OBSA = SOL[2] + (X-1) + CH1
      AX[1]=FIRST BIN OF DATA X[2]=LENGTH OF DATA
 [5]
      OC+((X-1)+CH1)-SOL[6]
 [6]
      OC \leftarrow OC, (OBS - COMP), ((-1 + SOL[2] + X) \neq CH1) - SOL[7]
 [8]
      PD+BR+$50 BIN OC+2+OC+CH1
[9]
      PD[0:]+PD[0:]-PD[0:0]
[10]
      PD + PD \times Q(\phi_0 PD)_07 \quad 4.75 \pm [/PD
[11]
      X+2+1+<sup>-</sup>1+,(7×(151)+50)•.×1 1
[12]
      Y+1.75+.PD[1:] •.×1 1
[13]
      START
      PLOTSET -.1 0
[14]
[15]
      AXISD 2 1.75 7 0 .14_.15
[16]
      AXISD 2 1.75 7 0 .7
                             -.05
[17]
      AXISD 2 1.75 4.75 90 .475 .15
[18]
      ('ZI4' □FMT .1×(111)×0 1/R+([/BR)-|/BR)LABELD 1.38.(1
     .75+.475×111),[.5]0
[19]
      ('MI-ELF5.1' [FMT 100×(1+L/BR)+.1×(110)×1+R)LABELD(2+
     .7×110),1.55,[.5]10p0
[20]
      1 SPLINE X.[.5]Y
[21]
      'NOISE LEVEL AS A PERCENTAGE OF COMPUTED INTENSITY' C
     ENTER 1 9 1.3 0
[22] (1 20p'NUMBER OF OCCURANCES')LABELD 1 3p1.2 2.8 90
```

[23] MEAN + (+/2 + 0C) + 4094[24] SIGMA+((0C*2)+.*.5)*4093 [25] N+0 L1:AXISD(2+7×(MEAN-1+L/BR)+1+R),(2+.475×N),.2375 90 [26] [27] \rightarrow (10 \neq N \neq N+1)/L1 AXISD T+(2+7×((MEAN-.5×SIGMA)-1+L/BR)+1+R),2.35,T,0,([28] T+7×SIGMA+1+R),.15 [29] AXISD T×(5p1), 1 (1 20p20+, 'MEAN ',, 'ME-EF8.5' [FMT MEAN)LABELD 1 3p2. [30] 5 6 0 (1 200201, 'SIGMA ', 'F8.6' [FMT SIGMA)LABELD 1 302.5 [31] 5.5 0 [32] PLOTEN D V ▼ PDPLOT S:X:N:□IO [1] APLOT OF THE NORMALIZED FIRST DERIVATIVES OF INTENSIT Y [2] AS IS STARTING BIN NUMBER [3] $\Pi T 0 \neq 0$ [4] $x + .75 \times x + (\rho x) \rho [/x + x - (\rho x) \rho]/x + PDER, COMP$ [5] X+X+(pX)p1.82+.78×16 [6] X+X,[.5]2+((7*p<u>COMP</u>)×1p<u>COMP</u>)∘.×6p1 Ē 7 Ī N+0 F 8 1 START [9] L1:1 SPLINE, QQX[;;N] [10] AXISD 2,(1.82+N×.78),.75 90 .75 -.1 [11] AXISD 2,(1.82+N×.78),.75 90 .75 [12] AXISD 9,(1.82+N×.78),.75 90 .75 .1 [13] AXISD 9,(1.82+N×.78),.75 90 .75 •.1 AXISD(2+7×SOL[8] # pCOMP),(2.56+N×.78),.04 90 F147 AXISD 2,(2.21+N×.78).7 0 [15] +(6≠N+N+1)/L1 [16] [17] AXISD(2+7×SOL[8] + pCOMP),1.75 .07 90 AXISD 2 1.75 7 0 1.4 .05 [18] [19] AXISD 2 1.75 7 0 1.4 .05 [20] 'TIME IN MILLISECONDS FROM BEGINNING OF DATA WINDOW' CENTER 1 9 1.3 0 [21] (06 40'COMPVELOTIMEPREIPOSTDIAM')LABELD 1.35.(2.02+.7 8×16).[.5]600 [22] (1 180'PARTIAL DERIVATIVES')LABELD 1 301.2 2.8 90 [23] ('LI4' [FMT S+(15)×.2× 1+0 COMP)LABELD(2+1.4×15).1.55. [.5]5p0 [24] PLOTEN D Δ ▼ PDPLOT2 S:X:N:□IO F17 APLOT OF NORMALIZED FIRST DERIVATIVES OF INTENSITY [2] AFOR TWO-STAR (DC2) SOLUTION F 3 1 AS IS STARTING BIN NUMBER Γ41 □I0+0 Ē5Ī $X + .52 \times X + (\rho X) \rho [/X + X - (\rho X) \rho [/X + (1 1 0 1 1 1 1 1 1 / PDER]],$ COMP [6] X+X+(pX)p1.82+.52×19

```
[7]
      x+x,[.5]2+((7*pCOMP)×1pCOMP) •.×9p1
 [8]
      N \neq 0
 [9]
      START
F101
      L1:1 SPLINE. QOX[::N]
Ē11]
      AXISD 2.(1.82+N×.52)..5 90 .5
                                      .1
F127
      AXISD 2.(1.82+N×.52)..5 90 .5 .1
                                      .1
      AXISD 9,(1.82+N×.52),.5 90 .5
[14]
      AXISD 9,(1.82+N×.52),.5 90 .5
                                      .1
F157
      AXISD(2+7×SOL2[8]*pOBS),(2.56+N×.52),.04 90
      AXISD(2+7×SOL[8]*pOBS),(2.56+N×.52),.04 90
[16]
[17]
      AXISD 2.(2.08+N×.52).7 0
[18]
      +(9≠N+N+1)/L1
      AXISD(2+7×SOL[8]+00BS),1.75 .07 90
[19]
F207
      AXISD(2+7×SOL2[8]+00BS),1.75 .07 90
      AXISD 2 1.75 7 0 1.4 .05
[21]
[22]
      AXISD 2 1.75 7 0 1.4
                             .05
[23]
      'TIME IN MILLISECONDS FROM BEGINNING OF DATA WINDOW'
     CENTER 1 9 1.3 0
     (09 40'COMPVEL2TIM2INT2DIA2VEL1TIM1INT1DIA1')LABELD 1
[24]
     .35,(2.02+.52×19),[.5]9p0
F257
     (1 18p'PARTIAL DERIVATIVES')LABELD 1 3p1.2 2.8 90
     ('LI4' □FMT S+(15)×.2× 1+pCOMP)LABELD(2+1.4×15),1.55,
[26]
     [.5]5n0
[27]
     PLOTEND
     Δ
    V POWERPLOT:R:X:I:DIO
 [1] APOWER SPECTRA OF MODEL CURVE, STAR+SKY, AND OCCULTAT
     TON
 [2]
      AEXECUTE FTPREP IN BEFORE POWERPLOT
 [3]
      □I0+0
 Γ41
      START
 [5]
      PLOTSET .3 0
 [6]
      AXISD 2 1 5 90 1 .1
 [7]
      AXISD 2 1.5 4 90 1 .05
 [8]
      AXISD 4.66 1 5 90 1 .1
 [9]
      AXISD 4.66 1.5 4 90 1 .05
[10]
      AXISD 7.32 1 5 90 1 .1
[11]
      AXISD 7.32 1.5 4 90 1 .05
[12]
      X+(500+1+FTOBS),(500+1+FTOUT),500+1+FTMOD
[13]
             -1+[10⊕([/x),L/x
      R+10*0
[14]
      X+(10@X) +10@R[0]
「15]
      I+1+ 1+, (5×+500)×(1501) •.×1 1
[16]
      1 SPLINE(2+2.5×,(500+X) •.×1 1),1+¢I
[17]
      1 SPLINE(4.66+2.5×, (500+1000+X) •.×1 1), 1+¢I
F187
      1 SPLINE(7.32+2.5×,(1000+x).×1 1),1+¢I
[19]
      AXIS 2 1 2.5 0.(2.5#-/10@R)..1
F207
      AXIS 4.66 1 2.5 0.(2.5+-/10@R)..1
[21]
      AXIS 7.32 1 2.5 0.(2.5*-/10@R)..1
[22]
      AXIS 2 6 2.5 0, (2.5+-/10@R), .1
[23]
      AXIS 4.66 6 2.5 0,(2.5+-/10@R),
[24]
      AXIS 7.32 6 2.5 0, (2.5#-/10@R).
                                       .1
[25]
      ('LI3' [FMT 100×15)LAXISD 1.8 6 1 270
[26]
      (1 18p'FREQUENCY IN HERTZ')LABELD 1 3p1.5 4.95 271
```

```
[27] (1 11p*OCCULTATION*)LAXISD 4.36 3 0 271
[28] (1 10p'STAR + SKY')LAXISD 7.02 3 0 271
[29] (1 110'MODEL CURVE')LAXISD 9.68 3 0 271
[30] PLOTEND
     V
    V RAWPLOT X: DIO; ZRO; SPA; NEG; POS; XPOS; STV; NUMM; N
 [1]
      APLOT OF THE RAW DATA, X (EVERY 4 POINTS)
 [2]
      \Pi T0 + 0
 [3]
      START
      PLOTSET -.4 0
 [4]
     AXISD 2 1.75 7 0 .4375 .15
AXISD 2 1.75 7 0 .875 .05
 [5]
 [6]
                               .05
      AXISD 2 1.75 4.75 90 .475 .15
 [7]
      AXISD 2.(1.75+4.75+20),4.275 90,(4.75+10), .08
 [8]
 [9]
      AXISD 9 1.75 4.75 90 .475 .15
[10]
      AXISD 9,(1.75+4.75+20),4.275 90 .475 .08
     ('ZF3.1' [FMT .1×111)LABELD 1.55,(1.75+.475×111),[.5]
[11]
[12]
      ('LI4' ∏FMT 4096×.1×111)LABELD 9.05,(1.75+.475×111),[
     .5]0
      (*LI4* [FMT 512×18)LABELD(2+.875×18),1.55,[.5]8p0
      'TIME IN MILLISECONDS FROM BEGINNING OF DATA WINDOW'
[14]
     CENTER 1 9 1.3 0
[15]
      (1 20p'NORMALIZED INTENSITY')LABELD 1 3pl.4 2.8 90
[16]
     (1 6p*COUNTS*)LABELD 1 3p9.8 4.9 270
     1 SPLINE(8192p0 0 1 0)/(2+7×(14096)+4095).1.75+4.75×X
[17]
     *4095
[18]
     PLOTEND
     V
    V Y+N SMOOTH D:□IO
 [1]
     AN-POINT UNWEIGHTED SMOOTHING OF THE VECTOR D
 [2] DIO+0
 [3]
      Y+(\rho D)\rho((\lfloor .5 \times N) + D), ((-N)+(+/(1N)\phi(N,\rho D)\rho D) + N), (-[.5 \times N)
     ) + D
```

V

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BIOGRAPHICAL SKETCH

Glenn H Schneider (no "." after the H, please note) was born to Elaine and Ira G. Schneider on October 12, 1955, an otherwise un-noteworthy day, except perhaps for the fact that it marked the 463rd anniversary of the discovery of America, as commonly reckoned. At the tender age of three he received his first telescope (Sears, 40X!), and learned early in life the meaning of light pollution.

He attended public school 97 and J.H.S. 135 (also known as the Frank D. Whalen Junior High School, though no one ever knew who Frank D. Whalen was) in New York City. Thanks to the closing of the school during the summer of 1969 he was permitted to retain possession of the school library's copies of <u>The Larousse Encyclopedia of Astronomy</u> and <u>Norton's Star Atlas</u>. Later that year he became a member of the Amateur Observers' Society of New York. These two events forever influenced his life and started him on the path to a career in astronomy.

He graduated from The Bronx High School of Science in June of 1972. Shortly thereafter, he organized an international solar eclipse expedition, in hopes of observing his second total solar eclipse. The clouds which hung ominously over the Gaspé Peninsula that day served to reinforce his devotion to shadow chasing.

In June of 1976, he earned a Bachelor of Science degree in physics from the New York Institute of Technology. Before entering the graduate program in astronomy at the University of Florida in September of 1977, he worked as an APL technical analyst for Warner Computer Systems, Inc.

After eight years, two trips to the South Pole, seemingly endless commuting between Gainesville, Paris, Noordwijc and Bristol, statistically improbable stretches of cloudy weather during scheduled observing runs, and being thwarted at every turn by wayward computers (micro and macro), he expects, finally, to receive the degree of Doctor of Philosophy in August of 1985. He has accepted a position at the Space Telescope Science Institute, working for the Computer Sciences Corporation as a Science and Mission Operations Astronomer.

He has observed every total solar eclipse since March, 1970, and intends to Keep chasing the moon's shadow well into the future. He has a definite preference for Dr. Brown's Celery Tonic (now called Cel-Ray Soda, alas), and egg creams made with Fox's U-Bet chocolate syrup and Good-Health seltzer.

At the moment he is single, but this malady is expected to be cured on June 7, 1985.

I certify that I have read this study and that in my opinion it conforms to acceptable standards of scholarly presentation and is fully adequate, in scope and quality, as a dissertation for the degree of Doctor of Philosophy.

John P. Dliver, Chairman Associate Professor of Astronomy

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Heinrich Eichhorn Professor of Astronomy

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Howard L. Cohen Associate Professor of Astronomy

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Héywood C. Smith' Associate Professor of Astronomy I certify that I have read this study and that in my opinion it conforms to acceptable standards of scholarly presentation and is fully adequate, in scope and quality, as a dissertation for the degree of Doctor of Philosophy.

Ralph G. Selfridge Professor of Computer and Information Sciences

This dissertation was submitted to the Graduate Faculty of the Department of Astronomy in the College of Liberal Arts and Sciences and to the Graduate School and was accepted as partial fulfillment of the requirements for the degree of Doctor of Philosophy.

August 1985

Dean, Graduate School